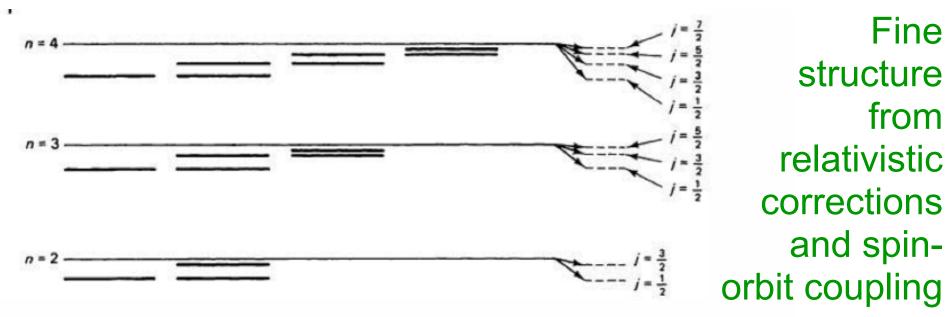
Let's move to Bound States

When we discuss bound states of two objects in central-force potential, kinetic energy and potential energy are ~the same. How does this compare to the rest energy of the objects?

Hydrogen ionization energy: 13.6 eV vs 0.5 MeV rest mass

Masses of b and c quarks are ~relatively large, so we can also consider them non-relativistically (which makes them much easier). NOT true for uds quarks

#### Briefest of overviews of hydrogen atom



Energy levels ~ -1/n<sup>2</sup>

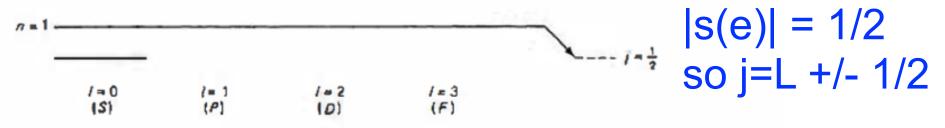
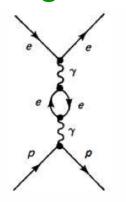


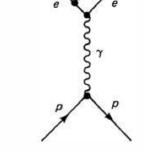
Fig. 5.2 Fine structure in hydrogen. The *n*th Bohr level (fine line) splits into *n* sublevels (dashed lines), characterized by  $j = \frac{1}{2}, \frac{3}{2}, \ldots, (n - \frac{1}{2})$ . Except for the last of these, two different values of *I* contribute to each level:  $l = j - \frac{1}{2}$  and  $l = j + \frac{1}{2}$ . Spectroscopists' nomenclature - S for l = 0, P for l = 1, D for l = 2, F for l = 3 - is indicated. All levels are shifted downward, as shown (the diagram is not to scale, however).

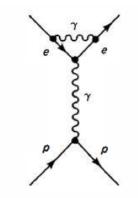
Griffiths

### Lamb shift

Lamb shift: Led to development of quantum electro-dynamics! QED corrections to the electronproton interaction break degeneracy of two levels with same n, j but different L (so  $2S_{1/2}$  and  $2P_{1/2}$  are not fully degenerate)







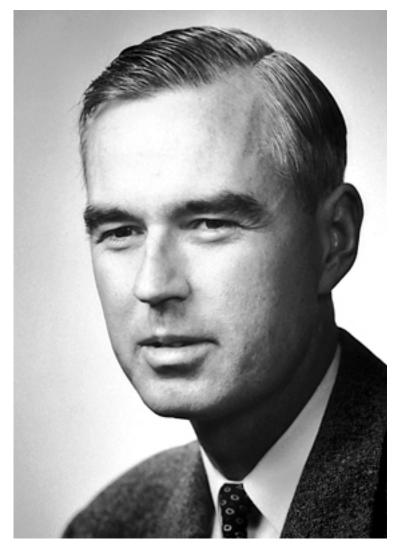
Vacuum polarization

Electron mass renormalization

Anomalous magnetic moment

Griffiths

Fig. 5.3 Some loop diagrams contributing to the Lamb shift.



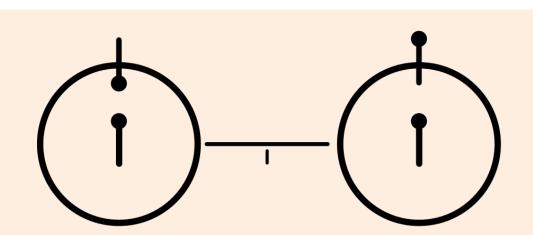
Willis Lamb

Spin-orbit coupling is principally due to spin of electron interacting with 'B field' from nucleus (fine structure). Much smaller is spin of nucleus interacting with 'B field' from electron. Goes as  $(m_e/m_p)^4$  hence hyperfine (and not fine) splitting.

For n=1, the difference in energy states of proton (e and p spins aligned vs anti-aligned, which is lower) is  $5.9 \,\mu eV = 1420$ MHz = 21 cm.Famous 21 cm line (penetrates dust!)

Lifetime of 21 cm is millions of years! Thankfully, enough hydrogen can provide this transition. Long lifetime = narrow width, so this is excellent for spectroscopy (Doppler shifts)

Used extensively in radio-astronomy, studying the early Universe, galaxy formation, measuring distances to objects, cosmology



Pioneer Plaque: 21 cm line defines distance and unit of time

Differences between quarkonium and hydrogen/ positronium: Don't really know the potential (strong force!) Also, interaction between quarks is large. Doesn't work for two light quarks, either

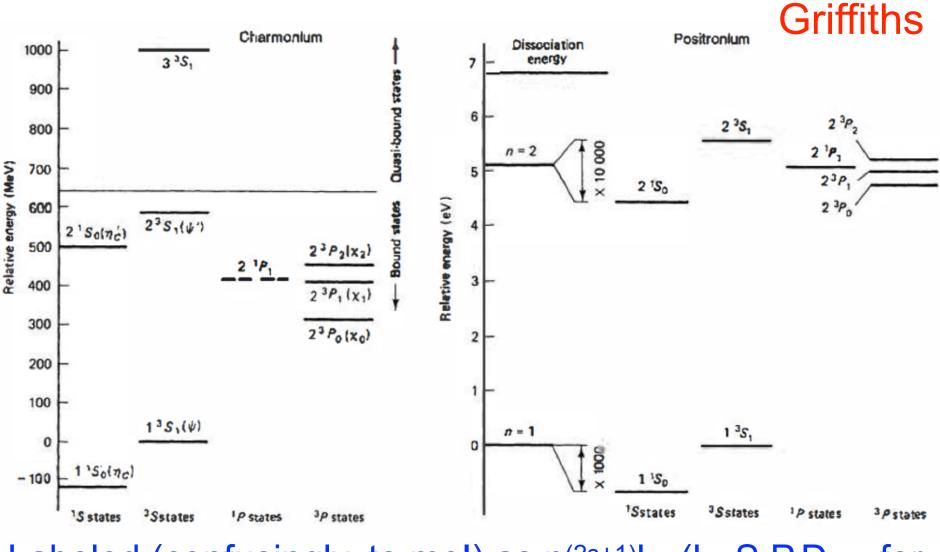
Instead of considering different states as energy levels of an atom, consider different bound states as different particles, each with a different mass. Start with mesons (much easier than baryons)

Q----Q'

At short distances, we know that QCD is not a strong force. Reasonable to start with 1/r potential.

At large distances, we know that force grows exponentially. Try V~kr (others could also work). Of course, k can be a function of r too!

