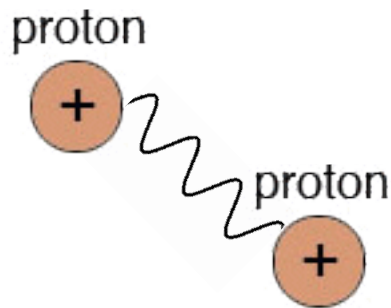


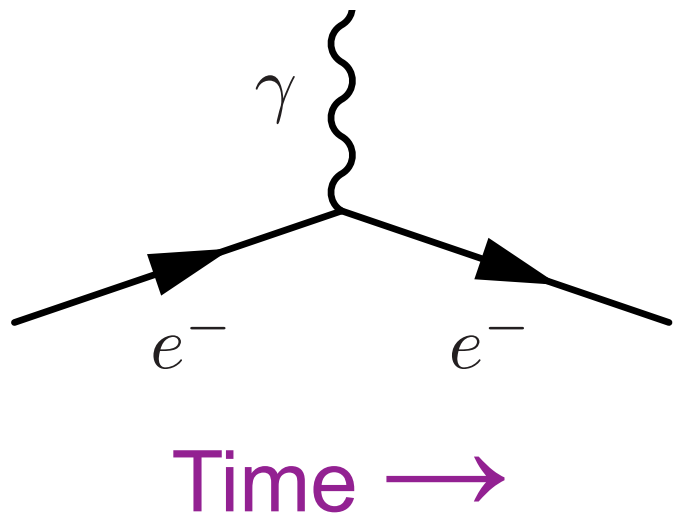
Photon as a particle (classical). But of course, what we really mean is that the field is quantized (here providing an attractive force), and the quantized unit of the field transmits some momentum from one object to another

But these forces can also be repulsive (more intuitive, actually)



Photon are still exchanged here, but now they transmit a repulsive force. It's a bit counterintuitive to think of forces being transmitted (or even better, '**mediated**') by particles, especially on a macroscopic scale, but we are anyway talking about single particles

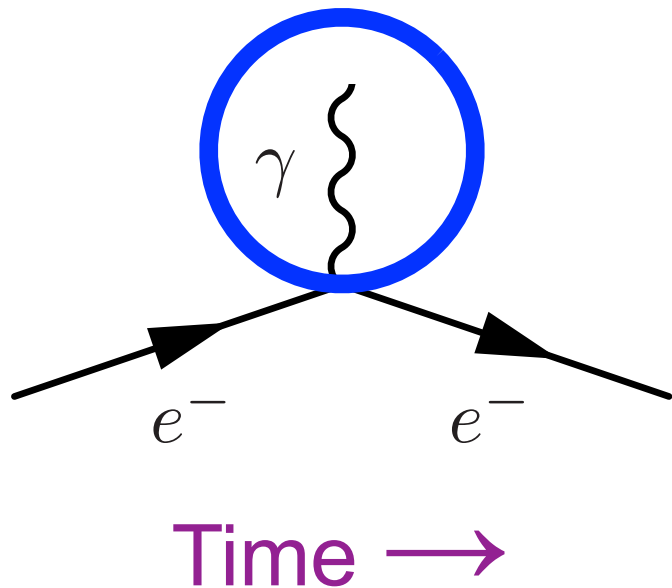
Feynman diagram



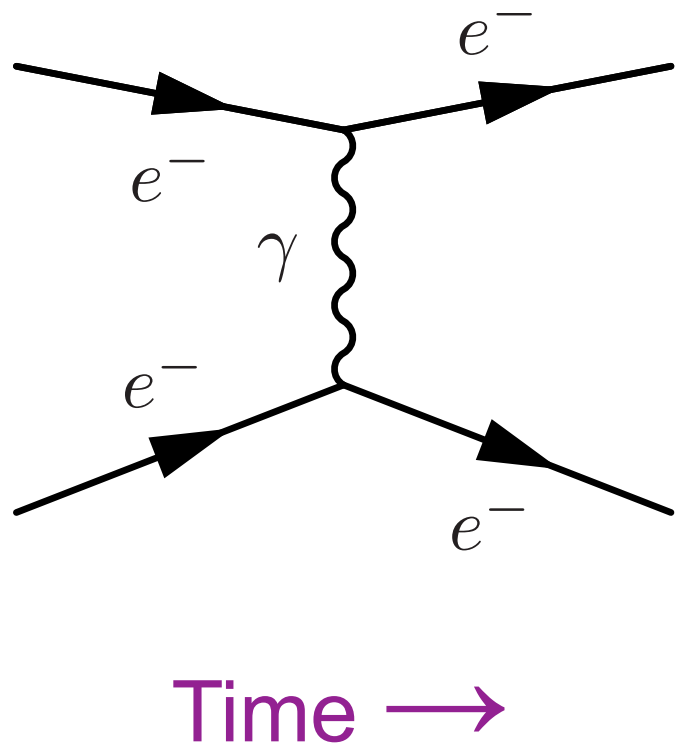
Think of time “passing” to the right, so initially, we start with an electron. And in the final state, we end with an electron, and at some intermediate time, there is a photon interacting with this electron (whether it emits it or absorbs it). We can combine multiple such diagrams

Feynman diagram

Note: Photon gets a squiggly line!

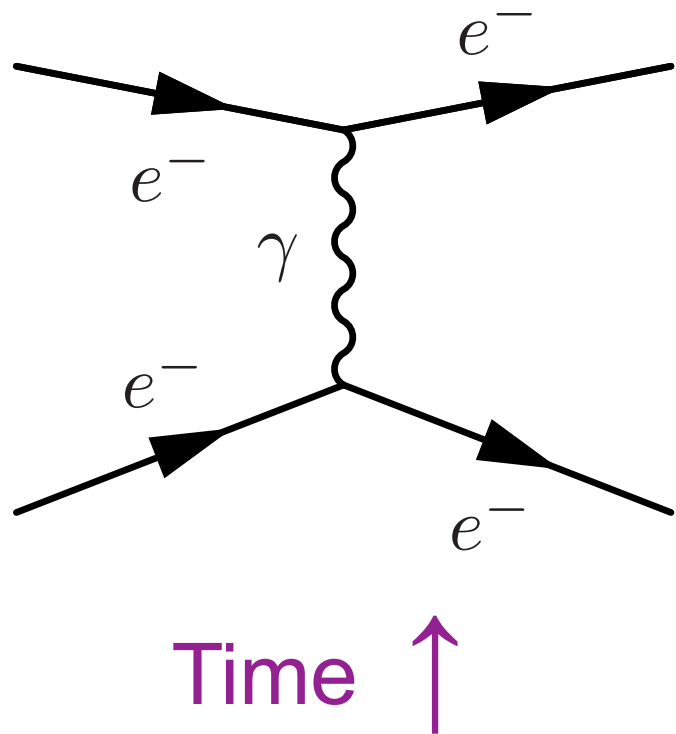


Think of time “passing” to the right, so initially, we start with an electron. And in the final state, we end with an electron, and at some intermediate time, there is a photon interacting with this electron (whether it emits it or absorbs it). We can combine multiple such diagrams



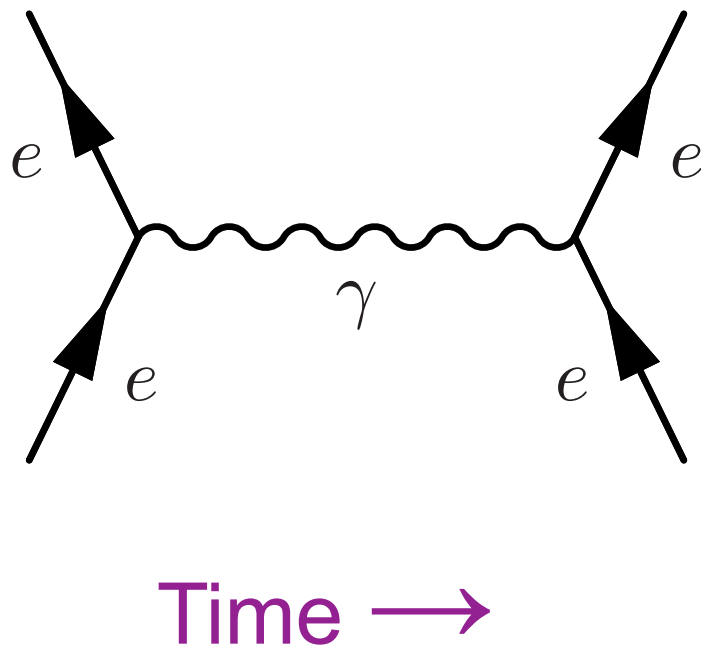
At some initial time, there were two separate electrons. At a later time, there were also two electrons. At an intermediate time, they were exchanging a photon between them (repulsive force)

Møller scattering



What happens if we just change the direction of time? Everyone turn your heads now :)

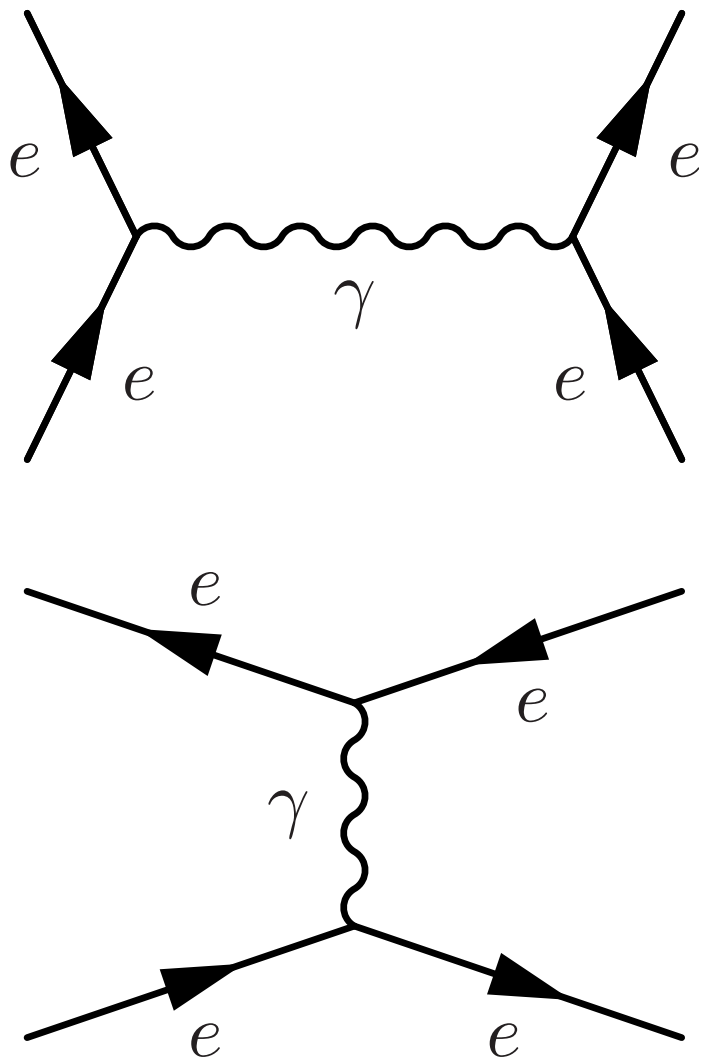
Does this remind you of anything from earlier in the course?



Electron vs positron
should now be more
obvious/automatic

Here I put things back such that time moves to the right. Note that in the initial time (on left) there is a fermion line moving backwards! This is an antiparticle moving forwards in time. And on the right there is a fermion line also going backwards, this is again an anti-particle

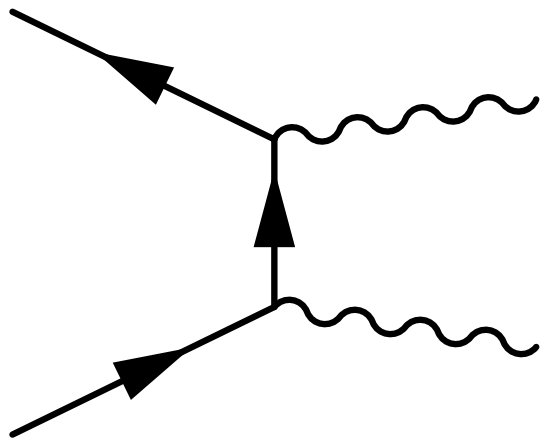
Bhabha scattering



Note the second diagram below... start with e^+e^- and end with e^+e^- ... these two diagrams are both needed (same initial and same final states, so they must interfere quantum mechanically!)

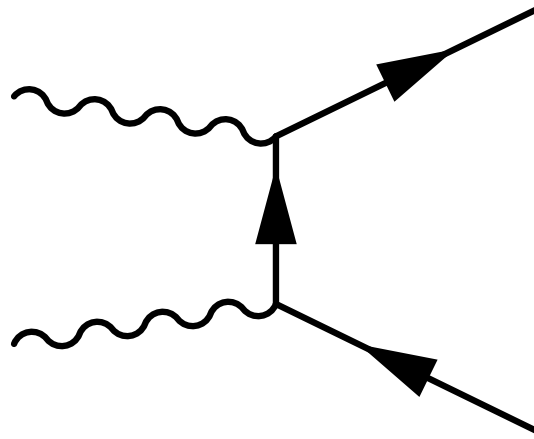
Feynman diagrams (note crossing symmetry again!)

