# Just a brief review at this stage before we delve in

#### **Relativistic kinematics**



#### **Inverse transformation**





Cannot define simultaneity in all reference frames

Lorentz Contraction (objects shrink)

Time dilation (moving clocks run slow)

What happens if u' = speed of light (which is ... what in our units?)

 $t'_{A} = t'_{B} + \gamma v \left( x_{B} - x_{A} \right) \checkmark$ 

 $L = \frac{L'}{\gamma}$ 

 $t = \gamma t$ 

 $u = \frac{u' + v}{1 + u'v}$ 

Velocity addition is not as simple as before

# Are folks familiar with fourvector notation?

## Get used to them :)



All we've done here is label the 0th component of space-time as time, and then the normal (x,y,z)

Note that x in the 4-vector is standard notation, but it is **not** the position x

$$x^{0'} = \gamma \left( x^0 - vx^1 \right)$$
$$x^{1'} = \gamma \left( x^1 - vx^0 \right)$$
$$x^{2'} = x^2$$
$$x^{3'} = x^3$$

Introduces a nice symmetry to the transformations, which leads us to re-writing it as...

$$x^{\mu'} = \sum_{\nu=0}^{3} \Lambda^{\mu}_{\nu} x^{\nu}$$

Where  $\Lambda$  is a 4x4 matrix. Let's get the coordinates of the matrix on the board...

$$\Lambda = \begin{bmatrix} \gamma & -\gamma v & 0 & 0 \\ -\gamma v & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

The symmetry really stands out here

$$x^{\mu'} = \sum_{\nu=0}^{3} \Lambda^{\mu}_{\nu} x^{\nu} = \Lambda^{\mu}_{\nu} x^{\nu}$$

Summation notation will be used a LOT The second form has the sum implicit, and is shorthand to save smart (read: lazy) physicists from typing and writing

# Define this invariant I. How does it transform? What is I'?

$$I = (x^{0})^{2} - (x^{1})^{2} - (x^{2})^{2} - (x^{3})^{2}$$

$$I' = (x^{0'})^{2} - (x^{1'})^{2} - (x^{2'})^{2} - (x^{3'})^{2}$$

$$I' = (x^{0'})^{2} - (x^{1'})^{2} - (x^{2})^{2} - (x^{3})^{2}$$

$$I' = \gamma^{2} (x^{0} - vx^{1})^{2} - \gamma^{2} (x^{1} - vx^{0})^{2} - (x^{2})^{2} - (x^{3})^{2}$$

$$I' = \gamma^{2} ((x^{0})^{2} - (x^{1})^{2} + v^{2} (x^{1})^{2} - v^{2} (x^{0})^{2} - 2vx^{0}x^{1} + 2vx^{0}x^{1}) - (x^{2})^{2} - (x^{3})^{2}$$

$$I' = \gamma^{2} ((x^{0})^{2} (1 - v^{2}) - (x^{1})^{2} (1 - v^{2})) - (x^{2})^{2} - (x^{3})^{2}$$

Now use 
$$\gamma^2 = \frac{1}{1 - v^2}$$

$$I = (x^{0})^{2} - (x^{1})^{2} - (x^{2})^{2} - (x^{3})^{2}$$
$$I' = \gamma^{2} \left( (x^{0})^{2} \left( 1 - v^{2} \right) - (x^{1})^{2} \left( 1 - v^{2} \right) \right) - (x^{2})^{2} - (x^{3})^{2} = I$$

# Now we see why we called I an invariant.





 $x_0 = x^0$   $x_1 = -x^1$   $x_2 = -x^2$  $x_3 = -x^3$  The names are not quite as important as remembering that they are different (and that it's easy to convert from one to another)

$$a^{\mu}b_{\mu} = a_{\mu}b^{\mu} = a^{0}b^{0} - a^{1}b^{1} - a^{2}b^{2} - a^{3}b^{3}$$

Four-vector version of the dot product (which is also an invariant but only for three spatial dimensions)

$$a^{\mu}b_{\mu} = a_{\mu}b^{\mu} = a^{0}b^{0} - a^{1}b^{1} - a^{2}b^{2} - a^{3}b^{3}$$