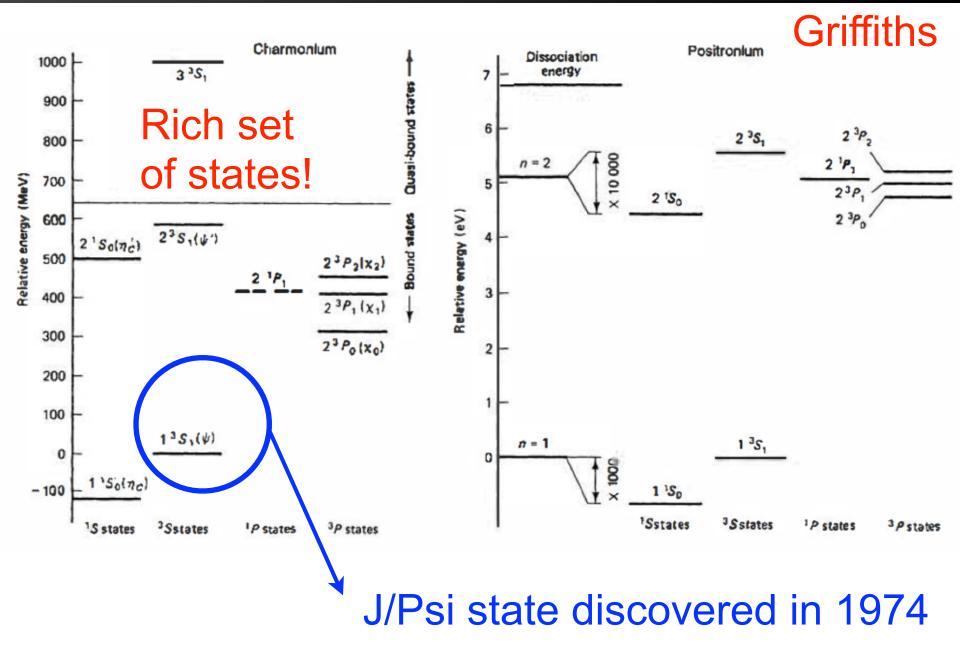
#### Charmonium



Labeled (confusingly, to me!) as  $n^{(2s+1)}L_J$  (L=S,P,D,... for 0,1,2... and s=0 or 1 for anti-aligned or aligned spins), with J=L+s

#### Charmonium





#### Why was J/Psi discovered first?

 $\eta$ 

 $I^{G}(J^{PC}) = 0^{+}(0^{-+})$ 

		Charged modes	
Г <sub>8</sub>	charged modes	(28.06±0.34) %	S=1.2
Г9	$\pi^+\pi^-\pi^0$	(22.73±0.28) %	S=1.2
Γ <sub>10</sub>	$\pi^+\pi^-\gamma$	( 4.60±0.16) %	S=2.1
$\Gamma_{11}$	$e^+e^-\gamma$	( 6.8 $\pm$ 0.8 ) $ imes$ 10 $^{-3}$	S=1.7
Γ <sub>12</sub>	$\mu^+\mu^-\gamma$	$(3.1 \pm 0.4)  imes 10^{-4}$	
Γ <sub>13</sub>	e <sup>+</sup> e <sup>-</sup>	$< 7.7 \times 10^{-5}$	CL=90%
Γ <sub>14</sub>	$\mu^+\mu^-$	$(5.8 \pm 0.8)  imes 10^{-6}$	
Γ <sub>15</sub>	e <sup>+</sup> e <sup>-</sup> e <sup>+</sup> e <sup>-</sup>	$< 6.9 \times 10^{-5}$	CL=90%
Г <sub>16</sub>	$\pi^+\pi^-e^+e^-$	$(4.2 \pm 1.2) \times 10^{-4}$	
Γ <sub>17</sub>	$\pi^+\pi^-2\gamma$	$< 2.0 \times 10^{-3}$	
Г <sub>18</sub>	$\pi^+\pi^-\pi^0\gamma$	$< 5 \times 10^{-4}$	CL=90%
Γ <sub>19</sub>	$\pi^{0} \mu^{+} \mu^{-} \gamma$	$< 3 \times 10^{-6}$	CL=90%

## From PDG

Due to C-parity!

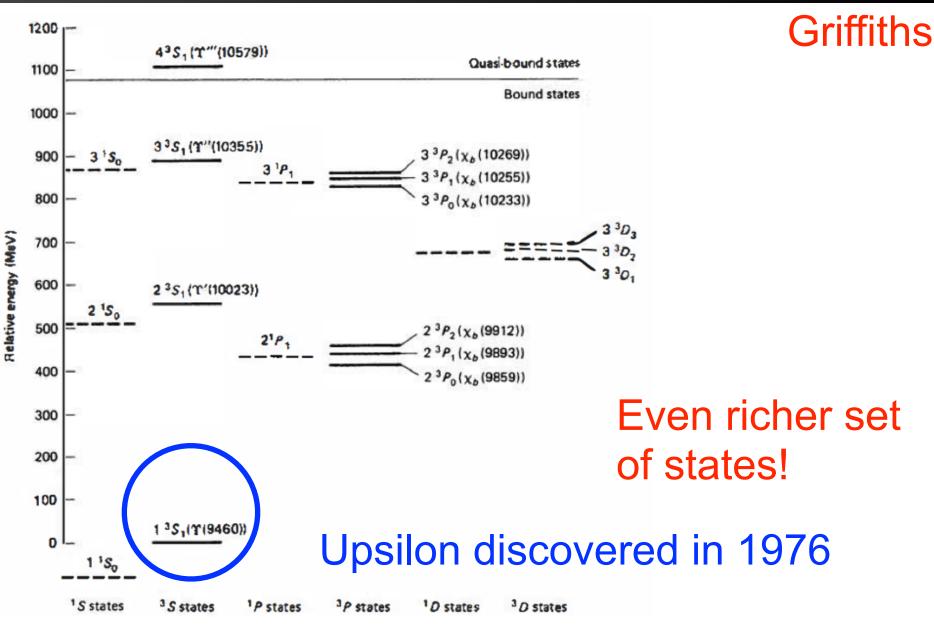
$$J/\psi(1S)$$

$$I^{G}(J^{PC}) = 0^{-}(1^{-})$$

#### $J/\psi(1S)$ DECAY MODES

	Mode	Fraction $(\Gamma_i/\Gamma)$	Scale factor/ Confidence level
Γ <sub>1</sub>	hadrons	(87.7 ±0.5 )%	
Γ <sub>2</sub>	$virtual\gamma  o hadrons$	$(13.50 \pm 0.30)$ %	
۲ <sub>3</sub>	ggg	(64.1 $\pm 1.0$ ) %	
Γ <sub>4</sub>	$\gamma g g$	(8.8 $\pm$ 0.5 )%	
Γ <sub>5</sub>	e <sup>+</sup> e <sup>-</sup>	( 5.94 $\pm 0.06$ )%	
Г <sub>6</sub>	$\mu^+\mu^-$	( 5.93 $\pm 0.06$ )%	

#### **Bottomonium**



Recall magnetic moment formula:  $\mu = -\frac{e}{m}\mathbf{S}$ 

Spin-spin interactions in hadrons have two components:

$$\mu_1 \cdot \mathbf{S_2} = -\frac{e}{m_1} \mathbf{S_1} \cdot \mathbf{S_2}$$
$$\mu_2 \cdot \mathbf{S_1} = -\frac{e}{m_2} \mathbf{S_2} \cdot \mathbf{S_1} = -\frac{e}{m_2} \mathbf{S_1} \cdot \mathbf{S_2}$$

Sum is then

$$-e\frac{m_1 + m_2}{m_1 m_2} \left( \mathbf{S_1} \cdot \mathbf{S_2} \right) = A(m_1 + m_2) \frac{1}{m_1 m_2} \left( \mathbf{S_1} \cdot \mathbf{S_2} \right)$$

$$(\mathbf{S_1} \cdot \mathbf{S_2}), \mathbf{S} = \mathbf{S_1} + \mathbf{S_2}$$
  
 $\mathbf{S}^2 = \mathbf{S_1}^2 + \mathbf{S_2}^2 + 2(\mathbf{S_1} \cdot \mathbf{S_2})$   
 $(\mathbf{S_1} \cdot \mathbf{S_2}) = \frac{1}{2} (\mathbf{S}^2 - \mathbf{S_1}^2 - \mathbf{S_2}^2)$ 

S<sub>1</sub> and S<sub>2</sub> = ±1/2  $S_1^2 = S_2^2 = (1/2)(1/2+1) = 3/4$   $S^2 = (1)(1+1) = 2$  (spins aligned) or  $S^2 = (0)(0+1) = 0$  (spins anti-aligned) So: S<sub>1</sub>·S<sub>2</sub> = 1/4 (spins aligned) or: S<sub>1</sub>·S<sub>2</sub> = -3/4 (spins anti-aligned)

$$M(m_1 - - m_2) = m_1 + m_2 + A(m_1 + m_2) \frac{1}{m_1 m_2} \left( \mathbf{S_1} \cdot \mathbf{S_2} \right)$$

