

QCD and Electroweak Physics at Hadron Colliders

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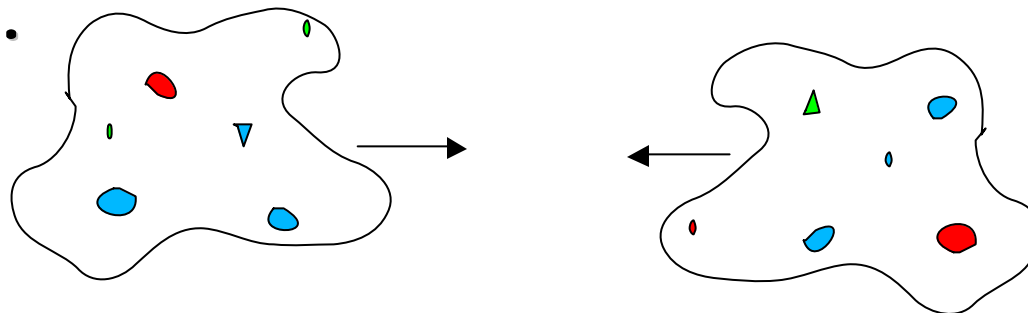


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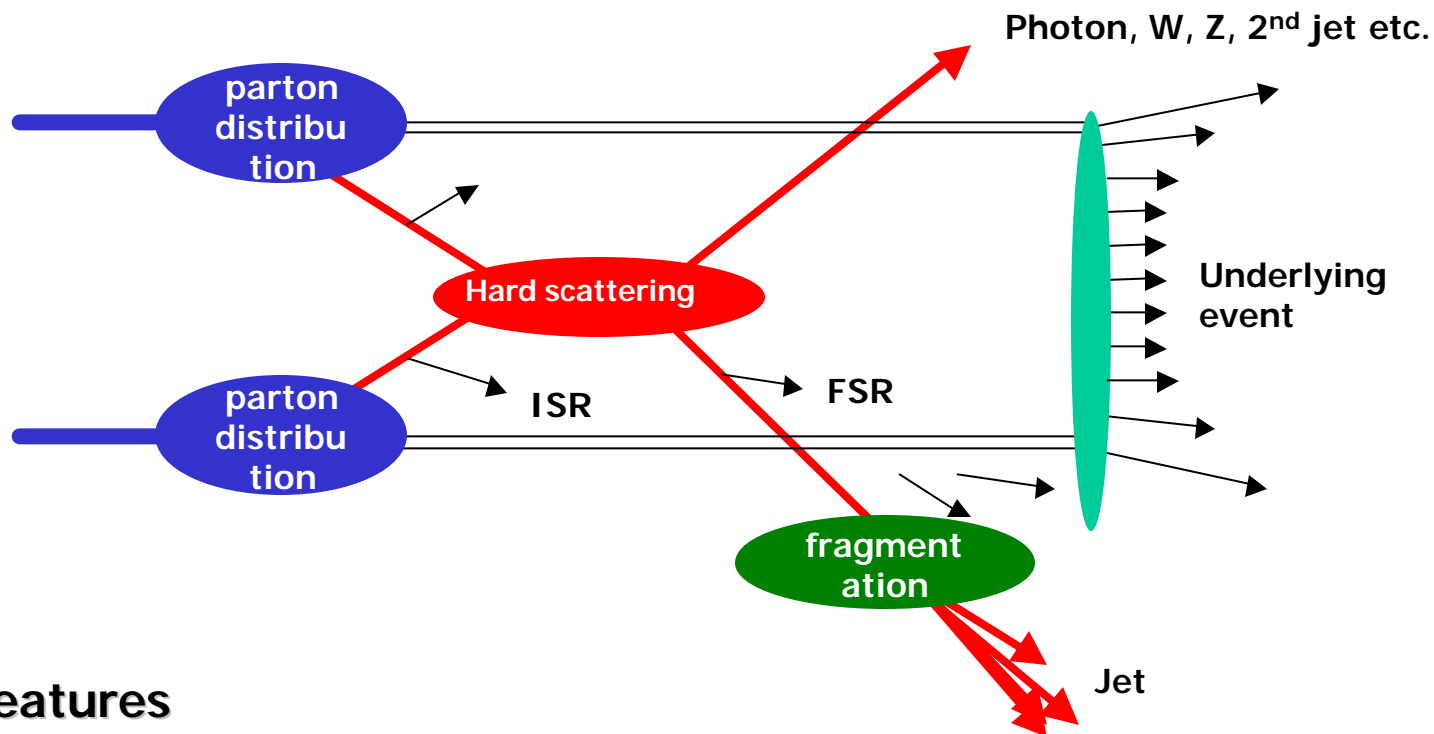
The Hadron Collider Idea

- Similar to a collision between two bags of rocks.