Pair annihilation



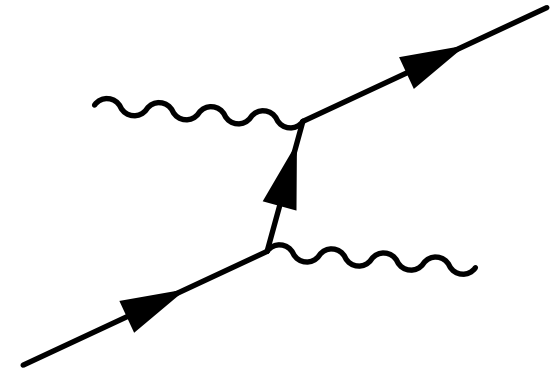
$$e^-e^+ \rightarrow \gamma\gamma$$

Pair production



$$\gamma\gamma \rightarrow e^-e^+$$

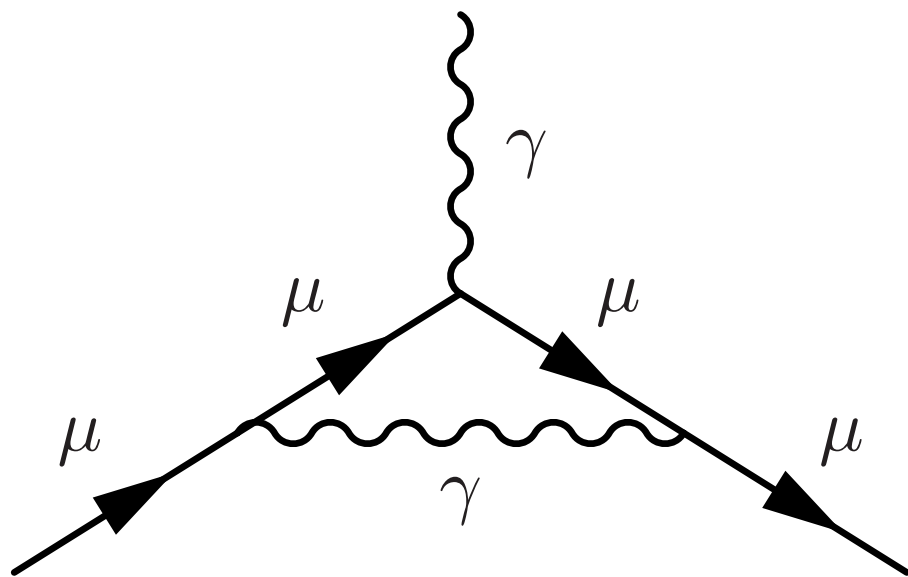
Compton scattering



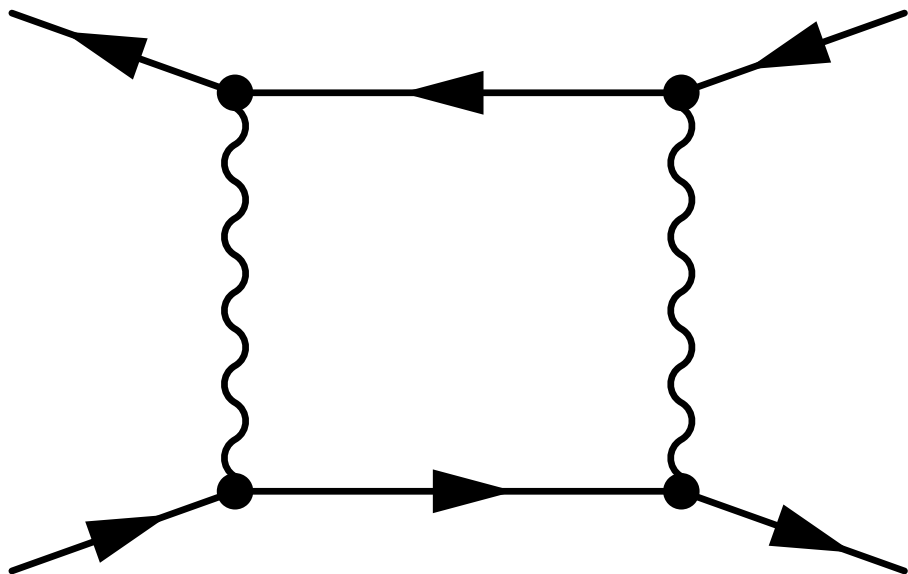
$$e^-\gamma \rightarrow \gamma e^-$$

Electrons here could be replaced by other objects with electric charge. Note that Feynman diagrams so far just look like pictorial ways to represent processes, but they can be used for much more as we'll see

We can start to combine these...

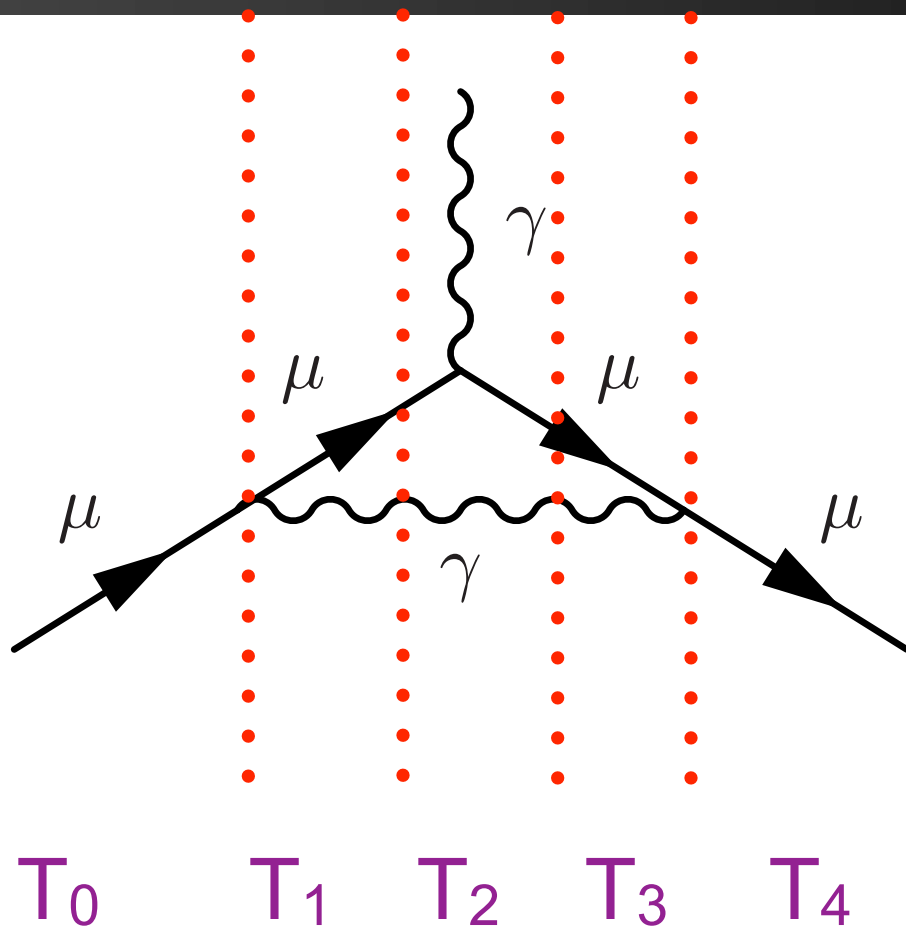


What is this? Does it remind you of any experiment that NIU works on? :)



Modification to Møller scattering (same initial and final states, different diagram)

Let's take a closer look at these diagrams

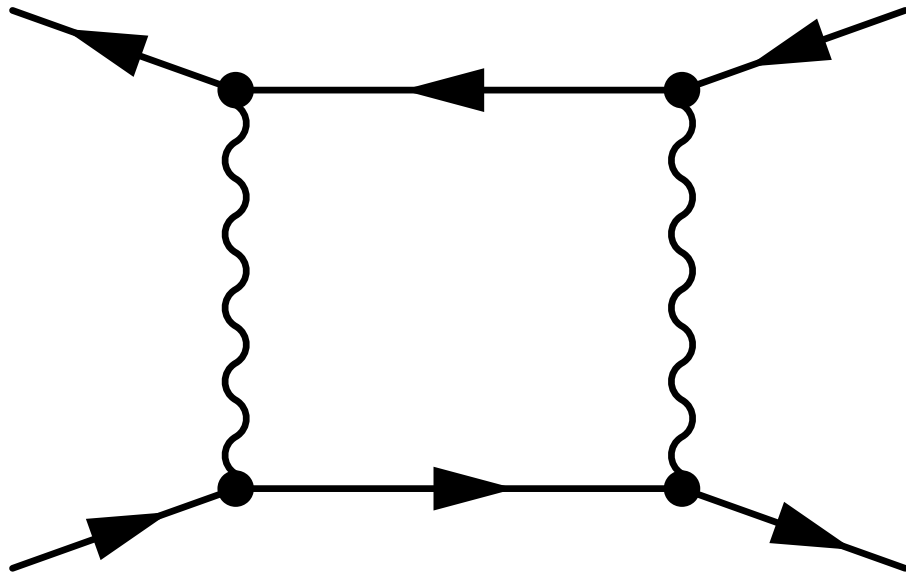


- T_0 : Single initial muon
- T_1 : Muon after interaction with photon, and that photon
- T_2 : Some external photon, muon after interaction with one or two photons, the internal photon
- T_3 : Muon after interaction with external photon, before reabsorbing the internal photon
- T_4 : Single final muon

Let's slice this up into different time bins

Can cleanly divide into **internal** and **external** particles

Let's take a closer look at these diagrams

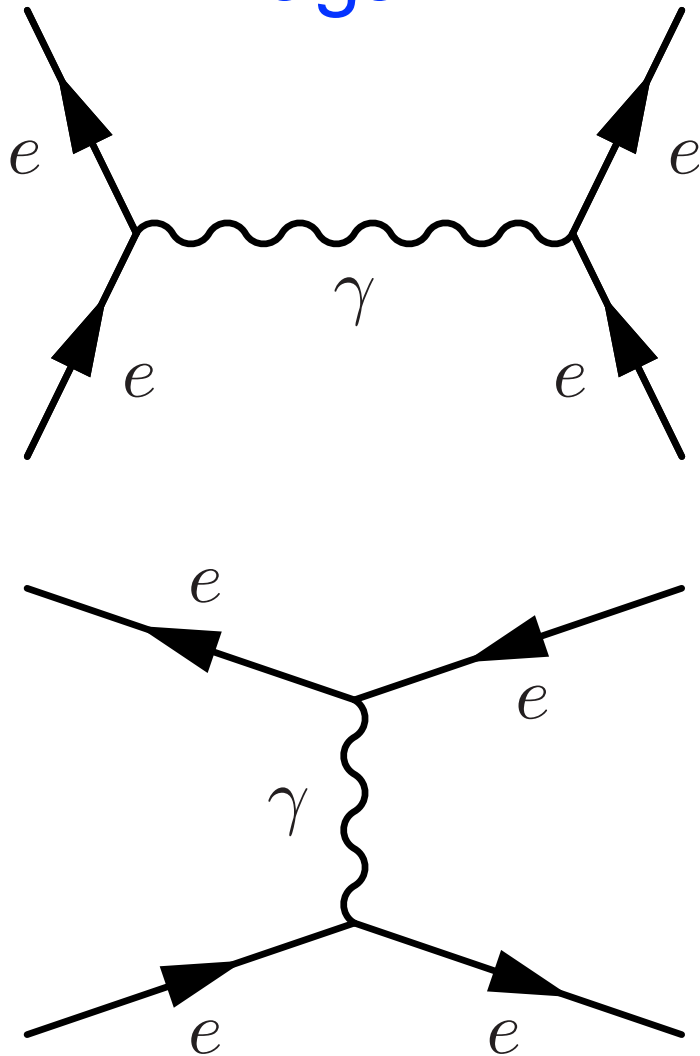


There are eight different particles in this diagram. Four of them are **external** to the diagram and **can be observed**.