- Mass M of
- meson

this

composed of quarks with mass m<sub>1</sub> and m<sub>2</sub> then generically looks like  $S_1 \cdot S_2 = 1/4$  (spins aligned) or:  $S_1 \cdot S_2 = -3/4$  (spins anti-aligned)

> Can try something simpler, and assume A is a constant (it surely is not, but maybe that is a reasonable approximation)

#### Fits for masses (from Griffiths)

Table 5.3 Pseudo scalar and vector meson masses.  $(MeV/c^2)$ 

Meson	Calculated	Observed	
π	139	138	
ĸ	487	496	
η	561	548	
p	775	776	
w	775	783	
K*	892	894	
<b>\$</b>	1031	1020	

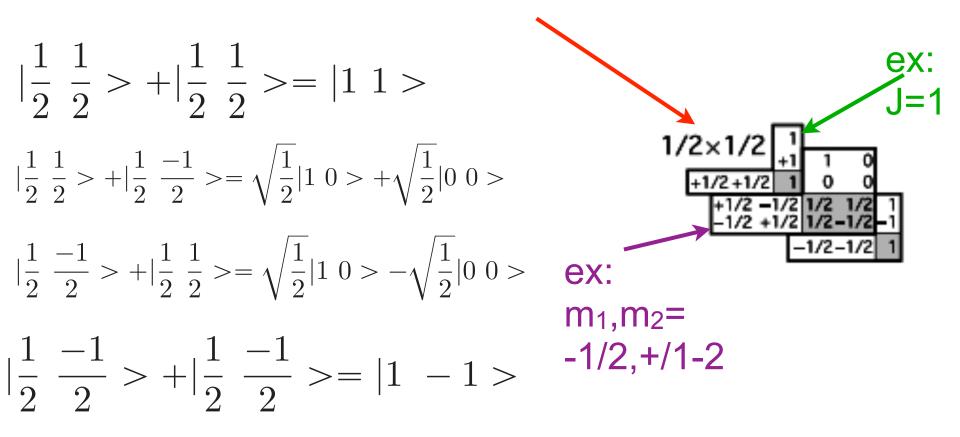
Very nice agreement! But need to be careful... For example:  $\eta = \frac{u\overline{u} + d\overline{d} - 2s\overline{s}}{\sqrt{6}}$  A lot more complicated - have three quarks, and thus three spins to add together. Most importantly, mesons are always composed of a quark and an anti-quark, ie never contain two of the same particle. In baryons, however (example: proton = uud), this no longer has to be true.

Regardless, though, baryons have half-integer spin (three quarks with s=+/- 1/2 can combine to give s= +/- 1/2 or +/- 3/2 only)

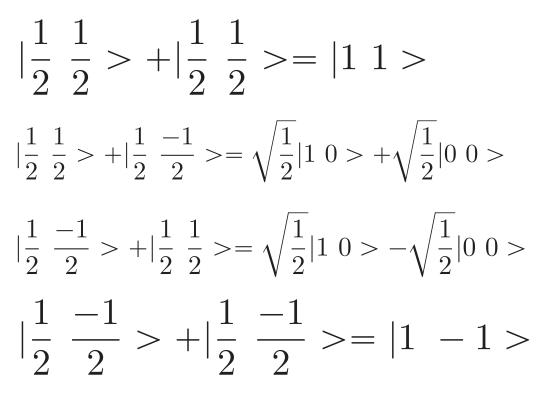
#### How to add three spins

To add three spins together, we first start by adding two of them together. Back to those C-G tables from the PDG ...

Combining two 1/2 x 1/2 particles



#### Let's rearrange



These are the  $|1 \ 1 >= |\frac{1}{2} \ \frac{1}{2} > + |\frac{1}{2} \ \frac{1}{2} >$ easy ones  $|1 \ -1 >= |\frac{1}{2} \ \frac{-1}{2} > + |\frac{1}{2} \ \frac{-1}{2} >$  Let's rearrange (can also use tables for this)

$$\begin{aligned} \left|\frac{1}{2} \ \frac{1}{2} > + \left|\frac{1}{2} \ \frac{1}{2} > = \left|1 \ 1 > \right. \right. \\ \left|\frac{1}{2} \ \frac{1}{2} > + \left|\frac{1}{2} \ \frac{-1}{2} > = \sqrt{\frac{1}{2}} \right|^{1} \ 0 > + \sqrt{\frac{1}{2}} \left|0 \ 0 > \right. \end{aligned}$$

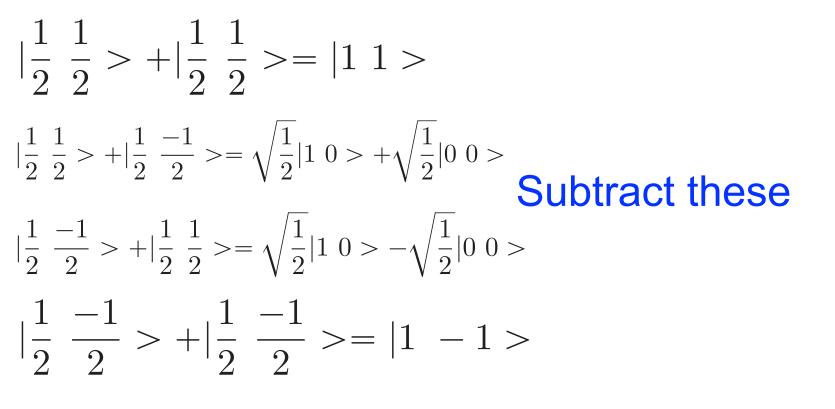
$$\begin{aligned} \left|\frac{1}{2} \ \frac{-1}{2} > + \left|\frac{1}{2} \ \frac{1}{2} > = \sqrt{\frac{1}{2}} \left|1 \ 0 > -\sqrt{\frac{1}{2}} \left|0 \ 0 > \right. \end{aligned}$$

$$\begin{aligned} \left|\frac{1}{2} \ \frac{-1}{2} > + \left|\frac{1}{2} \ \frac{1}{2} > = \sqrt{\frac{1}{2}} \left|1 \ 0 > -\sqrt{\frac{1}{2}} \left|0 \ 0 > \right. \end{aligned}$$

$$\begin{aligned} \left|\frac{1}{2} \ \frac{-1}{2} > + \left|\frac{1}{2} \ \frac{-1}{2} > = \left|1 \ -1 > \right. \end{aligned}$$

 $\sqrt{2}|1\ 0\rangle = \left(|\frac{1}{2}\ \frac{1}{2}\ \rangle + |\frac{1}{2}\ \frac{-1}{2}\ \rangle\right) + \left(|\frac{1}{2}\ \frac{-1}{2}\ \rangle + |\frac{1}{2}\ \frac{1}{2}\ \rangle\right)$  $|1\ 0\rangle = \frac{1}{\sqrt{2}}\left(|\frac{1}{2}\ \frac{1}{2}\ \rangle + |\frac{1}{2}\ \frac{-1}{2}\ \rangle\right) + \frac{1}{\sqrt{2}}\left(|\frac{1}{2}\ \frac{-1}{2}\ \rangle + |\frac{1}{2}\ \frac{1}{2}\ \rangle\right)$ 

#### Let's rearrange (can also use tables for this)

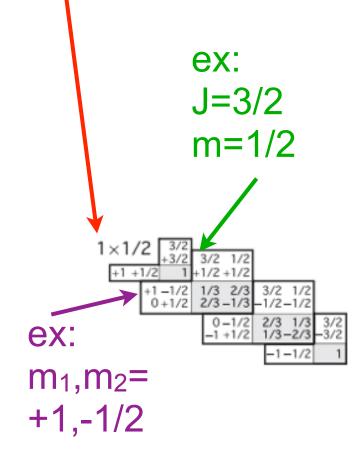


 $\sqrt{2}|0 \ 0 > = \left(|\frac{1}{2} \ \frac{1}{2} > +|\frac{1}{2} \ \frac{-1}{2} > \right) - \left(|\frac{1}{2} \ \frac{-1}{2} > +|\frac{1}{2} \ \frac{1}{2} > \right)$  $|0 \ 0 > = \frac{1}{\sqrt{2}} \left(|\frac{1}{2} \ \frac{1}{2} > +|\frac{1}{2} \ \frac{-1}{2} > \right) - \frac{1}{\sqrt{2}} \left(|\frac{1}{2} \ \frac{-1}{2} > +|\frac{1}{2} \ \frac{1}{2} > \right)$ 

When we add the third quark we will have to add spin 1/2 with either spin 0 or spin 1

## Combining spin 1 x 1/2 particles

We want the "inverse" of what we have been reading off. Can also use the tables for that!