Four-vector version of the dot product (which is also an invariant but only for three spatial dimensions)

$$a^{\mu}b_{\mu} = a^0b^0 - \vec{a}\cdot\vec{b}$$

$$a^2 = a^{\mu}a_{\mu} = a^0a^0 - \vec{a}^2$$

$$a^2 = a^{\mu}a_{\mu} = a^0a^0 - \vec{a}^2$$

- $a^2 > 0$  time-like
- $a^2 = 0$  light-like
- a<sup>2</sup> < 0 space-like (like normal spatial 3-vectors)



#### Conservation of momentum: $m_A v_A + m_B v_B = m_C v_C + m_D v_D$

Now imagine a frame moving with constant velocity V with respect to this frame  $v'_{A} = v_{A} - V$  $v'_{B} = v_{B} - V$  $v'_{C} = v_{C} - V$  $v'_{D} = v_{D} - V$ 

Now imagine a frame moving with constant velocity V with respect to this frame

 $m_A v'_A + m_B v'_B = m_C v'_C + m_D v'_D$ 

$$v'_{A} = v_{A} - V$$
$$v'_{B} = v_{B} - V$$
$$v'_{C} = v_{C} - V$$
$$v'_{D} = v_{D} - V$$

 $m_A v_A + m_B v_B - V(m_A + m_B) = m_C v_C + m_D v_D - V(m_C + m_D)$ And we know  $m_A v_A + m_B v_B = m_C v_C + m_D v_D$ 

$$V(m_A + m_B) = V(m_C + m_D)$$
$$m_A + m_B = m_C + m_D$$

Classical conservation of momentum (in all inertial reference frames) implies conservation of mass!

 $m_A + m_B = m_C + m_D$ 



# Conservation of momentum:

 $m_A v_A + m_B v_B = m_C v_C + m_D v_D$ 

Now imagine a frame moving with constant velocity V with respect to this frame. Note the use of new velocity addition formula!  $v'_{A} = \frac{v_{A} - V}{1 - v_{A}V}$  $v'_{B} = \frac{v_{B} - V}{1 - v_{B}V}$  $v'_{C} = \frac{v_{C} - V}{1 - v_{C}V}$  $v'_{D} = \frac{v_{D} - V}{1 - v_{D}V}$ 

#### Conservation of momentum in SR

$$m_{A}v_{A} + m_{B}v_{B} = m_{C}v_{C} + m_{D}v_{D}$$
$$m_{A}v_{A}' + m_{B}v_{B}' = m_{C}v_{C}' + m_{D}v_{D}'$$

$$m_A \frac{v_A - V}{1 - v_A V} + m_B \frac{v_B - V}{1 - v_B V} = m_C \frac{v_C - V}{1 - v_C V} + m_D \frac{v_D - V}{1 - v_D V} \qquad v'_C = \frac{v_C - V}{1 - v_C V}$$
$$v'_D = \frac{v_D - V}{1 - v_D V}$$

Velocities no longer nicely cancel, so we're stuck in the relativistic case! No longer have proof of mass conservation  $v'_A = \frac{v_A - V}{1 - v_A V}$ 

 $v'_B = \frac{v_B - V}{1 - v_B V}$ 

Nominal definition of momentum  $p = m \frac{dx}{dt}$ 

Not surprising that we run into trouble - both numerator and denominator are not Lorentz invariant! Let's make an ad hoc try and replace time with proper time

$$\tau = t/\gamma$$

$$p = m \frac{dx}{d\tau} = m \frac{dx}{d(t/\gamma)} = \gamma m \frac{dx}{dt}$$

# First frame

$$\frac{m_A v_A}{\sqrt{1 - v_A^2}} + \frac{m_B v_B}{\sqrt{1 - v_B^2}} = \frac{m_C v_C}{\sqrt{1 - v_C^2}} + \frac{m_D v_D}{\sqrt{1 - v_D^2}}$$

### New frame

$$\frac{m_A v'_A}{\sqrt{1 - (v'_A)^2}} + \frac{m_B v'_B}{\sqrt{1 - (v'_B)^2}} = \frac{m_C v'_C}{\sqrt{1 - (v'_C)^2}} + \frac{m_D v'_D}{\sqrt{1 - (v'_D)^2}}$$
$$\frac{v'_A = \frac{v_A - V}{1 - v_A V}}{v'_B = \frac{v_B - V}{1 - v_B V}}$$
$$v'_C = \frac{v_C - V}{1 - v_C V}$$
$$v'_D = \frac{v_D - V}{1 - v_D V}$$