- Occasionally two stones or “partons” will collide.
- But you’ll never know which two collide and how much energy they carry.
- For the proton the “partons” are quarks and gluons.

A Bit More Formal

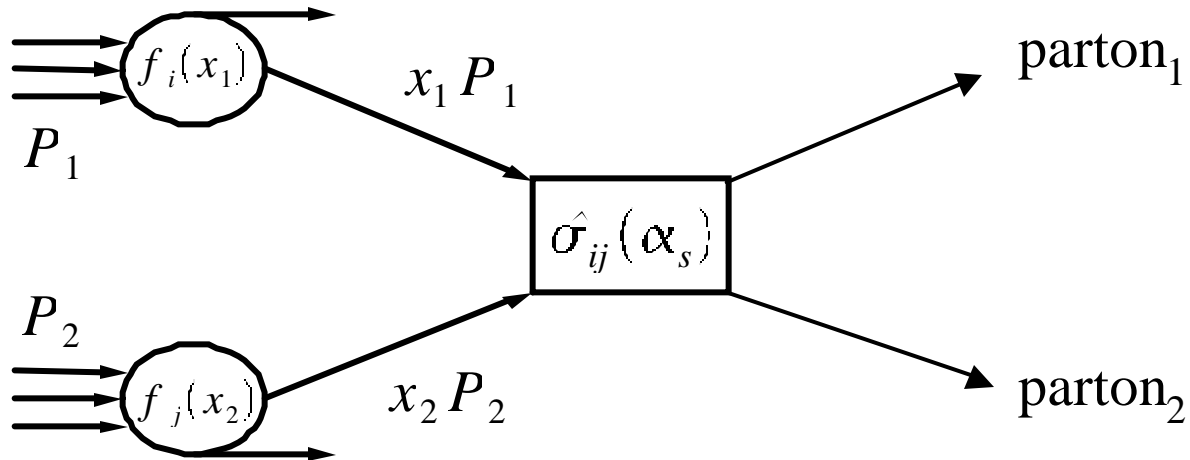


- **Features**

- The initial states are composed of quark and gluon distributions (absent for electron-electron machines)
- As with Rutherford scattering the hard scatter can be described quite well
- Particles can radiate from incoming and outgoing partons
- The un-scattered partons form an “underlying” event (also absent for electron-electron machines)

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Very Formal



$$S = \sum_{ij} \int dx_1 dx_2 f_i(x_1, m_F^2) f_j(x_2, m_F^2) \hat{\sigma}_{ij}(\alpha_s^m(m_R^2), x_1 P_1, x_2 P_2, \frac{Q^2}{m_F^2}, \frac{Q^2}{m_R^2})$$

Sum over
initial states

Parton
Distributions

Factorization
Scale

Point Cross
Section

Order α_s^m

Renormalization
Scale

QCD: Quantum ChromoDynamics

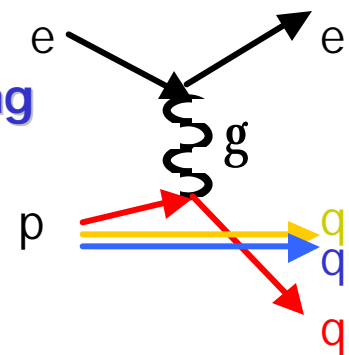
- A gauge theory (like electromagnetism)
 - describes fermions (quarks) which carry an SU(3) charge (color)
 - interact through the exchange of vector bosons (gluons)
- Interesting features:
 - gluons themselves have color
 - interactions are strong
 - coupling constant between quarks and gluons runs rapidly
 - weak at momentum transfers above a few GeV

$$\alpha_s(Q^2) = \frac{12\pi}{(33 - 2n_f) \ln Q^2 / \Lambda^2}$$

- Q^2 is the momentum transfer between the partons related to Mandelstam variable $t^2 = (p_1 - p_3)^2$