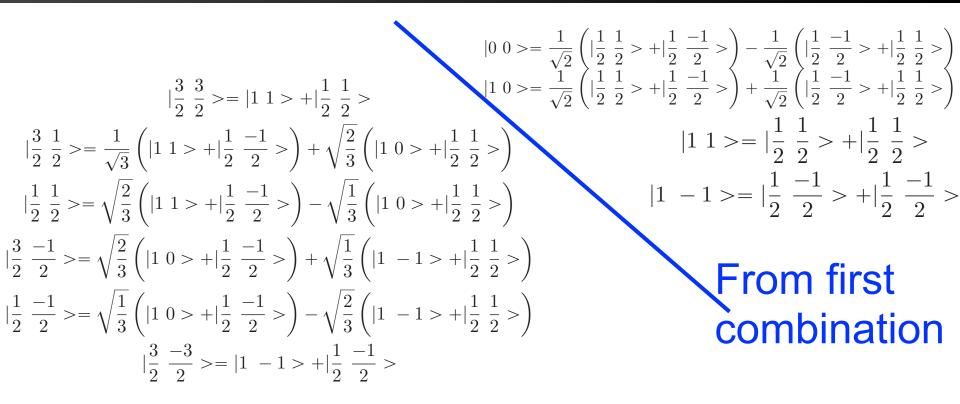
## Combining spin 1 x 1/2 particles

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$$\begin{aligned} |\frac{3}{2} \frac{3}{2} \rangle &= |1 1 \rangle + |\frac{1}{2} \frac{1}{2} \rangle \\ |\frac{3}{2} \frac{1}{2} \rangle &= \frac{1}{\sqrt{3}} \left( |1 1 \rangle + |\frac{1}{2} \frac{-1}{2} \rangle \right) + \sqrt{\frac{2}{3}} \left( |1 0 \rangle + |\frac{1}{2} \frac{1}{2} \rangle \right) \\ |\frac{1}{2} \frac{1}{2} \rangle &= \sqrt{\frac{2}{3}} \left( |1 1 \rangle + |\frac{1}{2} \frac{-1}{2} \rangle \right) - \sqrt{\frac{1}{3}} \left( |1 0 \rangle + |\frac{1}{2} \frac{1}{2} \rangle \right) \\ |\frac{3}{2} \frac{-1}{2} \rangle &= \sqrt{\frac{2}{3}} \left( |1 0 \rangle + |\frac{1}{2} \frac{-1}{2} \rangle \right) + \sqrt{\frac{1}{3}} \left( |1 - 1 \rangle + |\frac{1}{2} \frac{1}{2} \rangle \right) \\ |\frac{1}{2} \frac{-1}{2} \rangle &= \sqrt{\frac{1}{3}} \left( |1 0 \rangle + |\frac{1}{2} \frac{-1}{2} \rangle \right) - \sqrt{\frac{2}{3}} \left( |1 - 1 \rangle + |\frac{1}{2} \frac{1}{2} \rangle \right) \\ |\frac{3}{2} \frac{-3}{2} \rangle &= |1 - 1 \rangle + |\frac{1}{2} \frac{-1}{2} \rangle \end{aligned}$$

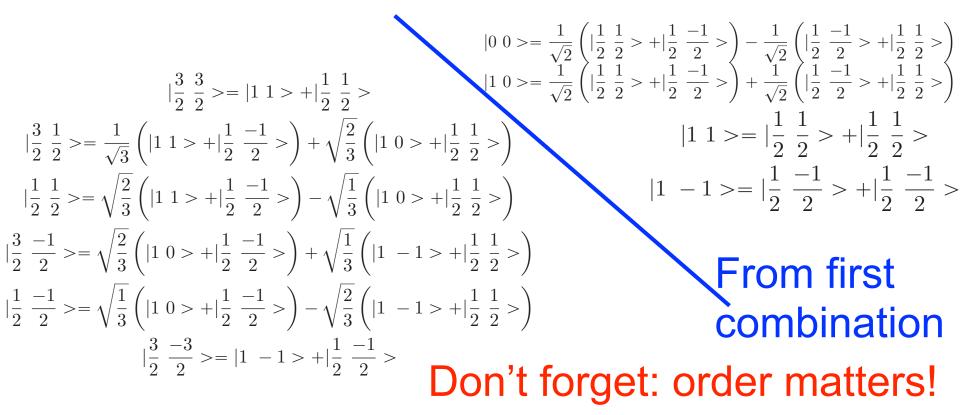
# Combining spin 0 x 1/2 particles is trivial

$$\begin{aligned} \left|\frac{1}{2} \frac{1}{2} \right| &>= \left|0 \ 0 > + \right| \frac{1}{2} \frac{1}{2} > \\ \left|\frac{1}{2} \frac{-1}{2} \right| &>= \left|0 \ 0 > + \right| \frac{1}{2} \frac{-1}{2} > \end{aligned}$$



$$\begin{aligned} |0 \ 0>+|\frac{1}{2} \ \frac{1}{2}> &= |\frac{1}{2} \ \frac{1}{2}> \\ |0 \ 0>+|\frac{1}{2} \ \frac{-1}{2}> &= |\frac{1}{2} \ \frac{-1}{2}> \end{aligned}$$

#### Let's introduce some nicer notation



$$|0 \ 0 > +|\frac{1}{2} \ \frac{1}{2} > = |\frac{1}{2} \ \frac{1}{2} > |\frac{1}{2} > |\frac{1}{2} > |\frac{1}{2} > |\frac{1}{2} > |\frac{1}{2} - \frac{1}{2} > |\frac{1}{2} \ \frac{1}{2} \ \frac{1}{2} > |\frac{1}{2} \ \frac{1}{2} \ \frac{1}{2} \ \frac{1}{2} > |\frac{1}{2} \ \frac{1}{2} \ \frac$$

$$|\frac{1}{2} \frac{1}{2} >= (\uparrow) \\ |\frac{1}{2} \frac{-1}{2} >= (\downarrow)$$

## Using the notation

$$\begin{split} |\frac{3}{2} \frac{3}{2} \rangle &= |1 1 \rangle + (\uparrow) \\ |\frac{3}{2} \frac{3}{2} \rangle &= \frac{1}{\sqrt{3}} \left( |1 1 \rangle + (\downarrow) \right) + \sqrt{\frac{2}{3}} \left( |1 0 \rangle + (\uparrow) \right) \\ |\frac{3}{2} \frac{1}{2} \rangle &= \frac{1}{\sqrt{3}} \left( |1 1 \rangle + (\downarrow) \right) - \sqrt{\frac{1}{3}} \left( |1 0 \rangle + (\uparrow) \right) \\ |\frac{1}{2} \frac{1}{2} \rangle &= \sqrt{\frac{2}{3}} \left( |1 1 \rangle + (\downarrow) \right) - \sqrt{\frac{1}{3}} \left( |1 0 \rangle + (\uparrow) \right) \\ |\frac{3}{2} \frac{-1}{2} \rangle &= \sqrt{\frac{2}{3}} \left( |1 0 \rangle + (\downarrow) \right) + \sqrt{\frac{1}{3}} \left( |1 - 1 \rangle + (\uparrow) \right) \\ |\frac{1}{2} \frac{-1}{2} \rangle &= \sqrt{\frac{1}{3}} \left( |1 0 \rangle + (\downarrow) \right) - \sqrt{\frac{2}{3}} \left( |1 - 1 \rangle + (\uparrow) \right) \\ |\frac{3}{2} \frac{-3}{2} \rangle &= |1 - 1 \rangle + (\downarrow) \end{split}$$

$$\begin{aligned} |\frac{1}{2} \ \frac{1}{2} > &= |0 \ 0 > + \uparrow \\ |\frac{1}{2} \ \frac{-1}{2} > &= |0 \ 0 > + \downarrow \end{aligned}$$