#### Some algebra in new frame

$$\frac{m_A v'_A}{\sqrt{1 - (v'_A)^2}} + \frac{m_B v'_B}{\sqrt{1 - (v'_B)^2}} = \frac{m_C v'_C}{\sqrt{1 - (v'_C)^2}} + \frac{m_D v'_D}{\sqrt{1 - (v'_D)^2}}$$

Let's look at first term:

 $\frac{m_A \left(\frac{v_A - V}{1 - v_A V}\right)}{\sqrt{1 - \left(\frac{v_A - V}{1 - v_A V}\right)^2}} =$  $\frac{m_A \left(\frac{v_A - V}{1 - v_A V}\right)}{\sqrt{\frac{(1 - v_A V)^2 - (v_A - V)^2}{(1 - v_A V)^2}}} =$  $\frac{m_A(v_A - V)}{\sqrt{(1 - v_A V)^2 - (v_A - V)^2}} =$  $\frac{m_A(v_A - V)}{\sqrt{1 + v_A^2 V^2 - 2v_A V - v_A^2 - V^2 + 2v_A V}} =$  $\frac{m_A(v_A - V)}{\sqrt{1 + v_A^2 V^2 - v_A^2 - V^2}}$ 

$$\frac{m_D v_D}{\sqrt{1 - (v'_D)^2}}$$

$$\frac{v_A - V}{1 - v_A V}$$

$$\frac{v_B}{v_B} = \frac{v_B - V}{1 - v_B V}$$

$$\frac{v_C}{v_C} = \frac{v_C - V}{1 - v_C V}$$

$$\frac{v_D}{v_D} = \frac{v_D - V}{1 - v_D V}$$

## Some algebra in new frame

$$\begin{aligned} \frac{m_A v'_A}{\sqrt{1 - (v'_A)^2}} + \frac{m_B v'_B}{\sqrt{1 - (v'_B)^2}} &= \frac{m_C v'_C}{\sqrt{1 - (v'_C)^2}} + \frac{m_D v'_D}{\sqrt{1 - (v'_D)^2}} \\ \text{Let's look at first term:} \\ \frac{m_A \left(\frac{v_A - V}{1 - v_A V}\right)}{\sqrt{1 - \left(\frac{v_A - V}{1 - v_A V}\right)^2}} &= \\ \frac{m_A (v_A - V)}{\sqrt{1 + v_A^2 V^2 - v_A^2 - V^2}} \\ \frac{m_A v_A}{\sqrt{(1 - v_A^2)(1 - V^2)}} - \frac{m_A V}{\sqrt{(1 - v_A^2)(1 - V^2)}} &= \\ \frac{m_A v_A}{\sqrt{1 - v_A^2}} \frac{1}{\sqrt{1 - V^2}} - \frac{m_A}{\sqrt{1 - v_A^2}} \frac{V}{\sqrt{1 - V^2}} = \end{aligned}$$

#### Some algebra in new frame



## So...

 $\begin{aligned} \frac{m_A v_A}{\sqrt{1 - v_A^2}} \frac{1}{\sqrt{1 - V^2}} &- \frac{m_A}{\sqrt{1 - v_A^2}} \frac{V}{\sqrt{1 - V^2}} + \\ \frac{m_B v_B}{\sqrt{1 - v_B^2}} \frac{1}{\sqrt{1 - V^2}} &- \frac{m_B}{\sqrt{1 - v_B^2}} \frac{V}{\sqrt{1 - V^2}} = \\ \frac{m_C v_C}{\sqrt{1 - v_C^2}} \frac{1}{\sqrt{1 - V^2}} &- \frac{m_C}{\sqrt{1 - v_C^2}} \frac{V}{\sqrt{1 - V^2}} + \\ \frac{m_D v_D}{\sqrt{1 - v_D^2}} \frac{1}{\sqrt{1 - V^2}} &- \frac{m_D}{\sqrt{1 - v_D^2}} \frac{V}{\sqrt{1 - V^2}} \end{aligned}$ 