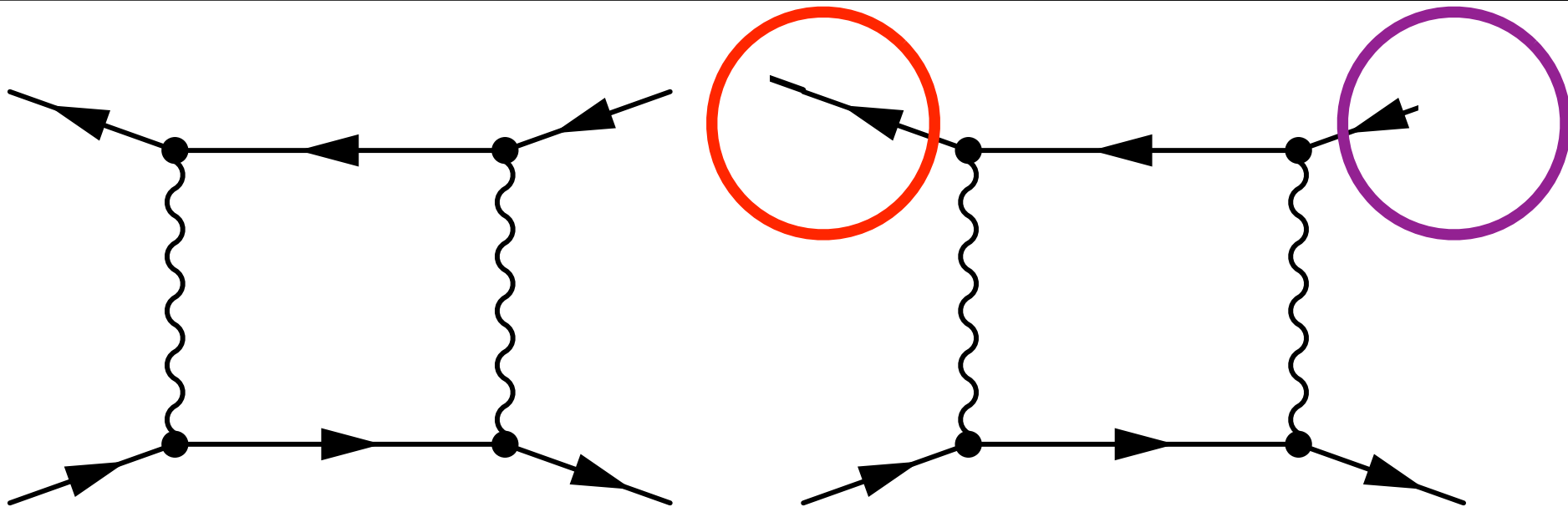
The other four are **internal** to the diagram and are called **virtual particles** and do **not** have to have on-shell mass

On virtual particles

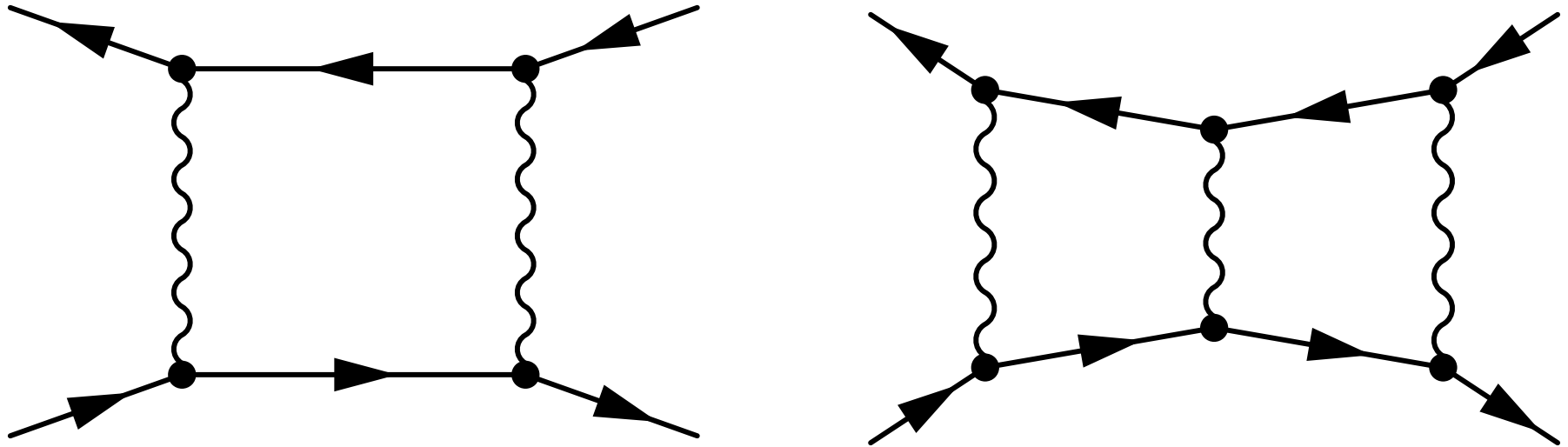
Bhabha scattering again



Virtual particles here (the photon in each diagram) cannot be observed, since the initial and final states are the same, so these must have interfering matrix elements

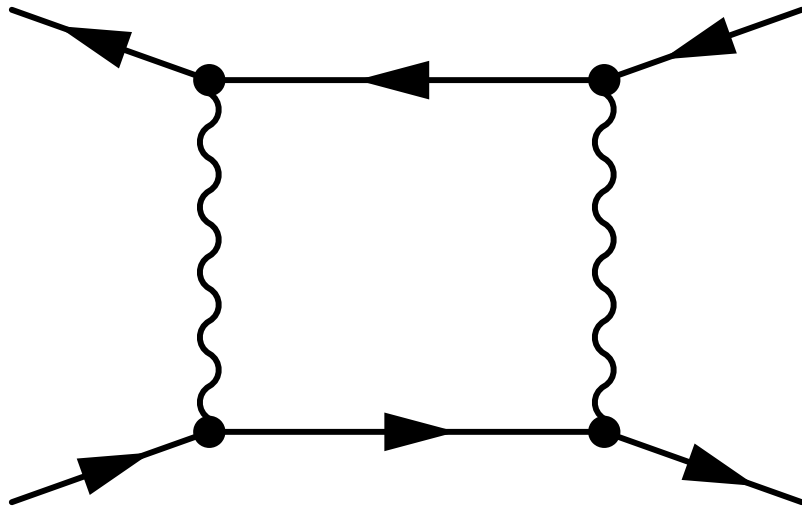


Have **extended one line** and **shrunk another**. Could have spaced things further apart or closer together - we do not differentiate between these sorts of things, nor to the angle of the lines, etc

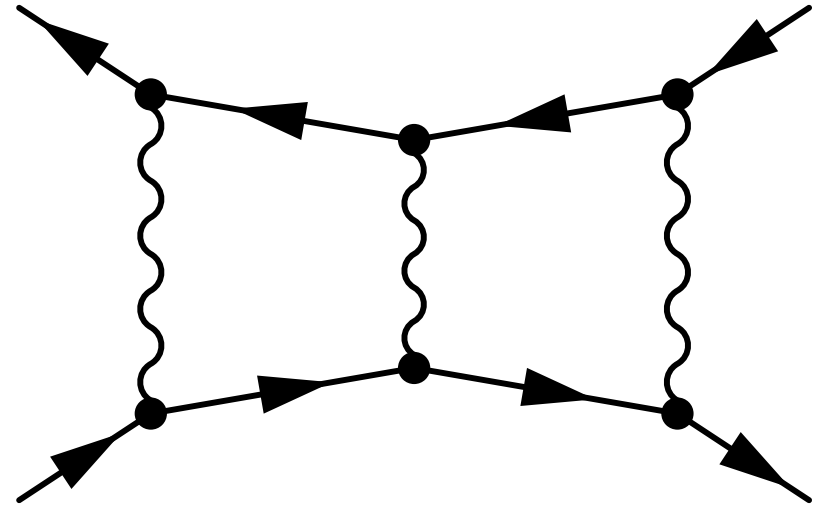


Again, can't distinguish these diagrams.
So the observable scattering amplitude
must be due to the sum of all such
diagrams: but there are an infinite
number of them!

On diagrams

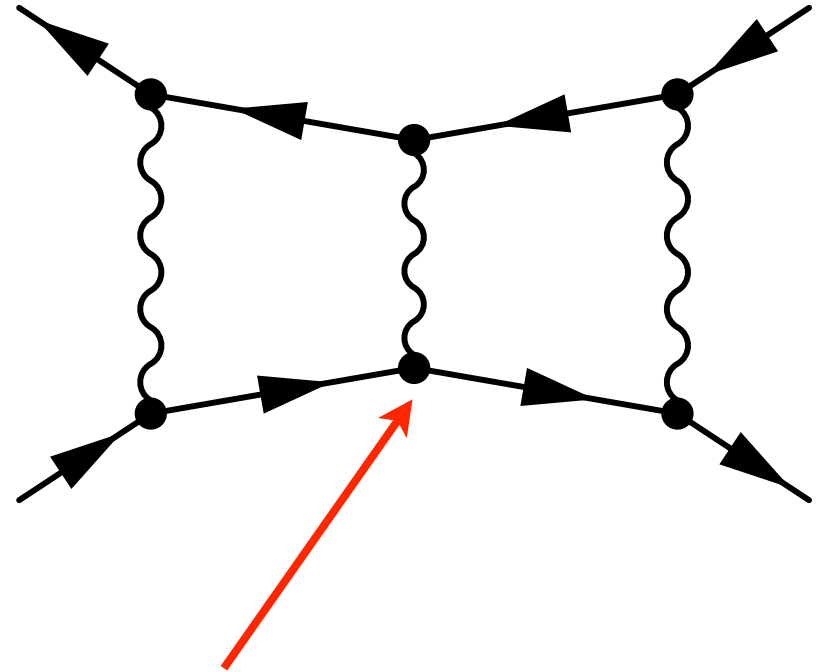
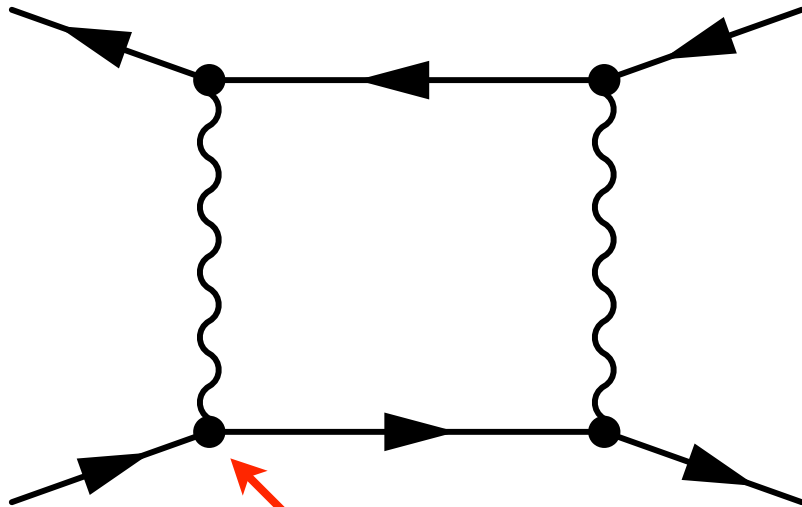


$$e^+e^- \rightarrow e^+e^- \rightarrow e^+e^-$$

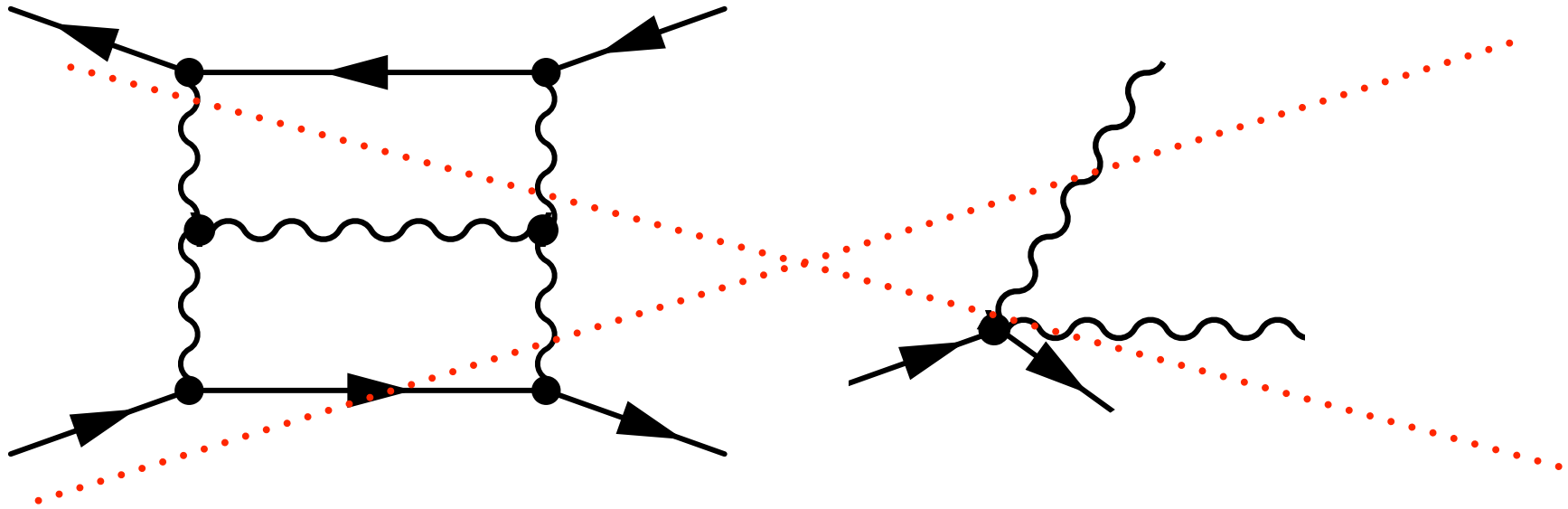


$$e^+e^- \rightarrow e^+e^- \rightarrow e^+e^- \rightarrow e^+e^-$$

$e^+e^- \rightarrow e^+e^-$ in both cases, since intermediate states are not observable. You might ask why we're bothering to draw pictures (I am not much of an artist), but they will shortly become an invaluable computation tool. Each of these represents a single number!