Phew

Putting it together

$$\begin{split} \left|\frac{3}{2} \ \frac{3}{2} \right| &\geq = \uparrow \uparrow \uparrow \\ \left|\frac{3}{2} \ \frac{1}{2} \right| &\geq = \frac{1}{\sqrt{3}} (\uparrow \downarrow \downarrow) + \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow \downarrow + \downarrow \uparrow) \uparrow\right) \\ \left|\frac{1}{2} \ \frac{1}{2} \right| &\geq = \sqrt{\frac{2}{3}} (\uparrow \uparrow \downarrow) - \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow \downarrow + \downarrow \uparrow) \uparrow\right) \\ \left|\frac{3}{2} \ \frac{-1}{2} \right| &\geq = \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow \downarrow + \downarrow \uparrow) (\downarrow)\right) + \sqrt{\frac{1}{3}} (\downarrow \downarrow \uparrow) \\ \left|\frac{1}{2} \ \frac{-1}{2} \right| &\geq = \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow \downarrow + \downarrow \uparrow) \downarrow\right) - \sqrt{\frac{2}{3}} (\downarrow \downarrow \uparrow) \\ \left|\frac{3}{2} \ \frac{-3}{2} \right| &\geq = \downarrow \downarrow \downarrow \end{split}$$

$$\begin{aligned} |\frac{1}{2} \ \frac{1}{2} > &= \frac{1}{\sqrt{2}} \left(\uparrow \downarrow - \downarrow \uparrow\right) \uparrow \\ \frac{1}{2} \ \frac{-1}{2} > &= \frac{1}{\sqrt{2}} \left(\uparrow \downarrow - \downarrow \uparrow\right) \downarrow \end{aligned}$$

#### How to interpret

$$\begin{split} |\frac{3}{2} \ \frac{3}{2} >=\uparrow\uparrow\uparrow\\ |\frac{3}{2} \ \frac{1}{2} >= \frac{1}{\sqrt{3}} (\uparrow\uparrow\downarrow) + \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)\uparrow\right)\\ |\frac{1}{2} \ \frac{1}{2} >= \sqrt{\frac{2}{3}} (\uparrow\uparrow\downarrow) - \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)\uparrow\right)\\ |\frac{3}{2} \ \frac{-1}{2} >= \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)(\downarrow)\right) + \sqrt{\frac{1}{3}} (\downarrow\downarrow\uparrow)\\ |\frac{1}{2} \ \frac{-1}{2} >= \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)\downarrow\right) - \sqrt{\frac{2}{3}} (\downarrow\downarrow\uparrow)\\ |\frac{3}{2} \ \frac{-3}{2} >=\downarrow\downarrow\downarrow$$

$$\begin{aligned} |\frac{1}{2} \ \frac{1}{2} > &= \frac{1}{\sqrt{2}} \left(\uparrow \downarrow - \downarrow \uparrow\right) \uparrow \\ |\frac{1}{2} \ \frac{-1}{2} > &= \frac{1}{\sqrt{2}} \left(\uparrow \downarrow - \downarrow \uparrow\right) \downarrow \end{aligned}$$

$$\begin{aligned} |\frac{3}{2} \ \frac{3}{2} > = \uparrow \uparrow \uparrow \\ |\frac{3}{2} \ \frac{1}{2} > = \frac{1}{\sqrt{3}} \left(\uparrow \uparrow \downarrow + \uparrow \downarrow \uparrow + \downarrow \uparrow \uparrow\right) \\ |\frac{3}{2} \ \frac{-1}{2} > = \frac{1}{\sqrt{3}} \left(\uparrow \downarrow \downarrow + \downarrow \uparrow \downarrow + \downarrow \downarrow \uparrow\right) \\ |\frac{3}{2} \ \frac{-3}{2} > = \downarrow \downarrow \downarrow \end{aligned}$$

Spin 3/2 states are easy to interpret: symmetric if we interchange any two quarks

#### How to interpret

$$\begin{split} |\frac{3}{2} \ \frac{3}{2} >=\uparrow\uparrow\uparrow\\ |\frac{3}{2} \ \frac{1}{2} >= \frac{1}{\sqrt{3}} (\uparrow\uparrow\downarrow) + \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)\uparrow\right)\\ |\frac{1}{2} \ \frac{1}{2} >= \sqrt{\frac{2}{3}} (\uparrow\uparrow\downarrow) - \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)\uparrow\right)\\ |\frac{3}{2} \ \frac{-1}{2} >= \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)(\downarrow)\right) + \sqrt{\frac{1}{3}} (\downarrow\downarrow\uparrow)\\ |\frac{1}{2} \ \frac{-1}{2} >= \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)(\downarrow)\right) - \sqrt{\frac{2}{3}} (\downarrow\downarrow\uparrow)\\ |\frac{3}{2} \ \frac{-3}{2} >=\downarrow\downarrow\downarrow$$

$$\begin{aligned} |\frac{1}{2} \ \frac{1}{2} > &= \frac{1}{\sqrt{2}} \left(\uparrow \downarrow - \downarrow \uparrow\right) \uparrow \\ |\frac{1}{2} \ \frac{-1}{2} > &= \frac{1}{\sqrt{2}} \left(\uparrow \downarrow - \downarrow \uparrow\right) \downarrow \end{aligned}$$

$$\begin{aligned} |\frac{1}{2} \ \frac{1}{2} > = \sqrt{\frac{2}{3}} (\uparrow \uparrow \downarrow) - \sqrt{\frac{1}{3}} \left( \frac{1}{\sqrt{2}} (\uparrow \downarrow + \downarrow \uparrow) \uparrow \right) \\ |\frac{1}{2} \ \frac{-1}{2} > = \sqrt{\frac{1}{3}} \left( \frac{1}{\sqrt{2}} (\uparrow \downarrow + \downarrow \uparrow) \downarrow \right) - \sqrt{\frac{2}{3}} (\downarrow \downarrow \uparrow) \end{aligned}$$

Two of spin 1/2 states are asymmetric under interchange of first and second quarks

#### How to interpret

$$\begin{split} |\frac{3}{2} \ \frac{3}{2} >=\uparrow\uparrow\uparrow\\ |\frac{3}{2} \ \frac{1}{2} >= \frac{1}{\sqrt{3}} (\uparrow\uparrow\downarrow) + \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)\uparrow\right)\\ |\frac{1}{2} \ \frac{1}{2} >= \sqrt{\frac{2}{3}} (\uparrow\uparrow\downarrow) - \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)\uparrow\right)\\ |\frac{3}{2} \ \frac{-1}{2} >= \sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)(\downarrow)\right) + \sqrt{\frac{1}{3}} (\downarrow\downarrow\uparrow)\\ |\frac{1}{2} \ \frac{-1}{2} >= \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} (\uparrow\downarrow+\downarrow\uparrow)\downarrow\right) - \sqrt{\frac{2}{3}} (\downarrow\downarrow\uparrow)\\ |\frac{3}{2} \ \frac{-3}{2} >=\downarrow\downarrow\downarrow$$

$$|\frac{1}{2} \frac{1}{2} \rangle = \frac{1}{\sqrt{2}} (\uparrow \downarrow - \downarrow \uparrow) \uparrow$$
$$|\frac{1}{2} \frac{-1}{2} \rangle = \frac{1}{\sqrt{2}} (\uparrow \downarrow - \downarrow \uparrow) \downarrow$$

$$\begin{split} |\frac{1}{2} \ \frac{1}{2} > &= \sqrt{\frac{2}{3}} \left(\uparrow\uparrow\downarrow\right) - \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} \left(\uparrow\downarrow+\downarrow\uparrow\right)\uparrow\right) \\ |\frac{1}{2} \ \frac{-1}{2} > &= \sqrt{\frac{1}{3}} \left(\frac{1}{\sqrt{2}} \left(\uparrow\downarrow+\downarrow\uparrow\right)\downarrow\right) - \sqrt{\frac{2}{3}} \left(\downarrow\downarrow\uparrow\right) \end{split}$$