# But from original frame

$$\frac{m_A v_A}{\sqrt{1 - v_A^2}} + \frac{m_B v_B}{\sqrt{1 - v_B^2}} = \frac{m_C v_C}{\frac{m_C v_C}{\sqrt{1 - v_C^2}}} + \frac{m_D v_D}{\sqrt{1 - v_D^2}}$$

Thus we get



#### Cancel the common term...

$$\frac{m_A}{\sqrt{1 - v_A^2}} + \frac{m_B}{\sqrt{1 - v_B^2}} = \frac{m_C}{m_C} + \frac{m_D}{\sqrt{1 - v_D^2}}$$

This quantity must be conserved then! But what is it?

$$\frac{m}{\sqrt{1-v^2}} = m\left(1-v^2\right)^{-1/2} \sim m(1+v^2/2)$$

So this is a statement of the conservation of energy! At small velocities, energy is equal to the rest mass plus the kinematic energy (note that when v = 0 the energy is not zero, but m!)

$$\frac{m}{\sqrt{1-v^2}} = m\left(1-v^2\right)^{-1/2} \sim m(1+v^2/2)$$

So this is a statement of the conservation of energy! At small velocities, energy is equal to the rest mass plus the kinematic energy

$$p^{\mu} = (E, p_x, p_y, p_z)$$
$$p^{\mu}p_{\mu} = E^2 - \mathbf{p}^2 = m^2$$

This way, p<sup>2</sup> is invariant, as it must be

For m = 0 (ex for: photons), |p|=E

$$p^{\mu} = (E, p_x, p_y, p_z)$$
$$p^{\mu}p_{\mu} = E^2 - \mathbf{p}^2 = m^2$$



As expected, for light, |v| = 1 and |p| = E

**Relativistic collisions** 

We just saw that energy and momentum (correctly defined) are conserved. For process A+B→C+D:

$$E_A + E_B = E_C + E_D$$
$$\vec{p}_A + \vec{p}_B = \vec{p}_C + \vec{p}_D$$

Conveniently written as:

$$p^\mu_A + p^\mu_B = p^\mu_C + p^\mu_D$$

Mass is clearly not conserved:

$$W^+ \to e^+ + \nu_e$$

What about a particle decay?  $A \rightarrow B + C$  $A_{\mu}A^{\mu} = m_A^2 \rightarrow (B + C)_{\mu}(B + C)^{\mu} = (E_B + E_C)^2 - (\mathbf{p_B} + \mathbf{p_C})^2$ 

Lets us estimate the mass of a parent particle from its decay products. Conversation of energy also tells us:

$$m_A = \sqrt{m_B^2 + \mathbf{p}_B^2} + \sqrt{m_C^2 + \mathbf{p}_C^2}$$

Particle decay products cannot be more massive than mass of original particle! (Remember - proton stability...)



Particle with total energy = 2m collides with an identical particle at rest. They stick together. What is the mass of the resulting lump? What is its velocity?

Conservation of E: 2m+m = 3m = E



$$E_1^2 - \mathbf{p}^2 = m^2$$
$$(2m^2) - m^2 = \mathbf{p}^2$$
$$\mathbf{p}^2 = 3m^2$$
$$|\mathbf{p}| = \sqrt{3}m$$

$$|\mathbf{p}| = |\mathbf{p}'| = \sqrt{3}m$$



$$E_1^2 - \mathbf{p}^2 = m^2$$
$$(2m)^2 - m^2 = \mathbf{p}^2$$
$$\mathbf{p}^2 = 3m^2$$
$$|\mathbf{p}| = \sqrt{3}m$$

$$|\mathbf{p}| = |\mathbf{p}'| = \sqrt{3}m$$



$$M^{2} = E^{2} - (\mathbf{p}')^{2} = 9m^{2} - 3m^{2}$$
$$M = \sqrt{6}m$$
Total mass not conserved

$$|v| = \frac{|\mathbf{p}|}{E} = \frac{\sqrt{3m}}{3m} = \frac{1}{\sqrt{3}}$$

Let's follow Griffiths Ex 3.3

# A pion at rest decays into a muon and a neutrino. What is the speed of the muon?



Let's follow Griffiths Ex 3.3

Before

Conservation of momentum:  $\mathbf{p}_{\mu} = -\mathbf{p}_{\nu}$ Conservation of energy:  $\mathbf{E}_{\mu} + \mathbf{E}_{\nu} = \mathbf{E}_{\pi}$ 

