

# Phase Space Propagation

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## Creating a Distribution

Consider a collection of  $N_{par}$  particles, each with a transverse coordinate  $x$  lying on a corresponding trajectory with slope  $x'$  relative to the ideal trajectory. Suppose the particles are distributed “normally” about the ideal trajectory, with rms spreads in position and slope of  $\sigma_x$  and  $\sigma_{x'}$ , respectively. To model this situation, we call up **R** on the computer and enter the following commands into a console window:

```
Npar = 10000
x = rnorm(Npar,0, 5)
xp = rnorm(Npar,0,2)
```

The “random normal” function in **R** creates coordinates  $x$  from a normal, or *Gaussian*, distribution with a desired mean of zero and an rms spread of  $\sigma_x = 5$ . (Let’s presume that the units are in mm.) The distribution of slopes  $x'$  should have an rms spread of  $\sigma_{x'} = 10$ . (Here, let’s presume that the units are in  $\text{mr} = 10^{-3}$  radians.)

We can see how well we have done by finding the rms of the above variable distributions, using the `var()` function (*variance*) in **R**, which finds  $\langle(x - \langle x \rangle)^2\rangle$ :

```
sqrt(var(x))
sqrt(var(xp))
```

We can also plot simple histograms of the distributions in  $x$  and  $x'$ , as well as a “phase space plot” as follows:

```
hist(x)
hist(xp)
plot(x,xp,pch=".")
```

As will be presented in class, and as can be found in the on-line notes (see Appendices on the web site), the phase space distribution can be characterized by a set of parameters, namely an “emittance”,  $\epsilon$  – which is related to the area in phase space occupied by the particles – and the Courant-Snyder parameters,  $\alpha$ ,  $\beta$ , and  $\gamma$ .

In addition to `var(x)` for finding the variance  $\langle x^2 \rangle$  of a variable, the covariance  $\langle xy \rangle$  of two variables is found using `cov(x,y)`. The function `CScalc()` defined in the code below computes the Courant-Snyder parameters and the emittance using the variance and covariance functions applied to the  $x$  and  $x'$  distributions created earlier.

```
CScalc = function(x,xp){
  epsx=sqrt(var(x)*var(xp)-cov(x,xp)^2)
  betx= var(x)/epsx
  alfx=-cov(x,xp)/epsx
  gamx= var(xp)/epsx
  c(alfx,betx,gamx,epsx)
}
```

A simple call to this function prints out four numbers:  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $\epsilon$ :

```
CScalc(x,xp)
```

- How does the emittance,  $\epsilon$ , found from `CScalc()` relate to the initial rms values  $\sqrt{\langle x^2 \rangle}$  and  $\sqrt{\langle x'^2 \rangle}$  generated earlier?

- The value of the covariance `cov(x, xp)` should be nearly zero for our example; check that this is so. What would a non-zero value tell you about the phase space distribution?

```
cov(x, xp)
```

## Propagating the Distribution

We know that as the particles move down the beam line unencumbered through a distance  $L$ , the angles of their trajectories will not change, but their transverse positions will change according to

$$x \leftarrow x + x' \cdot L.$$

Input the following code into your **R** program to see how the distribution changes...

```
# drift a distance L
L = 3
x = x + xp*L

hist(x)
hist(xp)
plot(x, xp, pch=".")
CScalc(x, xp)
```

You should notice that the Courant-Snyder parameters have changed; the value of  $\alpha$  changes significantly, indicating a correlation developing between  $x$  and  $x'$ .

- Why would this happen?
- How has the value of  $\epsilon$  changed (or has it)?

Next, since we see the extent of the  $x$  distribution getting very large, let's go through a thin lens quadrupole of focal length  $F$ . The magnet will alter the slopes of the trajectories according to

$$x' \leftarrow x' - x/F$$

while leaving the positions  $x$  unchanged (in the “thin lens” approximation). The code below performs this transformation and generates results.

```
# focus through a quadrupole, focal length F
F = 2
xp = xp - x/F

hist(x)
hist(xp)
plot(x, xp, pch=".")
CScalc(x, xp)
```

Notice that the correlation has changed – in fact, the covariance has changed sign for our chosen parameters.

## Create a system of lenses

We now have the tools to model the propagation of our distribution through a series of lenses. Let's create two lists, one a list of lens focal lengths  $F$  and the second a list of the spaces between these lenses,  $L$ . Let's assume these distances to be in meters (and remember,  $x$  is in mm and  $x'$  is in mrad).

```
L = c( 5, 2, 7, 4, 3, 3, 6, 8)
F = c(-8, 4, -9, 4, -7, 9, -5, 7)
```

We now write a code that models a particle being focused through a thin lens of focal length  $F_i$ , then propagating a distance  $L_i$ . We then loop this operation through the entire list of  $N$  lenses, where  $i \in [1, N]$ . At the end of each drift we compute the Courant-Snyder variables and keep track of the results. In the code below, the vector  $V$  is  $V = (\alpha, \beta, \gamma, \epsilon)$  and at the end of each loop we *bind* the new results with the previous results using the `rbind()` function (*bind by row*) in **R**. The final result will be a matrix of values where each row will be the values of  $V$  at the end of the  $i$ -th drift.

```
i = 0
V = CScalc(x,xp) # values at s=0
while(i < length(L)){
  i = i+1
  xp = xp-x/F[i]
  x = x+xp*L[i]
  V = rbind(V,CScalc(x,xp))
}
```

After executing the above code we can see the final result simply by typing

```
V
```

in the console window.

To ease the analysis we'll define new variables in order to capture the values found in each of the 4 rows of  $V$ :

```
alpha = V[,1]
beta = V[,2]
gamma = V[,3]
eps = V[,4]
```

Finally, we can make a plot of the beam envelope through our beam line. First, to get the path length coordinate, let  $s_{i+1} = s_i + L_i$ . The `cumsum()` function in **R** makes this fairly easy, though we must include the initial value of  $s = 0$  to get aligned with the proper number of values found in our loop above:

```
s = c(0,cumsum(L))
plot(s, beta ,typ="l",ylim=c(0,1.2*max(beta )))
plot(s,sqrt(beta),typ="l",ylim=c(0,1.2*max(sqrt(beta))))
```

- Why might we plot  $\sqrt{\beta}$ , as we did in the second plot, above rather than  $\beta$  ?
- Use **R** to plot the value of  $\beta\gamma - \alpha^2$  at each of the points in our beam line calculation.
- With the result of this plot, write the general relationship between  $\gamma$  and the other two Courant-Snyder parameters.