These last two spin 1/2 states are symmetric under interchange of first and second quarks We need our 3 quarks to satisfy Fermi-Dirac statistics (must be anti-symmetric under exchange of any two quarks)

For ground state (I=0), space wave function is symmetric. Left off with wave functions for spin, color and flavor. We will see that color wave function is necessarily anti-symmetric. That means that flavor x spin combination must be symmetric

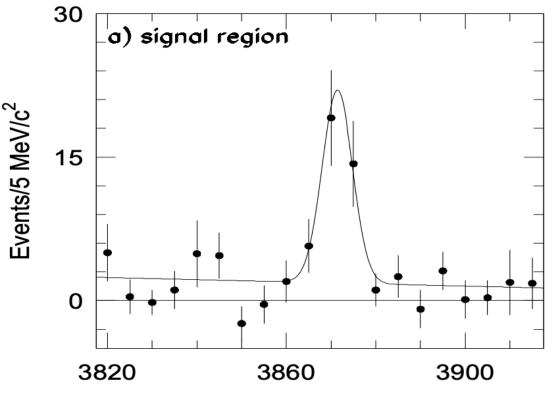
$$M(m_1m_2m_e) = m_1 + m_2 + m_3 + A'(\frac{(\mathbf{S_1} \cdot \mathbf{S_2})}{m_1m_2} + \frac{(\mathbf{S_1} \cdot \mathbf{S_3})}{m_1m_3} + \frac{(\mathbf{S_2} \cdot \mathbf{S_3})}{m_2m_3})$$

 $S_i \cdot S_j = 1/4$  (spins aligned) or:  $S_i \cdot S_j = -3/4$  (spins anti-aligned)

Using same formalism as we used for meson masses. Get a good fit, but find somewhat different quark masses in mesons and baryons (these are effective masses!) Conclusion: Much of the mass in, for example, the proton, comes from the energy of the QCD field and not from constituent quark masses themselves Why do hadrons come only in meson (q-qbar) and baryon (qqq) form? Why not also combinations of 4 quarks? Or 5 quarks?

You can combine two mesons to form a **tetraquark** that is colorless, or a meson and a baryon to form a colorless **pentaquark** 

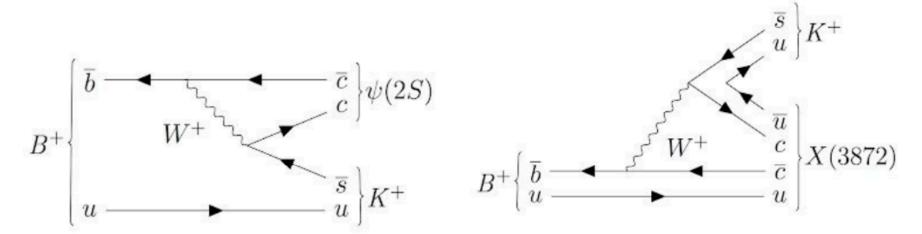
Such states seemed to come and go in terms of evidence for or against them but these days clearly seem to exist ...



X3872 observed by Belle (asymmetric e+e collider): first strong evidence in 2003 for an exotic quark! Intriguingly, mass is very close to the  $D^0\overline{D}^{0*}$  mass threshold: is a loosely bound "molecule"?

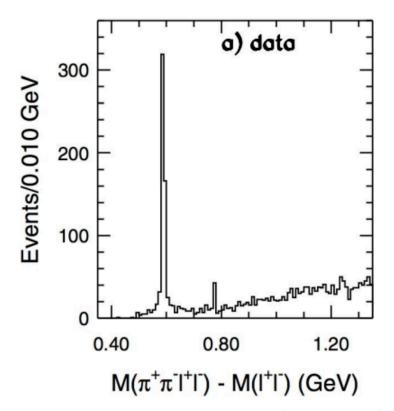
 $M(\pi^{+}\pi^{-}J/\psi)$  (MeV/c<sup>2</sup>)

#### 1808.04153



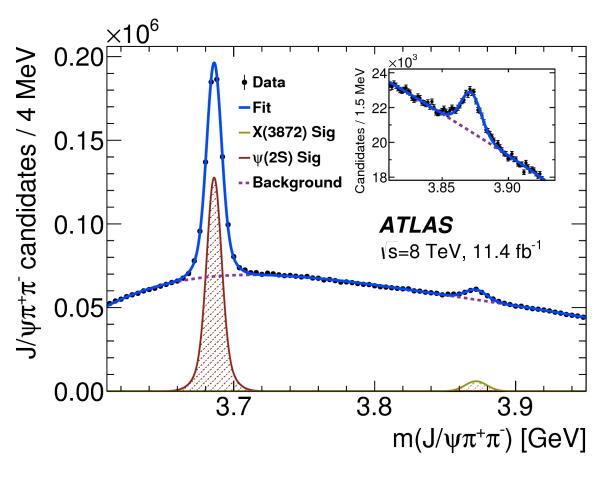
**Figure 6.** Feynman diagrams showing two different decays of a  $B^+$  meson. On the left is the decay to a final state containing the conventional  $\psi(2S)$  charmonium state and on the right is the decay to the exotic X(3872). The  $\psi(2S)$  and X(3872) both decay to  $J/\psi\pi^+\pi^-$ , such that the same particles are in the final state of both  $B^+$  meson decays.

$$\psi(2S) \to \pi^+ \pi^- J/\psi \to \pi^+ \pi^- l^+ l^-$$



**Figure 7.** Distribution of the difference between  $\pi^+\pi^-l^+l^-$  and  $l^+l^-$  invariant masses of  $B^+ \rightarrow J/\psi\pi^+\pi^-K^+$  candidates, where  $J/\psi \rightarrow l^+l^-$ , l being  $\mu$  or e. The narrow peaks near 0.6 and 0.8 GeV correspond to the  $\psi(2S)$  and X(3872) states, respectively. Reproduced from *Phys. Rev. Lett.* **91** (2003) 262001.

## BPHY-2015-03



We can study this at ATLAS and CMS (and of course at LHCb...)

# X(3872) in details

#### 2005.13419

# Compare with $m(D^0\overline{D}^{0^*}) = 3871.70 \pm 0.11 \text{ MeV}$

 $m_{\chi_{c1}(3872)} = 3871.695 \pm 0.067 \pm 0.068 \pm 0.010 \, {\rm MeV}$ 

#### Quite intriguing!

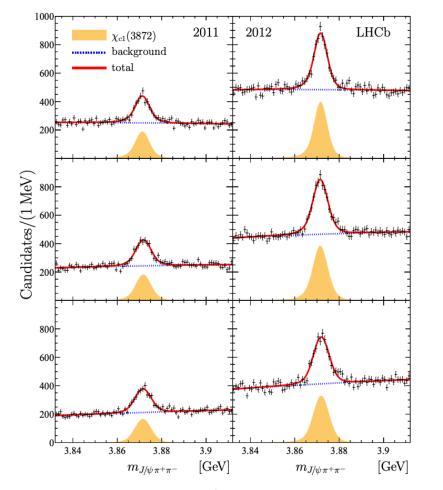


Figure 2: Mass distributions for  $J/\psi \pi^+\pi^-$  candidates in the  $\chi_{c1}(3872)$  region for (top) the low, (middle) mid and (bottom) high  $p_{\pi^+\pi^-}$  bins. The left (right)-hand plot is for 2011 (2012) data. The projection of the fit described in the text is superimposed.