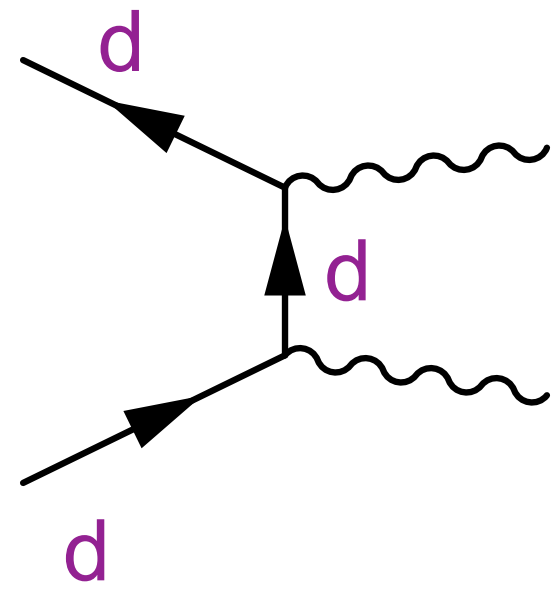
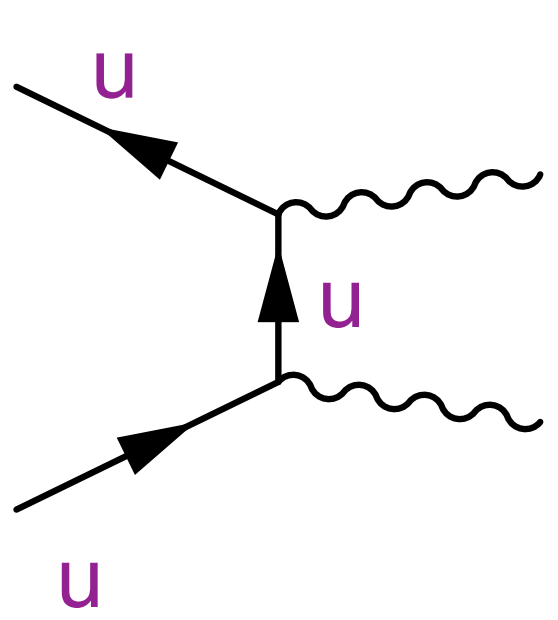
You may have noticed that I started putting circles at each interaction point - these are **vertices** and describe a point in space-time where an interaction occurs.



There is no such photon-photon-photon vertex in the SM (of course, you should know that already - the photon does not carry electric charge)

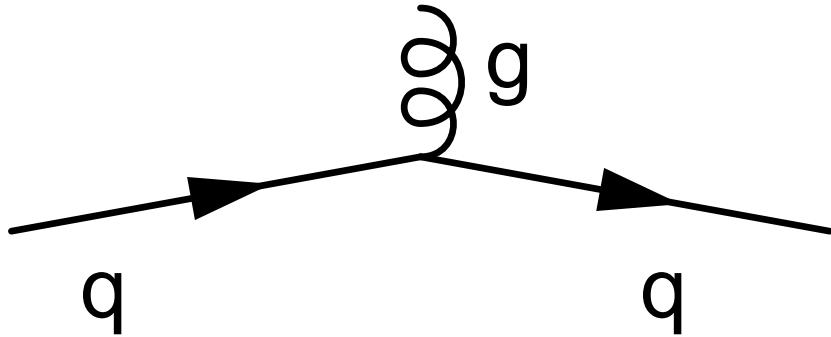
Muons can emit multiple photons, but not at the same vertex (note that the solid lines here represent any particle with electric charge)

Pair annihilation again, with quarks

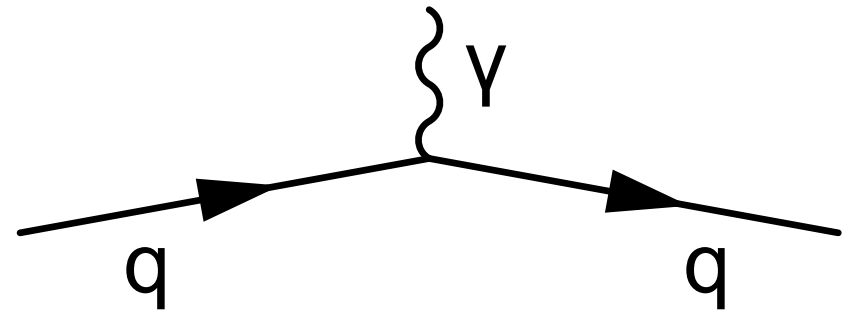


Now we can see why neutral pions decay 99% of the time to a pair of photons, and also why this happens so quickly (any good answers?)

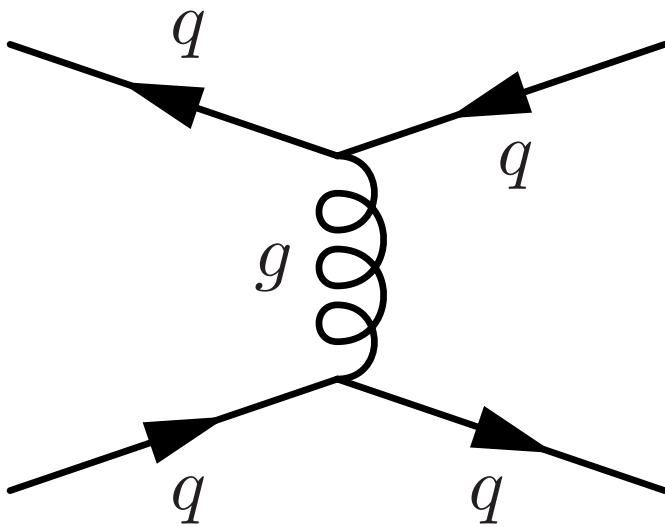
Start with common building block:



Note the similarity with the common building block of QED:

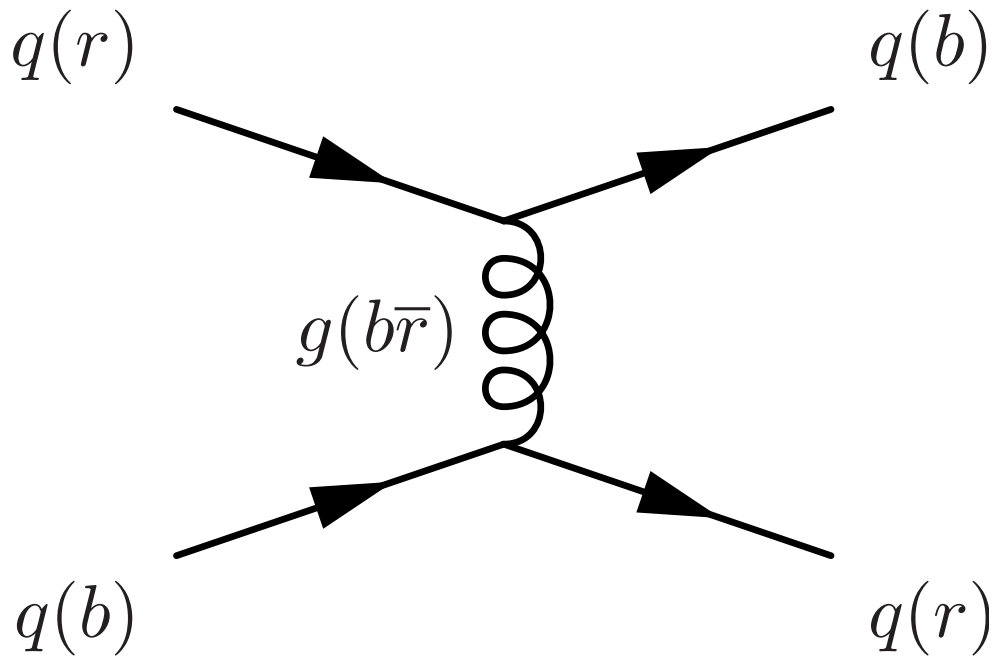


Glueon exchange between quark and anti-quark:



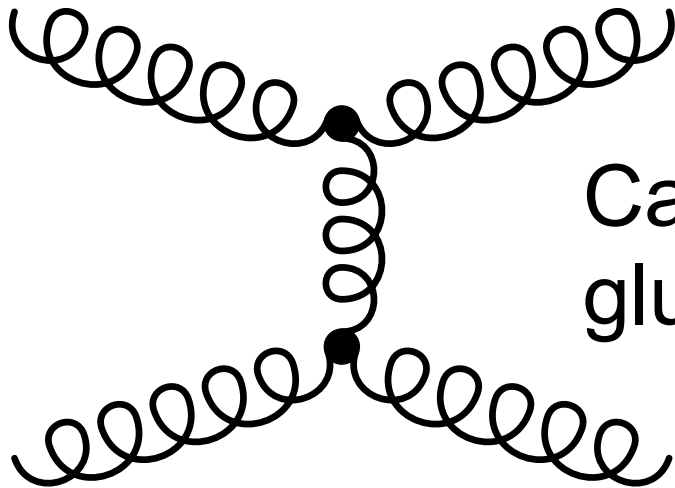
A few
fundamental
differences,
though, between
QED and QCD...

Gluon exchange between quarks:



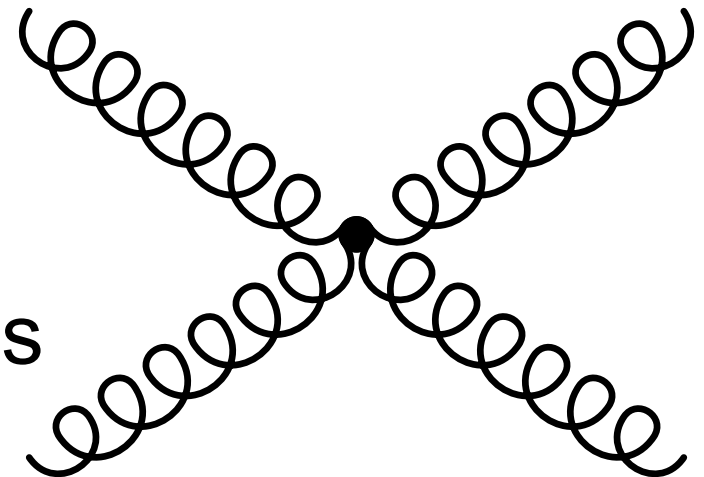
Remember that quarks have color (and so do gluons!) Try and follow the color lines (always a good thing to check)

Remember that gluons themselves have QCD color, so they self-couple!



Can have three-gluon vertices

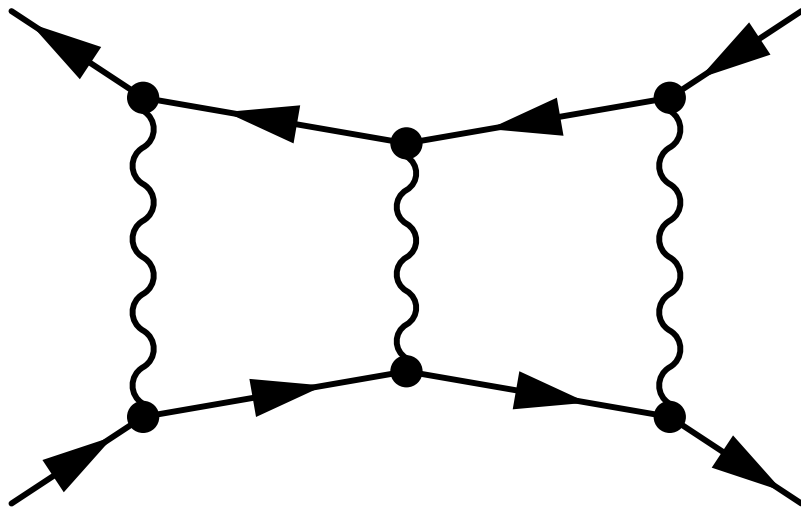
And four-gluon vertices



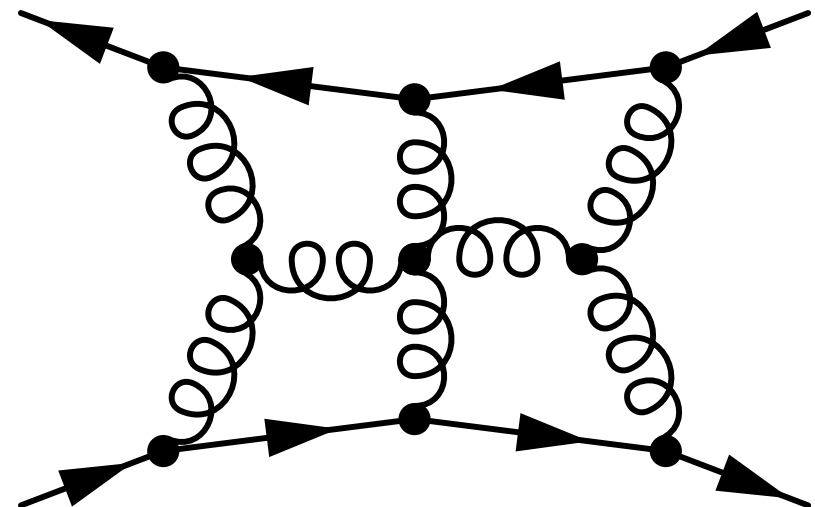
One other QED vs QCD differences

Each vertex describes an emission/absorption of a force carrier.

In QED, the probability of this is governed by the unit-less number $\alpha=1/137$, which is good, since it means the probability for many of these is small



In QCD, the probability of this is governed by the unit-less number $\alpha_s > 1$, which means the diagrams with more vertices contribute more!





Running of the coupling constant

The value of the coupling constant is indeed strong, but only at large distances. At small distances (such as inside of a proton), the strong coupling constant gets very small and we can calculate quantities

The property that the coupling gets weaker at smaller distances (aka at larger energies) is known as **asymptotic freedom**



Nobel prize for
Wilczek, Gross and
Politzer

As Griffiths points out...