Π

So muon and neutrino travel in opposite directions ("back to back")



## Apply conservation of energy:

$$E_{\mu} + E_{\nu} = E_{\pi}$$

$$\sqrt{m_{\mu}^2 + \vec{p}_{\mu}^2} + \sqrt{m_{\nu}^2 + \vec{p}_{\nu}^2} = \sqrt{m_{\pi}^2 + \vec{p}_{\pi}^2}$$

$$\sqrt{m_{\mu}^2 + \vec{p}_{\mu}^2} + |\vec{p}_{\nu}| = m_{\pi}$$
Pion was at rest
Neutrino mass ~0

# Apply conservation of energy:

#### Let's follow Griffiths Ex 3.3

$$\begin{aligned} |\vec{p}_{\mu}| &= \frac{m_{\pi}^2 - m_{\mu}^2}{2m_{\pi}} \\ E_{\mu}^2 &= \vec{p}_{\mu}^2 + m_{\mu}^2 = \frac{\left(m_{\pi}^2 - m_{\mu}^2\right)^2}{4m_{\pi}^2} + m_{\mu}^2 \qquad |\mathbf{v}| = |\mathbf{p}|/\mathbf{E} \\ E_{\mu}^2 &= \frac{m_{\pi}^4 + m_{\mu}^4 - 2m_{\pi}^2 m_{\mu}^2 + 4m_{\pi}^2 m_{\mu}^2}{4m_{\pi}^2} \\ E_{\mu}^2 &= \frac{m_{\pi}^4 + m_{\mu}^4 + 2m_{\pi}^2 m_{\mu}^2}{4m_{\pi}^2} \qquad |\mathbf{v}_{\mu}| = \frac{m_{\pi}^2 - m_{\mu}^2}{m_{\pi}^2 + m_{\mu}^2} \\ E_{\mu}^2 &= \frac{\left(m_{\pi}^2 + m_{\mu}^2\right)^2}{4m_{\pi}^2} \\ E_{\mu} &= \frac{m_{\pi}^2 + m_{\mu}^2}{2m_{\pi}} \end{aligned}$$

## Pion mass ~135 MeV Muon mass ~106 MeV

$$|v_{\mu}| = \frac{m_{\pi}^2 - m_{\mu}^2}{m_{\pi}^2 + m_{\mu}^2}$$

# v~0.25 (25% of the speed of light)

# Let's try using four-vector notation (remember that p<sup>2</sup>=m<sup>2</sup> but be careful that p here is the four-vector p)

$$p_{\pi} = p_{\mu} + p_{\nu}$$

$$p_{\nu} = p_{\pi} - p_{\mu}$$

$$p_{\nu}^{2} = (p_{\pi} - p_{\mu})^{2} = m_{\nu} = 0$$

$$0 = p_{\pi}^{2} + p_{\mu}^{2} - 2p_{\pi}p_{\mu}$$

$$0 = m_{\pi}^{2} + m_{\mu}^{2} - 2p_{\pi}p_{\mu}$$

What is  $p_{\pi}p_{\mu}$ ? The pion momentum is zero, so only first component (energy) contributes, and so  $p_{\pi}p_{\mu} = E_{\pi}E_{\mu}$ 

$$0 = m_{\pi}^{2} + m_{\mu}^{2} - 2m_{\pi}E_{\mu}$$
$$E_{\mu} = \frac{m_{\pi}^{2} + m_{\mu}^{2}}{2m_{\pi}}$$

Also, since pion was at rest,  $E_{\pi} = m_{\pi}$  so  $p_{\pi}p_{\mu} = m_{\pi}E_{\mu}$ 

#### Let's follow Griffiths Ex 3.3 (continued)

$$p_{\pi} = p_{\mu} + p_{\nu}$$

$$p_{\mu} = p_{\pi} - p_{\nu}$$

$$p_{\mu}^{2} = (p_{\pi} - p_{\nu})^{2} = m_{\mu}^{2}$$

$$p_{\mu}^{2} = p_{\pi}^{2} + p_{\nu}^{2} - 2p_{\pi}p_{\nu}$$

$$m_{\mu}^{2} = m_{\pi}^{2} - 2m_{\pi}|p_{\mu}|$$
$$|p_{\mu}| = \frac{m_{\pi}^{2} - m_{\mu}^{2}}{2m_{\pi}}$$

What is  $p_{\pi}p_{\nu}$ ? The pion momentum is zero, so only first component (energy) contributes, and so  $p_{\pi}p_{\nu} = E_{\pi}E_{\nu}$ Since pion was at rest,  $E_{\pi} = m_{\pi}$ . What is  $E_{\nu}$ ? It is the magnitude of the neutrino momentum, but that is the same as the magnitude of the muon momentum

# pp→ppp<del>p</del>

What is the threshold energy for this process?