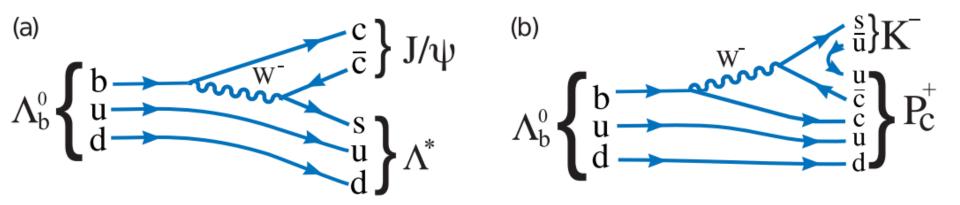
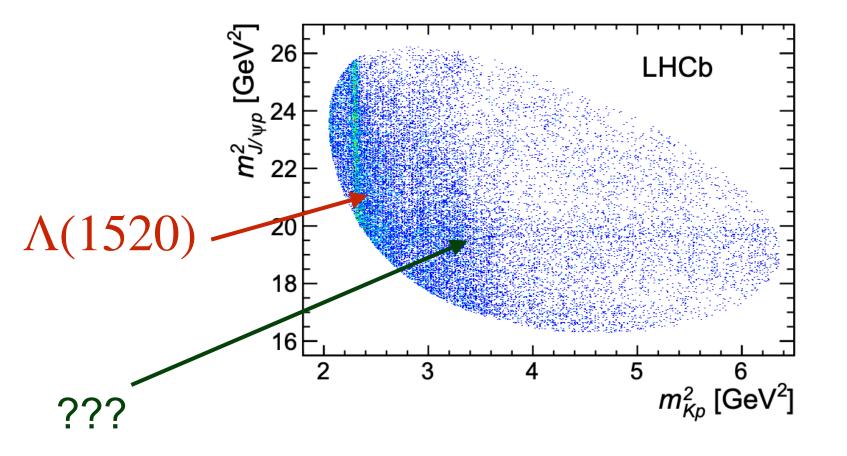


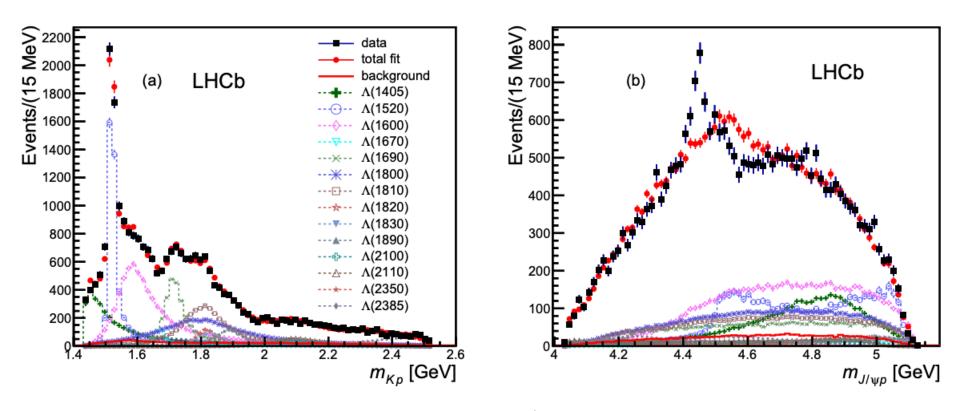
Figure 1: Feynman diagrams for (a)  $\Lambda_b^0 \to J/\psi \Lambda^*$  and (b)  $\Lambda_b^0 \to P_c^+ K^-$  decay.

LHCb studies  $\Lambda_b$  decays ( $\Lambda^* \to K^- p$ ) and found an interesting feature in the  $J/\psi$  + p mass distribution ( $J/\psi \to \mu^+ \mu^-$ )



Very useful for finding structure and understanding decay chains!

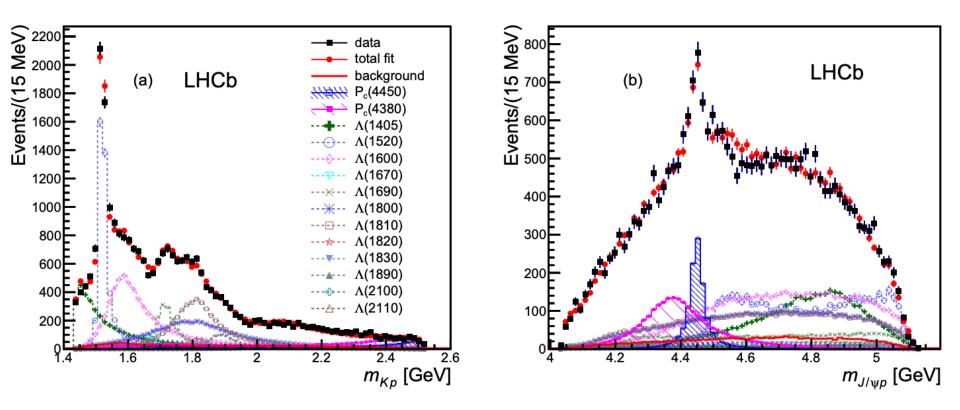
# Mass plots



Need to account for large numbers of  $\Lambda$  states and perform a full analysis accounting for QM interference and angular coefficients (up to J=9/2!). Tricky, as objects near kinematic limits combined with angular effects (and potentially interference) can cause resonant-like peaks that are not from real resonances, just **kinematic reflections**. But clearly the  $J/\psi$  + p distribution is not well modeled

# More exotic hadrons

#### 1507.03414

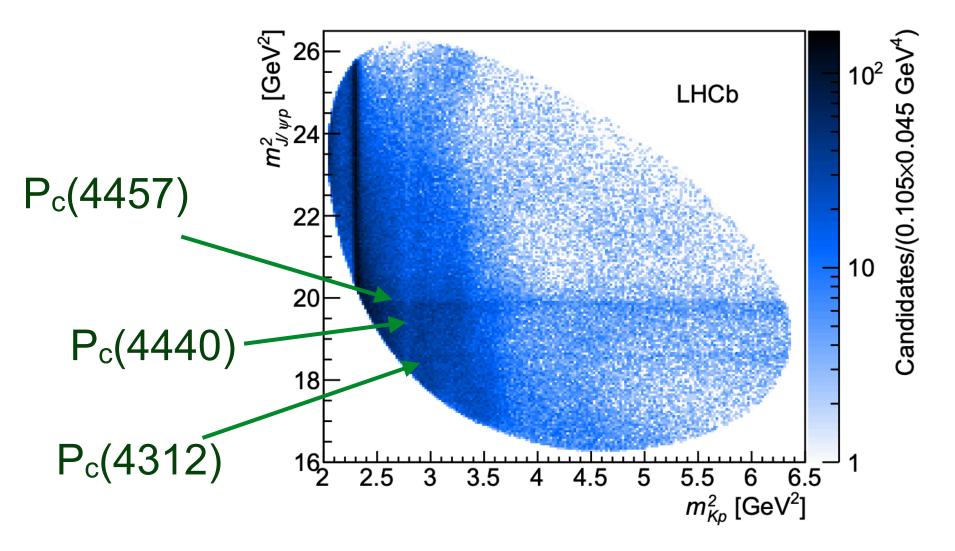


Including two new states gives a much better fit! Two P<sub>c</sub> states: Pc(4380),  $J^{P} = (3/2)^{-}$  and P<sub>c</sub>(4450),  $J^{P} = (5/2)^{+}$ 

Different binding mechanisms of pentaquark states are possible. Tight-binding was envisioned originally [3,4,35]. A possible explanation is heavy-light diquarks [36]. Examples of other mechanisms include a diquark-diquark-antiquark model [37,38], a diquark-triquark model [39], and a coupled channel model [40]. Weakly bound "molecules" of a baryon plus a meson have been also discussed [41].

