

Assignment: HW2 [40 points]

Assigned: 2009/09/05

Due: 2009/10/12

## Solutions

**P2.1** [1 + 3 + 3 + 3 + 2 = 12 points]

Let the motion of a point mass be governed by the law

$$\ddot{\mathbf{r}} = \dot{\mathbf{r}} \times \mathbf{a}, \quad \mathbf{a} = \text{const.} \quad (1)$$

- (a) Show that  $\dot{\mathbf{r}} \cdot \mathbf{a}$  is constant in time.
- (b) Reduce Eq. 1 to an inhomogeneous differential equation of the form  $\ddot{\mathbf{r}} + \omega^2 \mathbf{r} = f(t)$ .
- (c) Solve the above equation by using a particular function of the form  $\mathbf{r}_p = \mathbf{c}t + \mathbf{d}$ .
- (d) Express the constants of integration in terms of the initial values  $\mathbf{r}(0)$  and  $\dot{\mathbf{r}}(0)$ .
- (e) Describe the curve  $\mathbf{r}(t) = \mathbf{r}_g(t) + \mathbf{r}_p(t)$ , where  $\mathbf{r}_g(t)$  is the solution of the homogeneous equation  $\ddot{\mathbf{r}} + \omega^2 \mathbf{r} = 0$ .

**S2.1** (a)

$$\frac{d}{dt}(\dot{\mathbf{r}} \cdot \mathbf{a}) = \ddot{\mathbf{r}} \cdot \mathbf{a} = (\dot{\mathbf{r}} \times \mathbf{a}) \cdot \mathbf{a} = 0. \quad (2)$$

(b)

$$\frac{d}{dt}\ddot{\mathbf{r}} = \ddot{\mathbf{r}} \times \mathbf{a} = (\dot{\mathbf{r}} \times \mathbf{a}) \times \mathbf{a} = -(\mathbf{a} \cdot \mathbf{a})\dot{\mathbf{r}} + (\dot{\mathbf{r}} \cdot \mathbf{a})\mathbf{a}. \quad (3)$$

We have shown the second term to be constant. Thus, integrating this equation from 0 to  $t$ , we get

$$\ddot{\mathbf{r}}(t) - \ddot{\mathbf{r}}(0) = -\omega^2(\mathbf{r}(t) - \mathbf{r}(0)) + (\dot{\mathbf{r}}(0) \cdot \mathbf{a})\mathbf{a}t, \quad (4)$$

where  $\omega^2 = \mathbf{a} \cdot \mathbf{a} = a^2$ . Since  $\ddot{\mathbf{r}}(0) = \dot{\mathbf{r}}(0) \times \mathbf{a}$ , we may write

$$\ddot{\mathbf{r}}(t) + \omega^2 \mathbf{r}(t) = (\dot{\mathbf{r}}(0) \cdot \mathbf{a})\mathbf{a}t + \dot{\mathbf{r}}(0) \times \mathbf{a} + \omega^2 \mathbf{r}(0). \quad (5)$$

(c) The general solution of the homogeneous differential equation is

$$\mathbf{r}_g(t) = \mathbf{c}_1 \sin(\omega t) + \mathbf{c}_2 \cos(\omega t). \quad (6)$$

So, the particular solution is simply the RHS of Eq. 5 divided by  $\omega^2$ , which is of the form

$$\mathbf{r}_p(t) = \mathbf{c}t + \mathbf{d} \quad (7)$$

- (d) We can get the constants in Eqs. 6 and 7 by comparing them with the corresponding terms in Eq. 5 and using the initial conditions

$$\begin{aligned} \mathbf{c}_1 &= \frac{1}{\omega^3} (\mathbf{a}^2 \dot{\mathbf{r}}(0) - (\dot{\mathbf{r}}(0) \cdot \mathbf{a}) \mathbf{a}) = \frac{1}{\omega^3} \mathbf{a} \times (\dot{\mathbf{r}}(0) \times \mathbf{a}), \\ \mathbf{c}_2 &= \frac{1}{\omega^2} \dot{\mathbf{r}}(0) \times \mathbf{a}, \\ \mathbf{c} &= \frac{1}{\omega^2} (\dot{\mathbf{r}}(0) \cdot \mathbf{a}) \mathbf{a}, \\ \mathbf{d} &= \frac{1}{\omega^2} \dot{\mathbf{r}}(0) \times \mathbf{a} + \mathbf{r}(0). \end{aligned} \quad (8)$$

- (e) It represents a helix winding around the vector  $\mathbf{a}$ .

**P2.2** [2 + 5 = 7 points]

Use variational calculus to find

- (a) The shortest connection between two points  $(x_1, y_1)$  and  $(x_2, y_2)$  on the 2-dimensional Euclidean plane.  
 (b) The shape of a line of uniform mass density supported between  $(x_1, y_1)$  and  $(x_2, y_2)$  in a uniform gravitational field  $\mathbf{g} = g\hat{\mathbf{e}}_y$ .