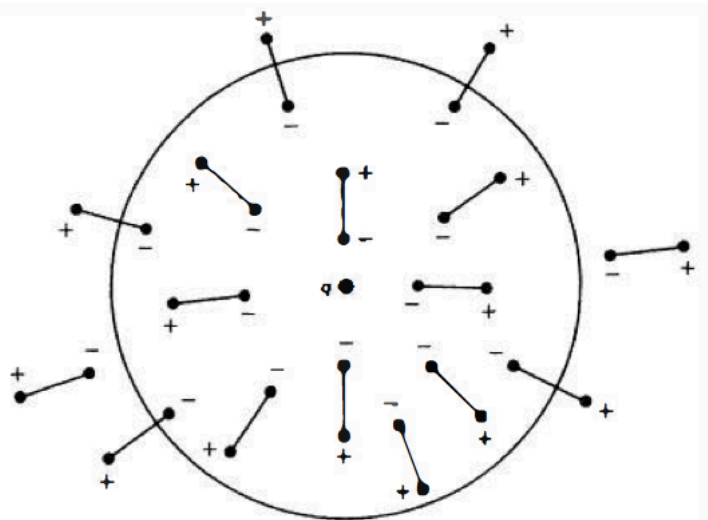


Fig. 2.1 Screening of a charge q by a dielectric medium.

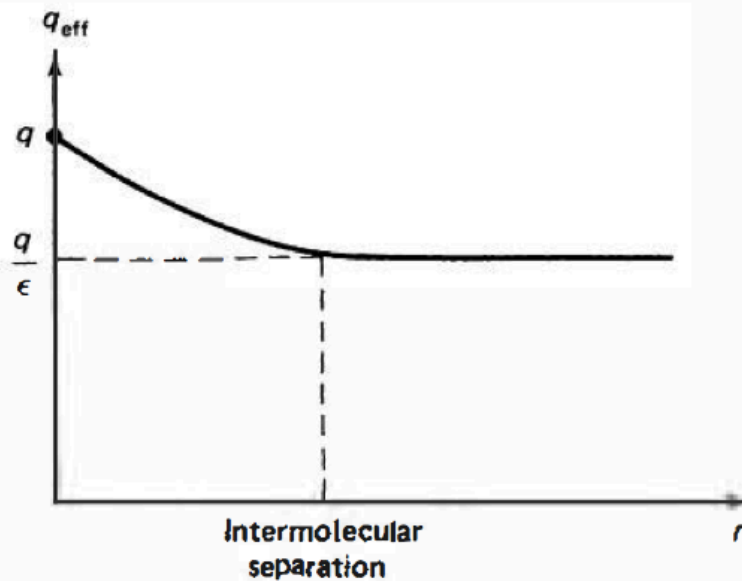
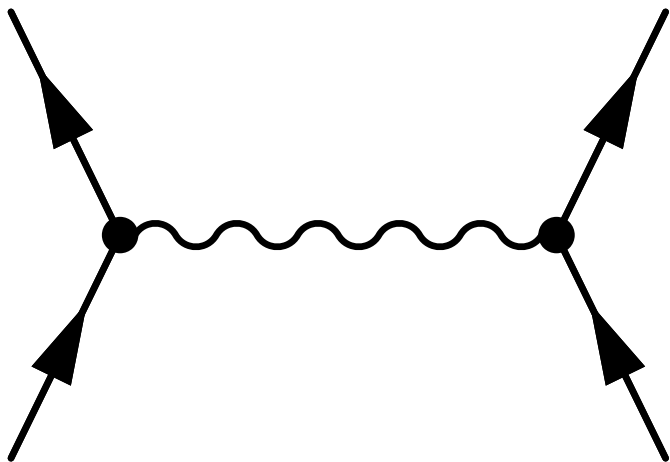


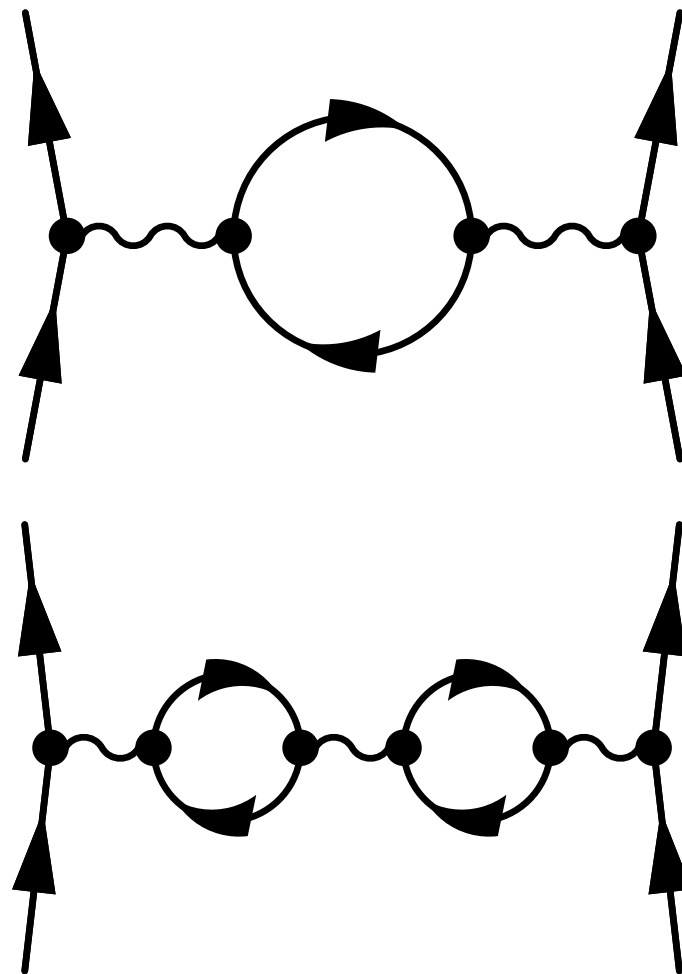
Fig. 2.2 Effective charge as a function of distance.

In classical E&M, the coupling is given by the charge q , and this is reduced at large distances by a dielectric

Bhabha scattering again



Virtual diagrams contributing



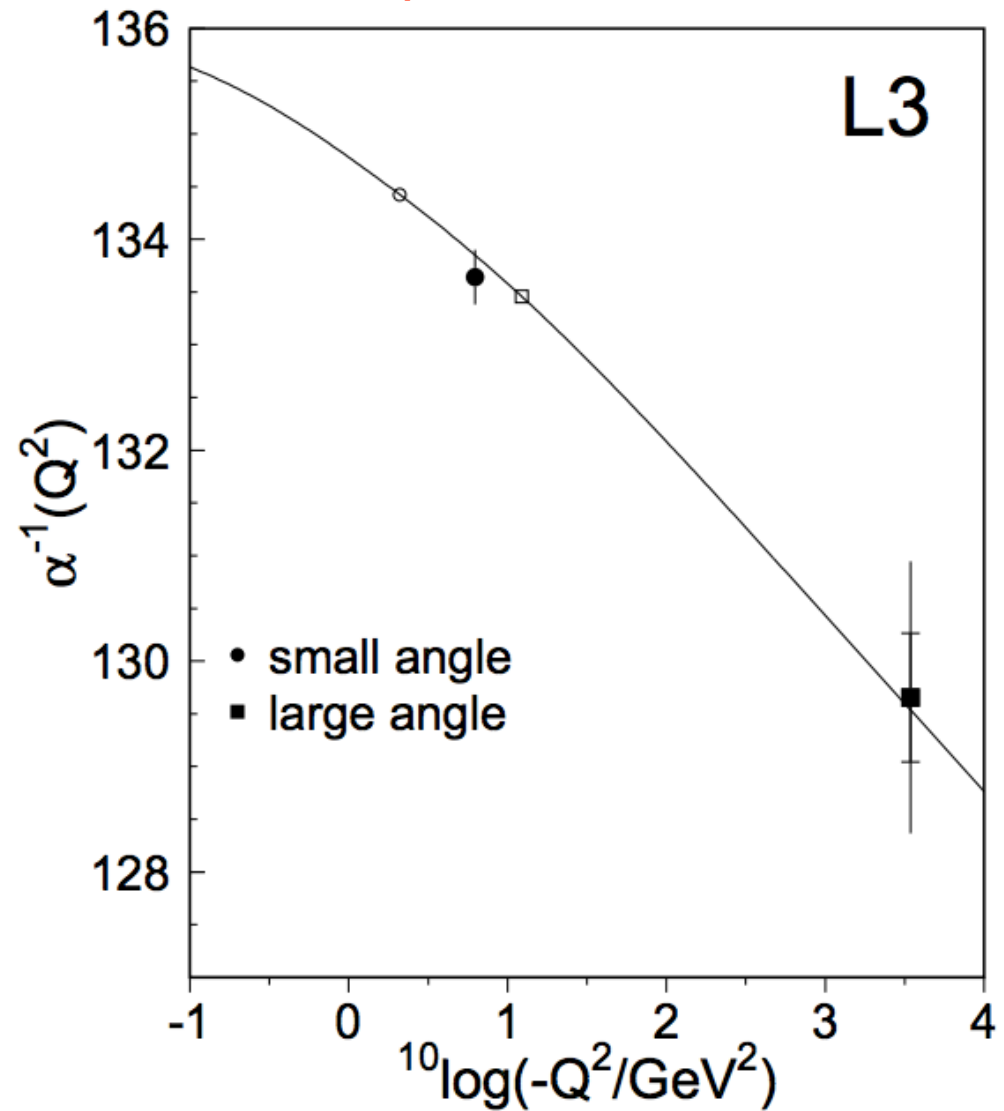
Running of the E&M fine structure constant

~1/137 at zero
momentum
transfer

~1/127 at the Z
boson mass!