At threshold, all four final state objects are at rest (if they were not at rest we could reduce the CoM collision energy and reduce their momentum) Invariant:

Initial 
$$(p_1 + p_2)^2 = (E + m)^2 - p^2$$
  
 $p_1=(E,p)$   $p_2=(m,0)$ 



$$16m^{2} = (E+m)^{2} - \mathbf{p}^{2} \qquad \mathbf{p}^{2} = E^{2} - m^{2}$$
$$16m^{2} = E^{2} + 2mE + m^{2} - E^{2} + m^{2}$$
$$14m^{2} = 2mE \rightarrow E = 7m$$

# What if we didn't fire at a stationary target?



# In this example, have to give protons a lot more kinetic energy for the fixed target version! Of course, the accelerator for the beams is much more complex :)

#### GZK cutoff

Greisen–Zatsepin–Kuzmin: Over long enough distances, high-energy cosmic rays coming from the Universe should interact with the cosmic microwave background

<u>http://</u> <u>www2.astro.psu.ed</u> <u>u/users/nnp/cr.html</u>



$$\gamma + p \to p\pi^{0}$$

$$p_{\gamma} + p_{p} = p_{p'} + p_{\pi^{0}}$$

$$(p_{\gamma} + p_{p})^{2} = (p_{p'} + p_{\pi^{0}})^{2}$$

$$p_{\gamma}^{2} + p_{p}^{2} + 2p_{\gamma} \cdot p_{p} = p_{p'}^{2} + p_{\pi^{0}}^{2} + 2p_{\pi^{0}} \cdot p_{p'}$$

$$m_{p}^{2} + 2p_{\gamma} \cdot p_{p} = m_{p}^{2} + m_{\pi^{0}}^{2} + 2p_{\pi^{0}} \cdot p_{p'}$$

$$2p_{\gamma} \cdot p_{p} = m_{\pi^{0}}^{2} + 2p_{\pi^{0}} \cdot p_{p'}$$

At threshold, pion and proton in final state are at rest and their momenta are zero (and energies equal to their masses)

$$2p_{\gamma} \cdot p_p = m_{\pi^0}^2 + 2m_{\pi^0}m_p$$

At threshold, have a head-on collision. Assume highly relativistic particles and we get...

$$2p_{\gamma} \cdot p_p = m_{\pi^0}^2 + 2m_{\pi^0}m_p$$

## Head-on collision with energies >> m<sub>p</sub>

$$p_{\gamma} = (E_{\gamma}, E_{\gamma}, 0, 0)$$
$$p_{p} = (E_{p}, -E_{p}, 0, 0)$$
$$2p_{p} \cdot p_{\gamma} = 2E_{p}E_{\gamma} - 2(-E_{p})E_{\gamma} = 4E_{p}E_{\gamma}$$

Want to solve for unknown cosmic rate proton energy. Photon energy is CMB (6x10<sup>-4</sup> ev). Plug in and we get ~2x10<sup>20</sup> ev

Similarly have  $\gamma + p \rightarrow n\pi^+$ 

#### GZK cutoff

Of course, we used an "average" CMB photon, and ignored a full calculation of the kinematics (which goes through delta resonances) and assumed protons. Nevertheless, distribution doesn't go to zero! Where are these sources coming from? Nearby sources? Or new physics?

http:// www2.astro.psu.ed u/users/nnp/cr.html



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Griffiths 3.3, 3.4, 3.6, 3.15, 3.16, 3.17, 3.25 (good for you to start getting familiar with the PDG, I suspect)

Also: At the HERA collider, 27.5 GeV electrons collided head-on with 920 GeV protons. What was the center of mass energy?