**S2.2** In the general case, where the action functional is the integral of a scalar function  $f(y, y', x)$ , with  $y' = \frac{dy}{dx}$ , the Euler-Lagrange equation reads

$$\frac{\partial f}{\partial y} = \frac{d}{dx} \frac{\partial f}{\partial y'} \quad (9)$$

Multiplying by  $y'$  and adding the term  $y'' \frac{\partial f}{\partial y'}$  to both sides we get

$$y' \frac{\partial f}{\partial y} + y'' \frac{\partial f}{\partial y'} = \frac{d}{dx} \left( y' \frac{\partial f}{\partial y'} \right) \quad (10)$$

If  $f$  does not depend explicitly on  $x$ , then  $\frac{\partial f}{\partial x} = 0$ . Adding  $\frac{\partial f}{\partial x}$  to the LHS of the above equation gives  $\frac{df(y, y')}{dx}$ , while the RHS is left unchanged:

$$\frac{df(y, y')}{dx} = \frac{d}{dx} \left( y' \frac{\partial f}{\partial y'} \right), \quad (11)$$

or,

$$y' \frac{\partial f}{\partial y'} - f(y, y') = \text{const.} \quad (12)$$

- (a) We have to minimize the arc length

$$L = \int ds = \int_{x_1}^{x_2} \sqrt{1 + y'^2} dx. \quad (13)$$

So,  $f(y, y') = \sqrt{1 + y'^2}$ , and Eq. 12 gives

$$y' \frac{\sqrt{y'}}{\sqrt{1 + y'^2}} - \sqrt{1 + y'^2} = \text{const.} \quad \text{or,} \quad y' = \text{const.} \quad (14)$$

Thus,  $y = ax + b$ . Applying the boundary conditions  $y(x_1) = y_1$ ,  $y(x_2) = y_2$  we get

$$y(x) = \frac{y_2 - y_1}{x_2 - x_1}(x - x_1) + y_1. \quad (15)$$

- (b) Equilibrium under gravity is reached when the  $y$  coordinate of the center of mass is at its lowest point. Thus, we must minimize the functional

$$L = \int y ds = \int_{x_1}^{x_2} y \sqrt{1 + y'^2} dx. \quad (16)$$

Once again, we use Eq. 12 with  $f(y, y') = y\sqrt{1 + y'^2}$  and get

$$\frac{yy'^2}{\sqrt{1 + y'^2}} - y\sqrt{1 + y'^2} = -\frac{y}{\sqrt{1 + y'^2}} = C_1 = \text{const}, \quad (17)$$

or,

$$\frac{y^2}{C_1^2} - y'^2 = 1. \quad (18)$$

This is a separable differential equation whose general solution is

$$y(x) = C_1 \cosh\left(\frac{x}{C_1} + C_2\right) \quad (19)$$

The constants  $C_1$  and  $C_2$  are determined by the boundary conditions  $y(x_1) = y_1$ ,  $y(x_2) = y_2$ .

**P2.3** [2 + 3 = 5 points]

If a conservative force field in 3-dimensional space is axially symmetric, and we choose the  $z$  axis of a cylindrical coordinate system  $\{r, \phi, z\}$  along the axis of symmetry, then show that

- (a) The corresponding potential has the form  $U = U(r, z)$ ,  
 (b) The force always lies in a plane containing the  $z$  axis.

**S2.3** (a) A general conservative force in 3 dimensions can be expressed as

$$\nabla U(r, \phi, z) = \frac{\partial U}{\partial r} \hat{\mathbf{e}}_r + \frac{1}{r} \frac{\partial U}{\partial \phi} \hat{\mathbf{e}}_\phi + \frac{\partial U}{\partial z} \hat{\mathbf{e}}_z \quad (20)$$

Since the force is axially symmetric, the coefficient of  $\hat{\mathbf{e}}_\phi$  must be zero:

$$\frac{\partial U}{\partial \phi} = 0 \quad \implies \quad U = U(r, z). \quad (21)$$

- (b) The force lies in the plane spanned by the unit vectors  $\hat{\mathbf{e}}_r$  and  $\hat{\mathbf{e}}_z$ , which always contains the  $z$  axis.

**P2.4** [3 + 2 = 5 points]

In Problem P1.4 (see Homework assignment 1), is

$$H \equiv \dot{\theta} \frac{\partial L}{\partial \dot{\theta}} - L \quad (22)$$

conserved? Evaluate  $H$  and compare it to the energy  $E \equiv T + V$ .

**S2.4** Recall the Lagrangian

$$L = \frac{m}{2}(R^2\dot{\theta}^2 + R^2\Omega^2 \sin^2 \theta) + mgR \cos \theta \quad (23)$$

and note that it is time-independent. Hence,  $H$  is conserved. Direct calculation yields

$$H(\theta, p_\theta) = \frac{m}{2} \left( \frac{p^2}{mR^2} - R^2\Omega^2 \sin^2 \theta \right) - mgR \cos \theta. \quad (24)$$

This is not the energy. The energy  $E$  is

$$E \equiv T + V = \frac{m}{2}(R^2\dot{\theta}^2 + R^2\Omega^2 \sin^2 \theta) - mgR \cos \theta = H + mR^2\Omega^2 \sin^2 \theta. \quad (25)$$

Since  $H$  is constant in time, but  $\theta$  may vary,  $E$  is not constant in time. Energy must enter and leave the system to keep the hoop rotating at a constant speed, and the term  $mR^2\Omega^2 \sin^2 \theta$  represents this varying amount of energy.