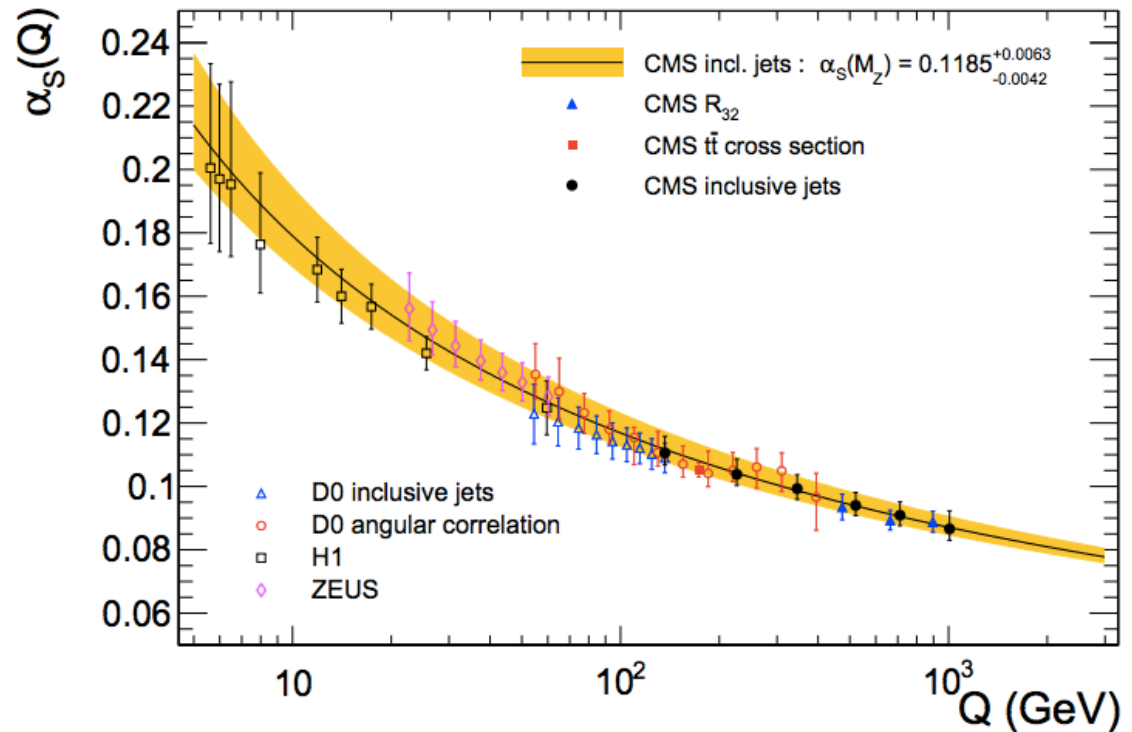
These are vacuum
polarization effects

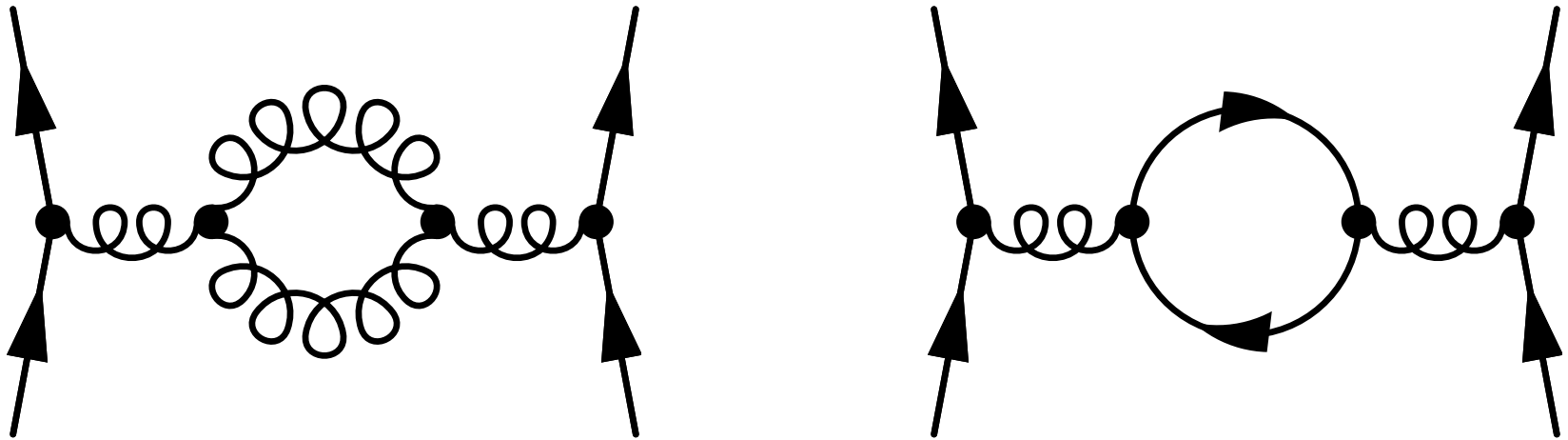
hep-ex/0002035v1



At 1 GeV, it's already ~ 0.5 , and approaches 1 at ~ 250 MeV

arXiv:1410.6765





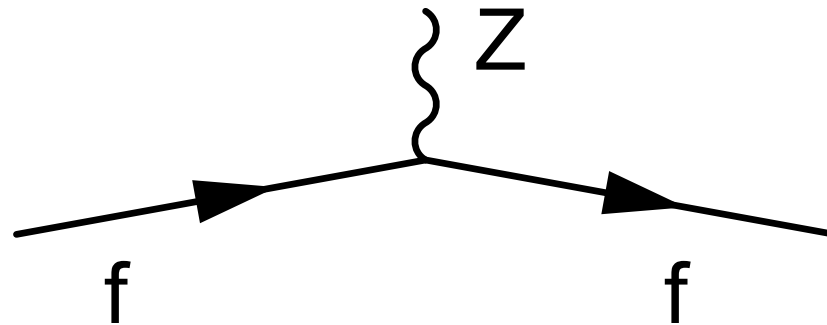
Effects from quark polarization and gluon polarization give opposite effects!

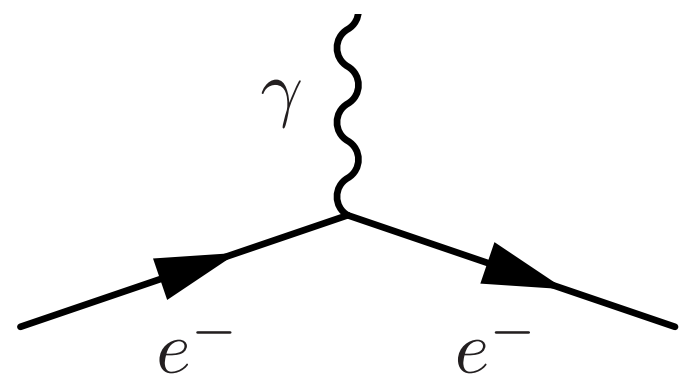
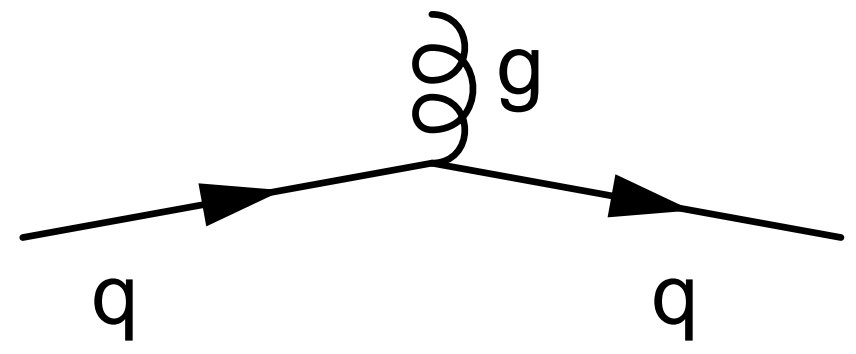
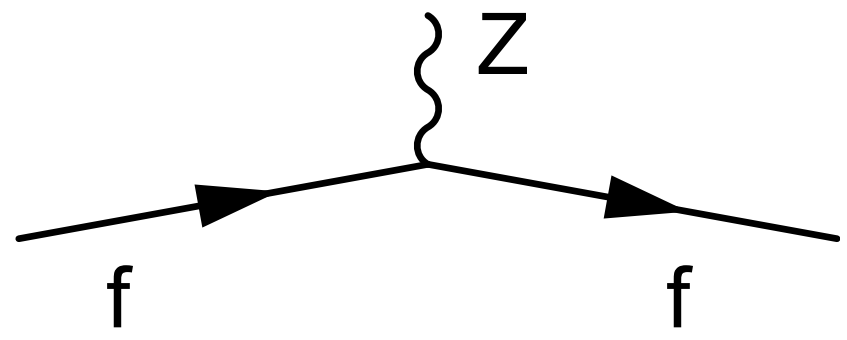
$a=2f-11n$ (f =number of flavors, n =number colors)

defines whether coupling increases or decreases at short distances

$f=6, n=3 \rightarrow a < 0 \rightarrow$ asymptotic freedom (upper limit on no more than 17 generations of quarks)

where f here
is any fermion (Z boson couples to all quarks,
charged leptons and neutral leptons)

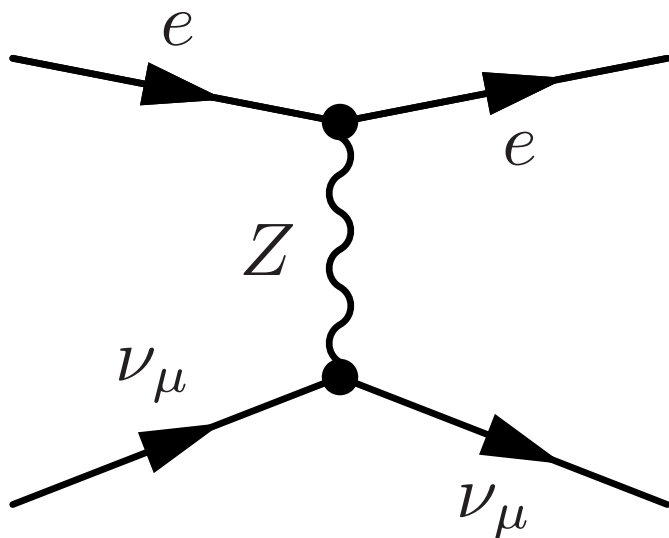




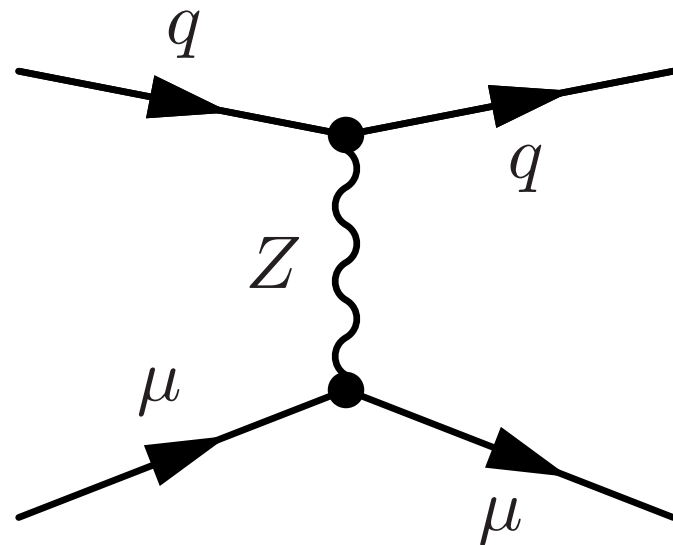
Griffiths uses a jagged straight line for EW bosons, but I will not. Just be careful, regardless of your choice!

What processes can the Z boson mediate?

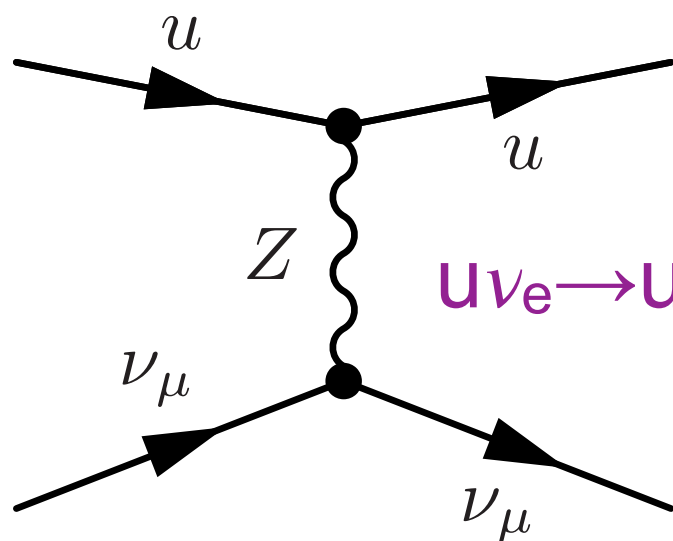
$e^- \nu_e \rightarrow e^- \nu_e$



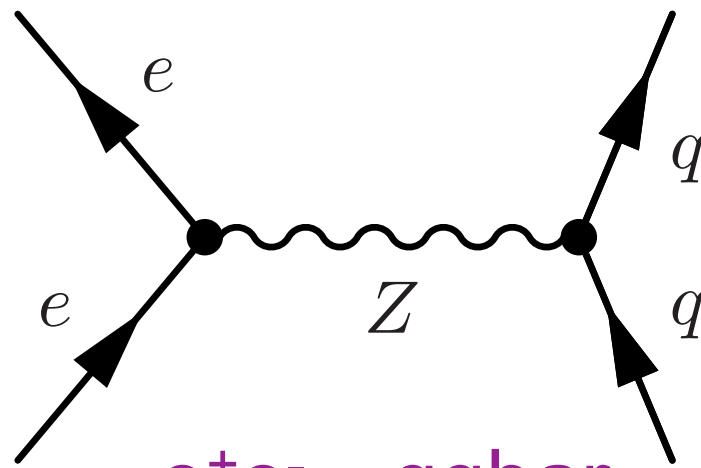
$\mu^- q \rightarrow \mu^- q$



$U \nu_e \rightarrow U \nu_e$

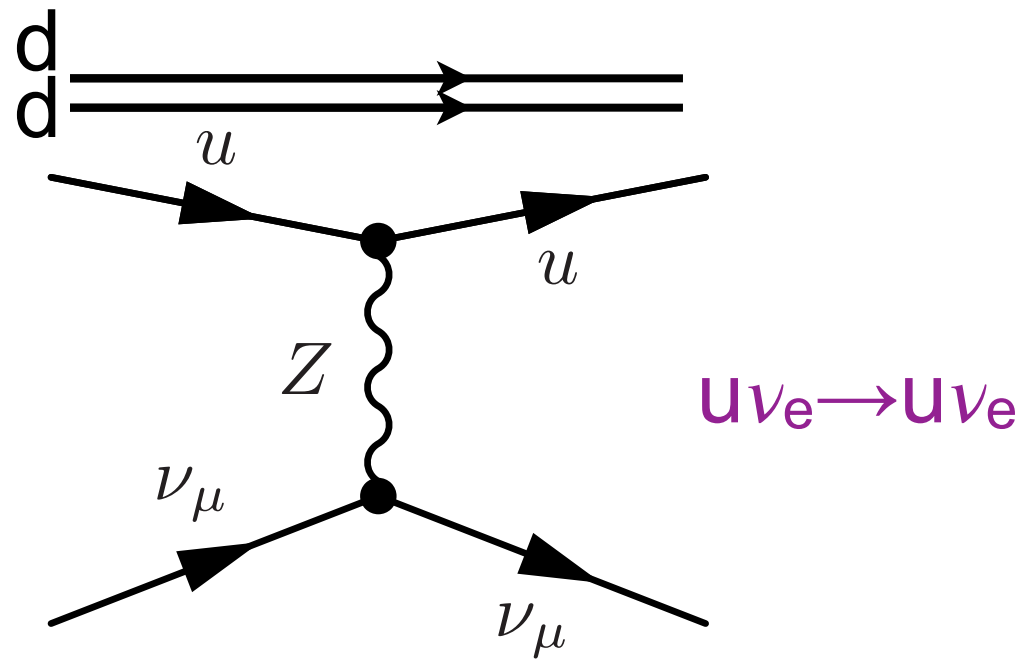


$e^+ e^- \rightarrow q q \text{bar}$

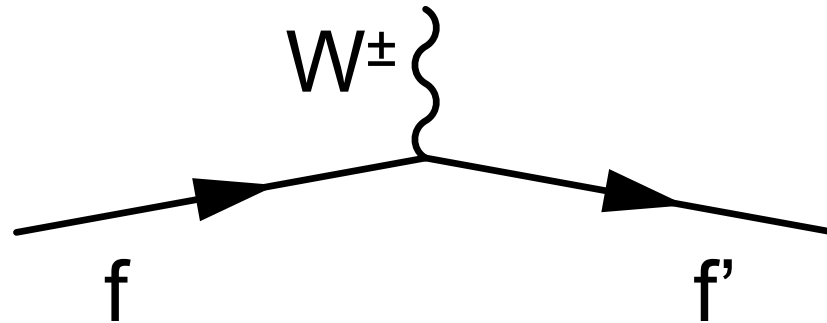


What processes can the Z boson mediate?