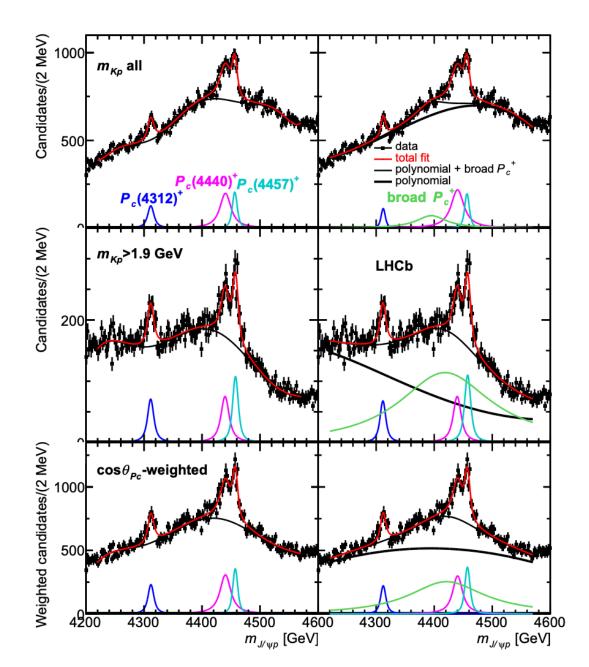
## More recent updates

#### 1904.03947



#### More recent updates

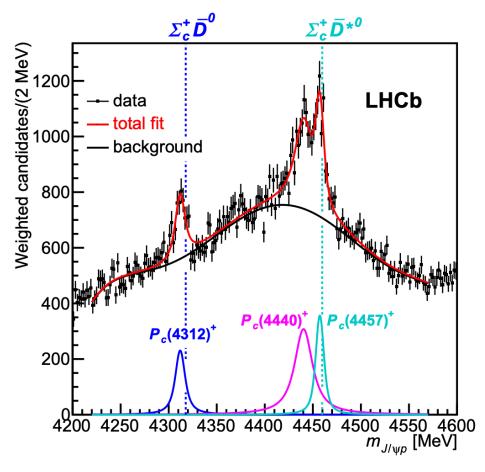
1904.03947

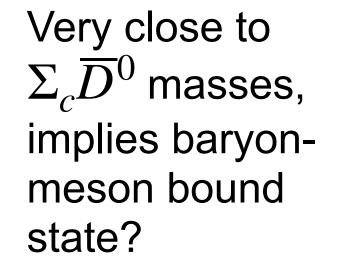


Fit in separate bins of  $\cos \theta_{PC}$ (cos of angle between K- and  $J/\psi$ ) in P<sub>C</sub> rest frame

### More recent updates

1904.03947





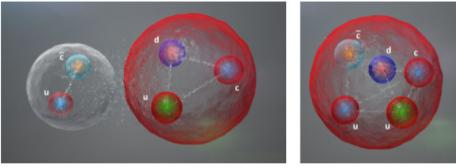
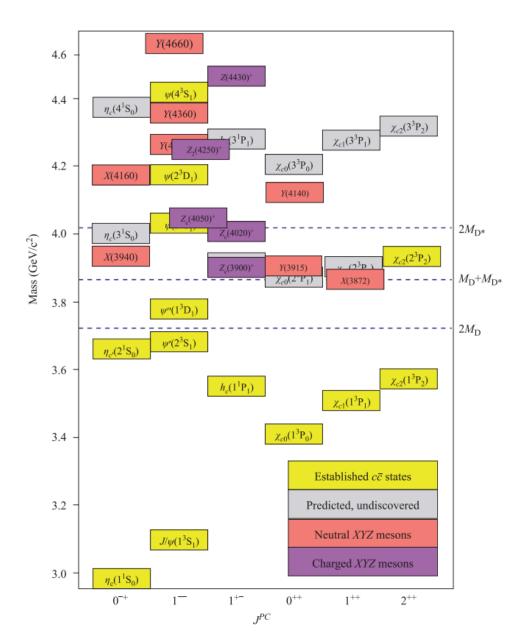


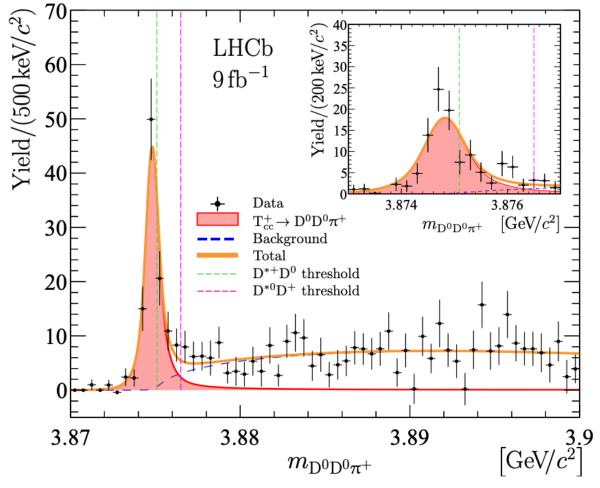
Figure 3. Possible quark combinations to make mesons, baryons and pentaquarks. Copyright CERN.

# 1411.7738



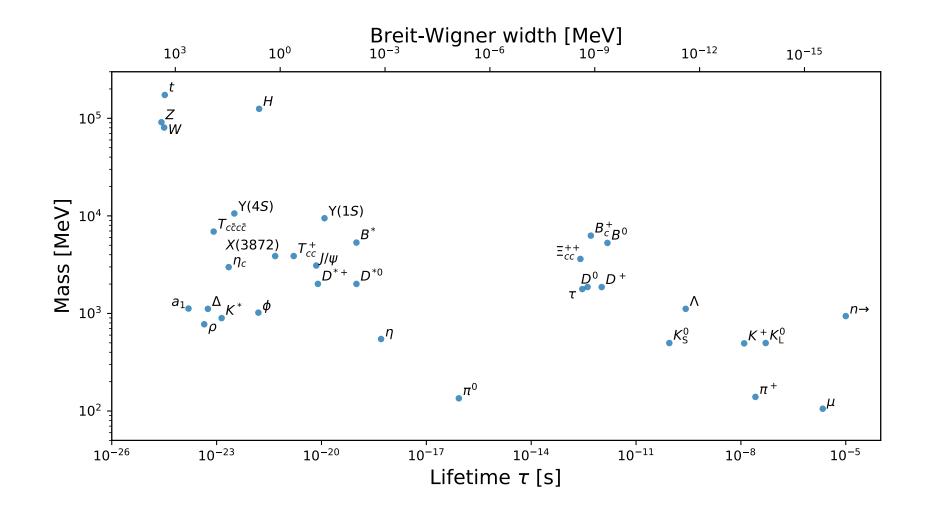
Not always obvious how the new states fit together. Molecular picture requires something to bind the pieces of the molecule. Pion carries isospin, so it cannot bind hadrons with I == 0. And it is a pseudoscalar, so it cannot bind two other pseudoscalars and conserve parity. Perhaps rho contributes

#### 2109.01038



Very close to threshold, so (relatively) longlived, again supporting the molecular picture. Found using (for both)  $D^0 \rightarrow K^- \pi^+$ 

Figure 1: The  $D^0D^0\pi^+$  mass distribution. The  $D^0D^0\pi^+$  mass distribution where the contribution of the non- $D^0$  background has been statistically subtracted. The result of the fit described in the text is overlaid.



# Let's talk about this!

I'd like for you to begin to think about your final project. You should put in some effort to prepare it, so time to start now. You need to pick a single **analysis** or result to present (something public in the past ~year)

http://arxiv.org/archive/hep-ex

Minos ATLAS CMS LHCb **ALICE SNO** Belle-2 AMS LIGO VIRGO LUX Dark Energy Survey **Pierre Auger** X17 Fermi Gamma Ray Telescope

Pick an analysis/paper and email it to me for approval no later than 1 week from today. I'll ask for some progress reports in the future, but for now I just want you to pick something interesting (and not what you work on for research, and not what someone else will be talking about). Talk to me if you need help picking a topic!

Want a 25 minute presentation on the topic! You should be including theory background if possible, as well as information on the detector, the analysis, the background estimation, and the significance of the result. We will all discuss the presentation for 10 minutes after you're done (aka ask you some questions)

I also reserve the right to reject papers that are too broad, too narrow, too old, or too out of focus for this course

- Note 0: It's not first-come/first-served, but instead we will flip coins or play rock-paper-scissors for who gets which topic.
- Note 1: If you dropped by my office to discuss the paper, that does not count as fulfilling your homework assignment. Please sent it to me by email
- Note 2: This really counts as a homework assignment. So don't miss the deadline! If you do, you get points off and I get to pick a topic for you :)

Note 3: Only one topic per person, so you might want to have 1-2 backups in mind (class is small so maybe that's OK). I'll let you know class after it's due where we have duplicates (and alternates due then class after that)

Note 4: You are not covering an "experiment", but rather a single analysis/limit/measurement. So I want you to have the paper you will be reading in mind. If there is no physics result, then this does not count