**P2.5** [3 + 6 + 2 = 11 points]

The Lagrangian for a particle of mass  $m$  and electric charge  $e$  moving in the  $xy$  plane is given by

$$L = \frac{m}{2}(\dot{x}^2 + \dot{y}^2) + eEy - \frac{eB}{c}y\dot{x}.$$

This describes the motion of the particle in a uniform electric field  $E$  in the  $y$  direction and a uniform magnetic field  $B$  in the  $z$  direction.  $c$  is the speed of light.

- Write down the Euler-Lagrange equations.
- Find the Hamiltonian. Simplify the expression (eliminating  $\dot{x}$  by making use of a first integral of motion). What can you say about the general allowed motions in  $y(t)$ ?
- At  $t = 0$ ,  $x = y = 0$ . What critical value  $v_c$  must  $\dot{x}(0)$  take in order for the particle to go in a uniform motion? What is the corresponding value of  $\dot{y}$ ? (You can either make use of the results in part (a) or that in part (b).)

**S2.5** (a) The Euler-Lagrange equations are

$$\frac{\partial L}{\partial x} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} = 0 \quad \Rightarrow \quad \frac{d}{dt} \left( m\dot{x} - \frac{eB}{c}y \right) = 0 \quad \Rightarrow \quad \ddot{x} = \frac{eB}{mc}\dot{y} \quad (26)$$

and

$$\frac{\partial L}{\partial y} - \frac{d}{dt} \frac{\partial L}{\partial \dot{y}} = 0 \quad \Rightarrow \quad \frac{d}{dt}(m\dot{y}) = eE - \frac{eB}{c}\dot{x} \quad \Rightarrow \quad \ddot{y} = -\frac{eB}{mc}\dot{x} + \frac{eE}{m}. \quad (27)$$

(b)

$$p_x = \frac{\partial L}{\partial \dot{x}} = m\dot{x} - \frac{eB}{c}y \quad \Rightarrow \quad \dot{x} = \frac{p_x}{m} + \frac{eB}{mc}y. \quad (28)$$

$$p_y = \frac{\partial L}{\partial \dot{y}} = m\dot{y} \quad \Rightarrow \quad \dot{y} = \frac{p_y}{m}. \quad (29)$$

Hence, the Hamiltonian is

$$\begin{aligned} H &= \dot{x}p_x + \dot{y}p_y - L \\ &= \frac{p_x^2}{m} + \frac{eB}{mc}p_x y + \frac{p_y^2}{m} - \frac{m}{2} \left[ \left( \frac{p_x}{m} + \frac{eB}{mc}y \right)^2 + \left( \frac{p_y}{m} \right)^2 \right] \\ &\quad - eEy + \frac{eB}{c}y \left( \frac{p_x}{m} + \frac{eB}{mc}y \right) \\ &= \frac{1}{2m}(p_x^2 + p_y^2) + \left( \frac{eB}{mc}p_x - eE \right) y + \frac{e^2 B^2}{2mc^2} y^2. \end{aligned} \quad (30)$$

Since  $x$  does not explicitly appear in the Hamiltonian, the corresponding momentum  $p_x$  must be conserved:

$$\dot{p}_x = -\frac{\partial H}{\partial x} = 0 \quad \Rightarrow \quad p_x = \text{const.} \quad (31)$$

Using Eqs. 28 and 29 in another Hamilton's equation,

$$\ddot{y} = \frac{\dot{p}_y}{m} = -\frac{1}{m} \frac{\partial H}{\partial y} = -\frac{eB}{m^2 c} p_x + \frac{eE}{m} - \frac{e^2 B^2}{m^2 c^2} y \quad (32)$$

or,

$$\ddot{y} + \frac{e^2 B^2}{m^2 c^2} y = \frac{e}{m} \left( E - \frac{p_x}{mc} B \right) = \text{const.} \quad (33)$$

Equation 33 represents a simple harmonic oscillator along  $y$ , with a characteristic frequency  $\omega_0 = \frac{eB}{mc}$ , acted on by a constant driving force. Note that the same equation could be obtained by integrating both sides of Eq. 26 once with respect to  $t$  to get  $\dot{x}$  and substituting it in Eq. 27.

(c) Uniform motion  $\Rightarrow \quad \ddot{x} = 0; \ddot{y} = 0$ . By Eqs. 27 and 26, this means

$$v_c = \dot{x} = \frac{cE}{B} \quad \text{and} \quad \dot{y} = 0. \quad (34)$$