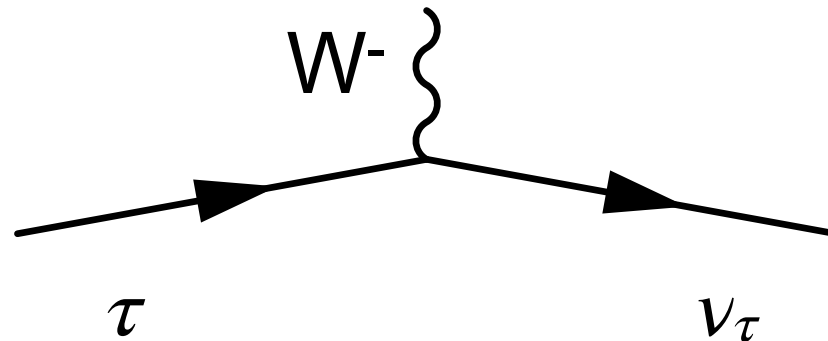
We of course don't scatter the bare quarks, but quarks inside hadrons (here inside a neutron)



Remember that W bosons are the only ones involved in change of flavor or type of matter



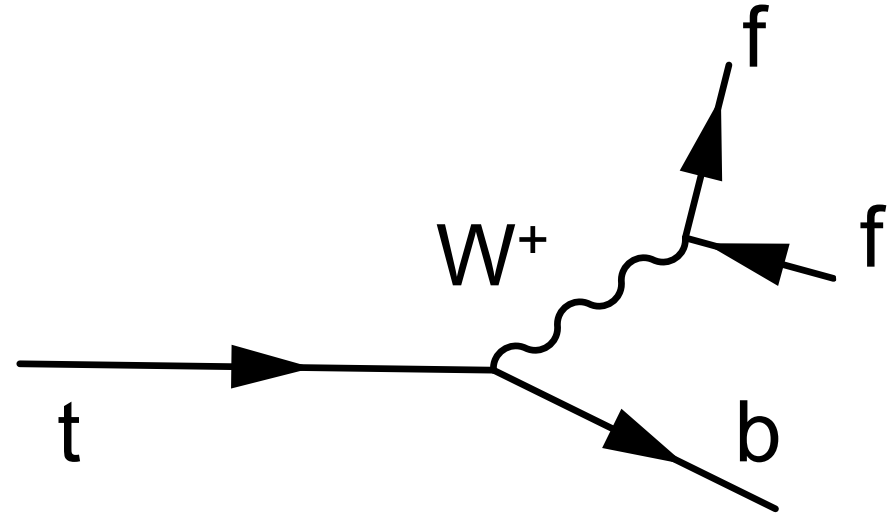
Remember that W bosons are the only ones involved in change of flavor or type of matter. Example, a tau converts into a tau neutrino and a W boson:



Charged weak interactions

Top quarks decay nearly $\sim 100\%$ of the time to a W boson and a bottom quark. The W boson can decay to a quark-antiquark pair, or to a charged lepton and a neutrino

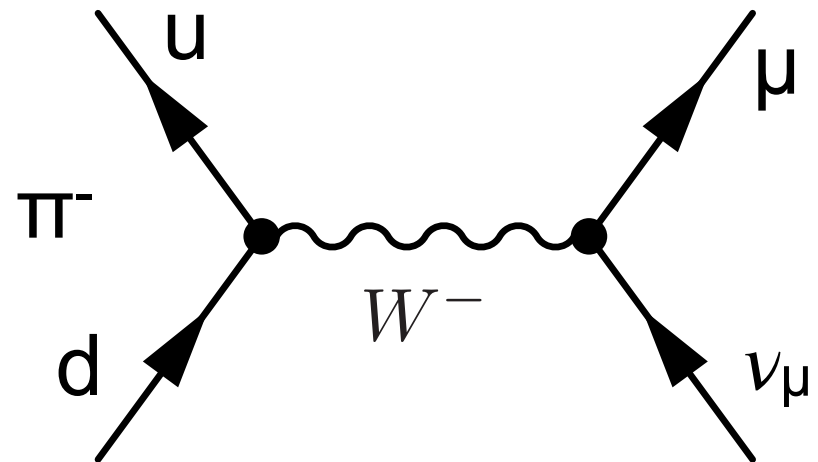
Check lepton
flavor numbers
and EM
charge!



Charged weak interactions

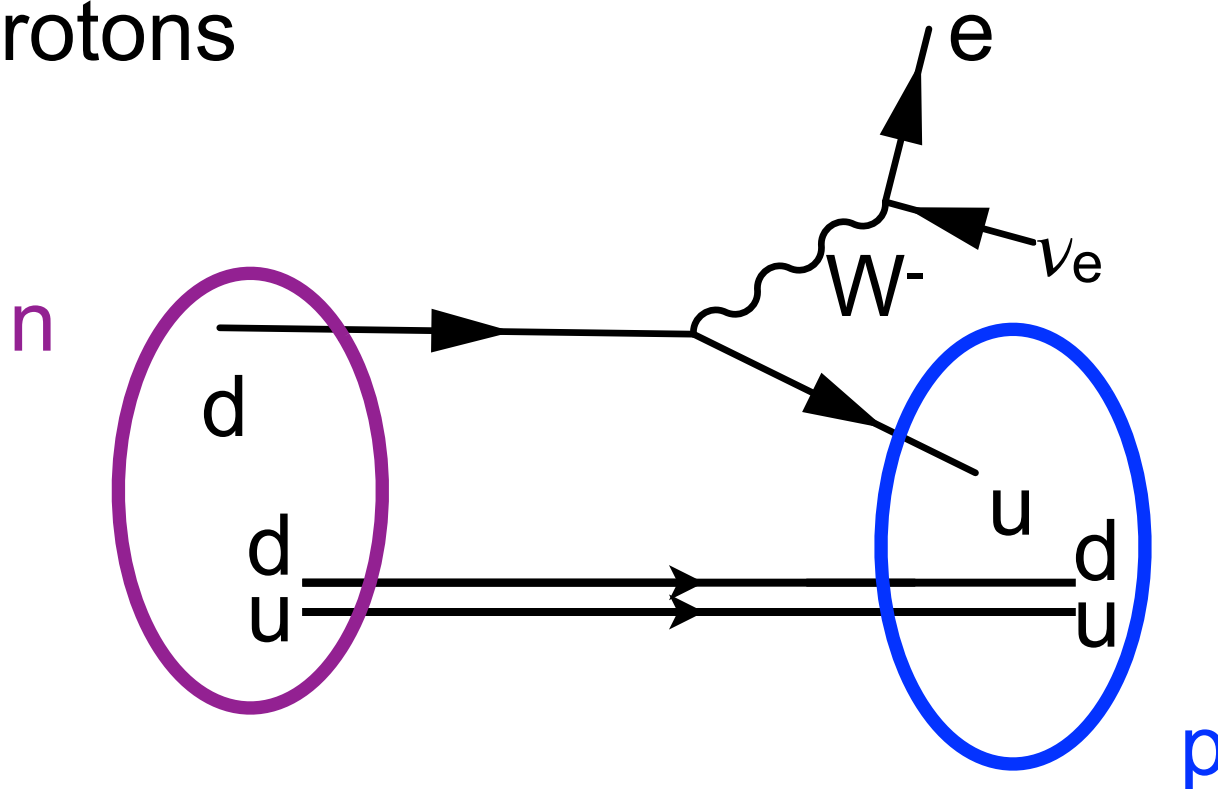
This is how a charged pion can decay to a muon (and an anti-muon neutrino). Could have guessed W boson involvement, because it involves a change of flavor (here charged pion decay is much slower than neutral pion decay)

Check lepton flavor numbers and EM charge!



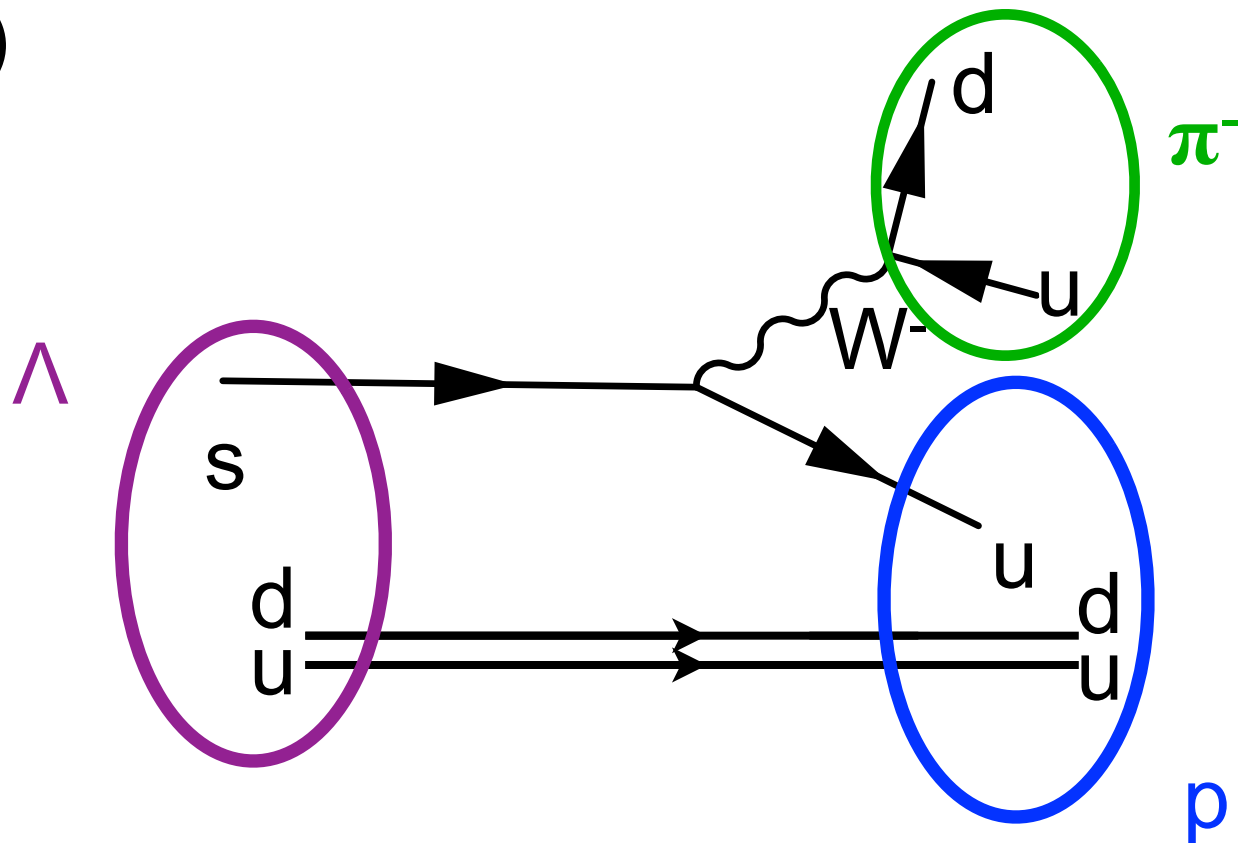
What processes can the W boson mediate?

We of course don't scatter the bare quarks, but quarks inside hadrons (here inside a neutron)
This is how neutrons beta-decay to protons



The important question...

How do we account for strangeness-changing weak interactions? (We know they have to come via the weak force, but so far we have only seen decays within a matter generation)



The Kobayashi-Maskawa matrix

The weak force couples not to eigenstates of quarks that we observe, but to a slightly skewed set instead:

What we observe

$$\begin{pmatrix} u \\ d \end{pmatrix} \quad \begin{pmatrix} c \\ s \end{pmatrix} \quad \begin{pmatrix} t \\ b \end{pmatrix}$$

What weak force couples to

$$\begin{pmatrix} u \\ d' \end{pmatrix} \quad \begin{pmatrix} c \\ s' \end{pmatrix} \quad \begin{pmatrix} t \\ b' \end{pmatrix}$$

$$\begin{pmatrix} u' \\ d' \\ s' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} u \\ d \\ s \end{pmatrix}$$

What we observe

$$\begin{pmatrix} u \\ d \end{pmatrix} \quad \begin{pmatrix} c \\ s \end{pmatrix} \quad \begin{pmatrix} t \\ b \end{pmatrix}$$

What weak force couples to

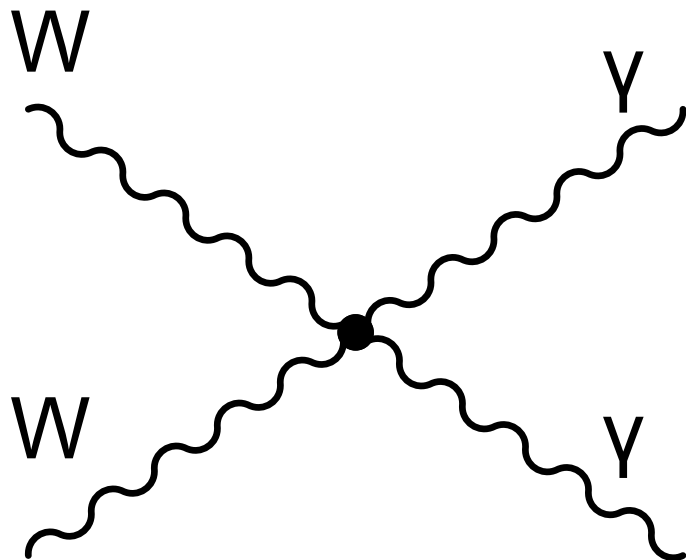
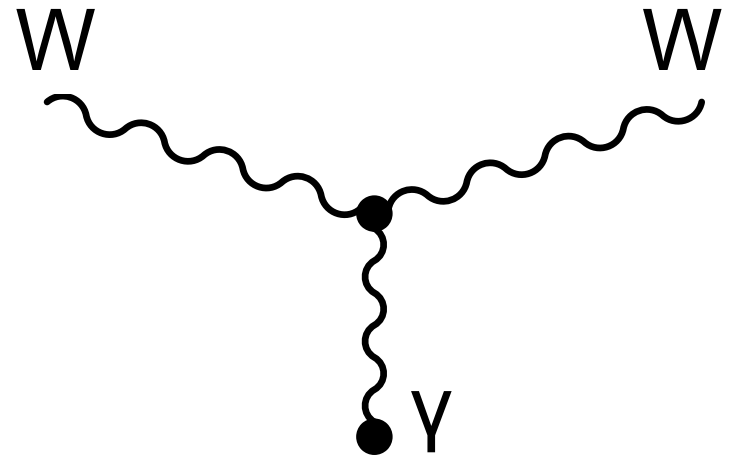
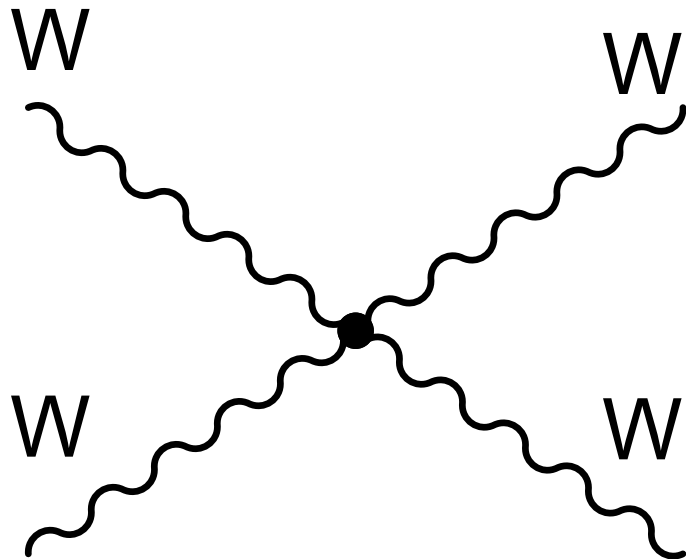
$$\begin{pmatrix} u \\ d' \end{pmatrix} \quad \begin{pmatrix} c \\ s' \end{pmatrix} \quad \begin{pmatrix} t \\ b' \end{pmatrix}$$

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

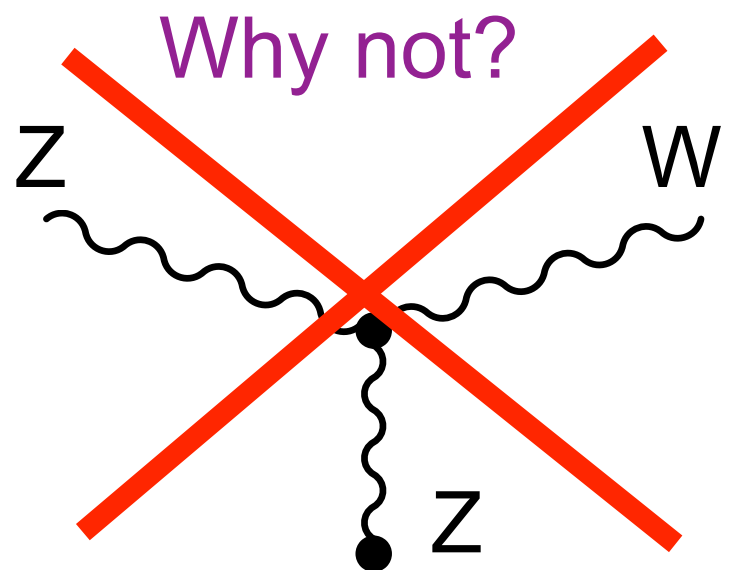
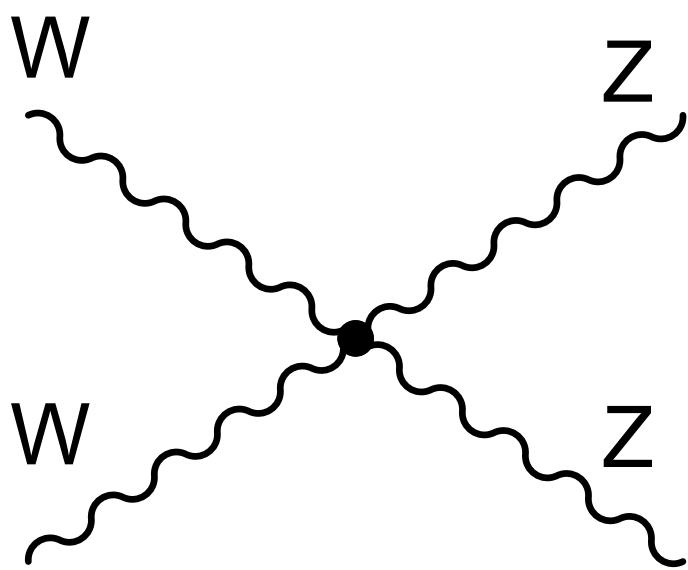
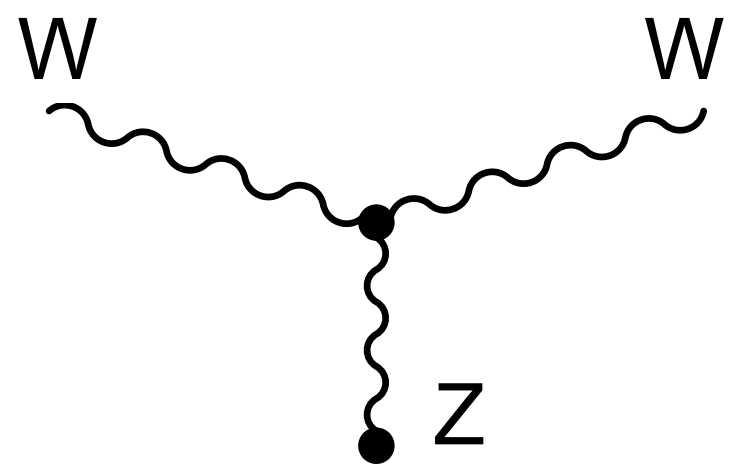
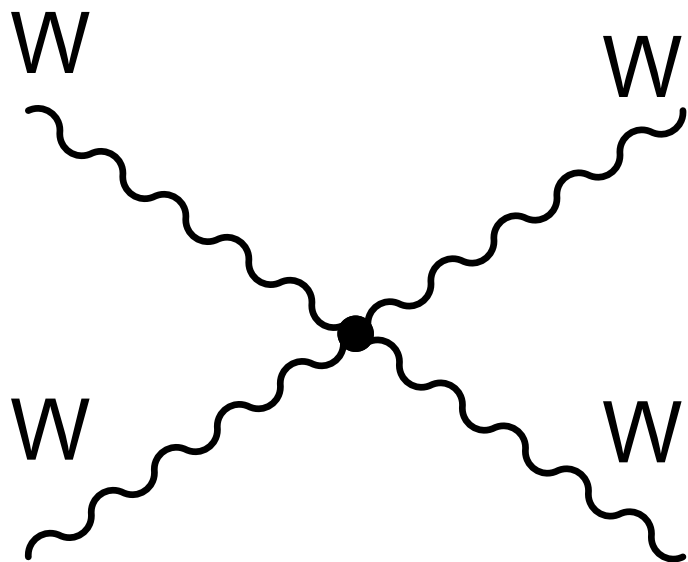
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} 0.974 & 0.225 & 0.0041 \\ 0.225 & 0.986 & 0.041 \\ 0.0084 & 0.040 & 1.021 \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix} \quad \text{2014 PDG}$$

Off-diagonal numbers are small, which has implications for which decays happen more/less frequently. **The matrix has to be unitary.** Note that I've cheated a bit, because there are phases in the matrix that I have ignored so far

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} 0.974 & 0.225 & 0.0041 \\ 0.225 & 0.986 & 0.041 \\ 0.0084 & 0.040 & 1.021 \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix} \quad \begin{matrix} 2014 \\ \text{PDG} \end{matrix}$$



In general, can always replace a photon by a Z boson in a diagram.
What about the reverse?



Particles... like to decay to lighter objects, unless there is a reason that they can't decay (such as a conservation law)

What is the most obvious reason for non-decay?

$$p^+ \rightarrow n e^+ \nu_e$$

Hint, look at the energy in the
rest frame

Momentum conservation

Angular momentum conservation

Conservation of charge

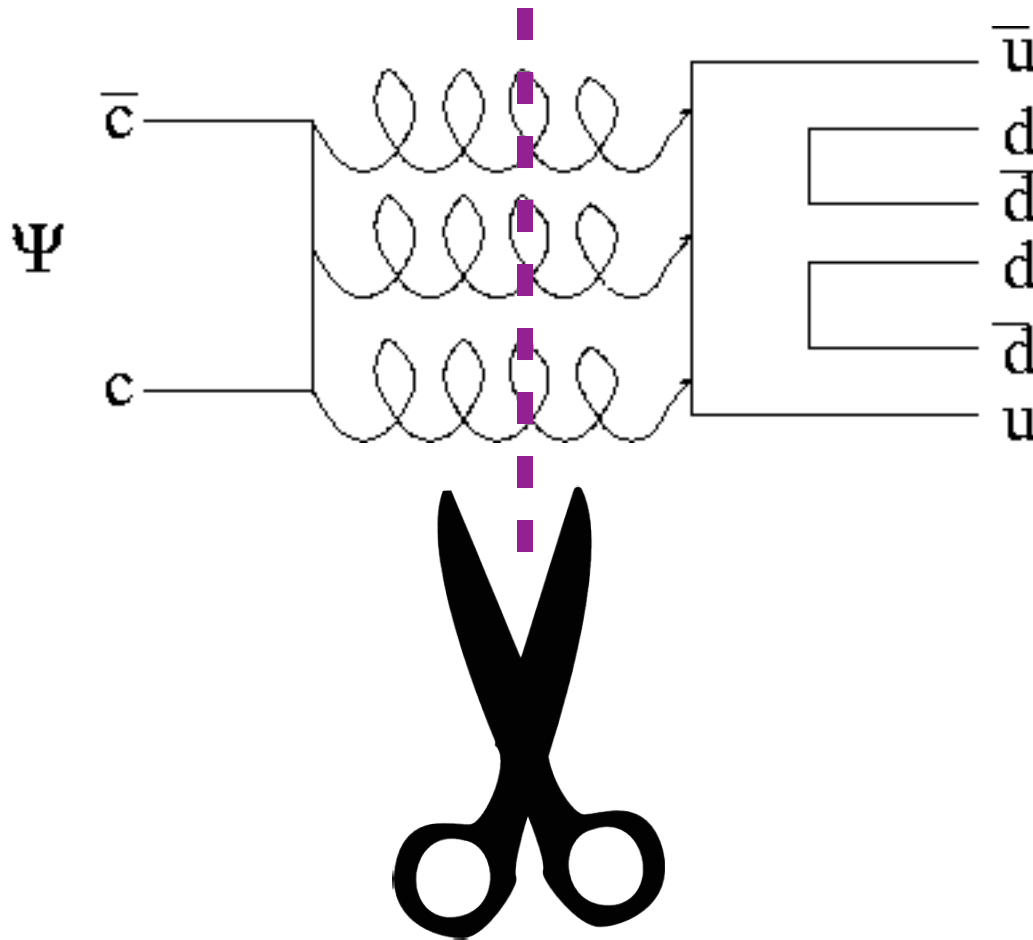
Conservation of color

Conservation of baryon number

Conservation of lepton (electron/muon/tau)
number

Quark flavor (not quite conserved)

<http://www.ippp.dur.ac.uk/~krauss/Lectures/QuarksLeptons/QCD/Equations/jpsidecay.gif>



OZI rule explains why the J/Psi takes so long to decay.
 What does this have to do with **asymptotic freedom**?

Photons can't decay (they already have zero mass - nothing lighter to decay into)

Protons are the lightest baryon, so baryon conservation tells us that they are stable

This is why the world is made up of protons, neutrons (which are stable inside nuclei), electrons and neutrinos

Decays are probabilistic - you never know when an object will decay, but you can say on average how long a type of object will take to decay

$$\mu: 2.2 \times 10^{-6} \text{ s}$$

$$\pi^+: 2.6 \times 10^{-8} \text{ s}$$

$$\pi^0: 8.3 \times 10^{-17} \text{ s}$$

$$\tau: 2.9 \times 10^{-13} \text{ s}$$

$$n \text{ (free)}: 880 \text{ s}$$

$$\text{top quark}: 4 \times 10^{-25} \text{ s}$$

Many of these objects will decay to multiple final states with different, predicted **branching ratios**

Why so different?

Strong decays: 10^{-23} s (top quark)

EM decays: 10^{-16} s (π^0)

Weak decays: longer (neutron and muon)

The more (less) phase space for a decay, the faster (slower) it will occur.

See the free neutron





Photos: Internet screengrabs

OK, not a real physics explanation, but a good way to remember things