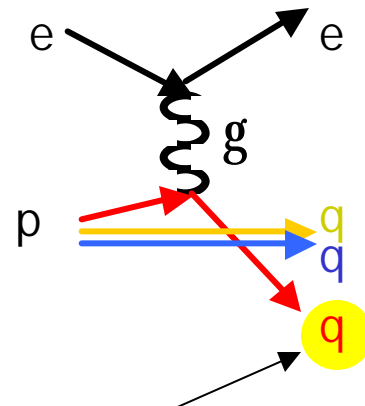
A Bit of Quark History

- These features lead to a picture where quarks and gluons are bound inside hadrons if left to themselves
- Behave like “free” particles if probed at high momentum transfer: asymptotic freedom (This year’s noble!)
- This is exactly what was seen in deep inelastic scattering experiments in the late 1960’s which led to the genesis of QCD
 - electron beam scattered off nucleons in a target
 - electron scattered from point-like constituents inside the nucleon
 - $\sim 1/\sin^4(\theta/2)$ behaviour like Rutherford scattering
 - other (spectator) quarks do not participate



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One More Ingredient: Fragmentation

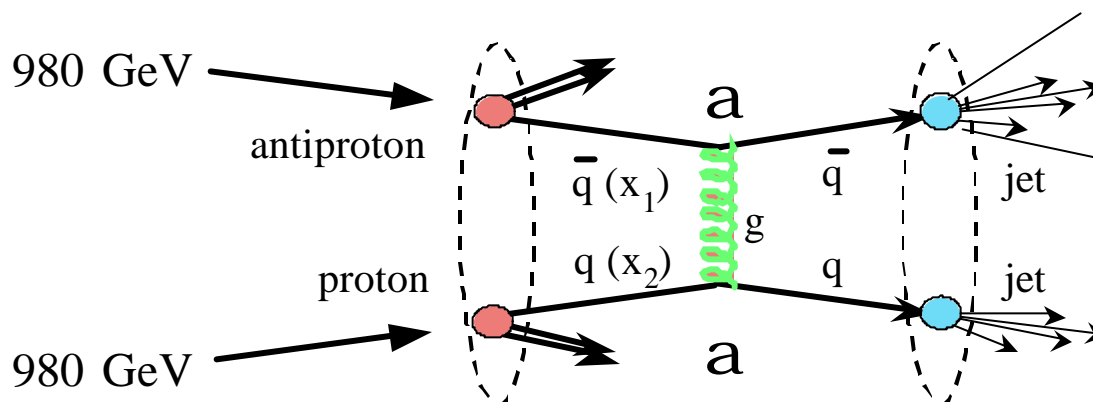


So what happens to a quark that is knocked out of the proton?

α_s is large (large distances = low momentum scales)

- Strong field = lots of gluon radiation between scattered quark and color-ed remnant of the nucleon
- Energy required for separation is large compared to pair production energy.
- Quarks and gluons produced in the “wake” of the outgoing quark recombine to form a “spray” of roughly collinear, colorless hadrons: a jet
- This process is known as: “fragmentation” or “hadronization”

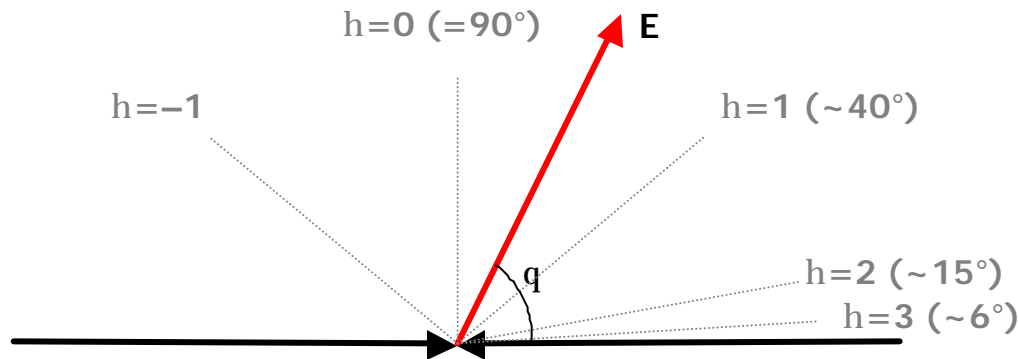
A Descriptive Example: Jet Production



- Parton distribution functions determine incoming energy
Collider energy (center of mass) = 1.96 TeV
Parton energy (center of mass) = $x_1 x_2 s$
- Probability of gluon exchange dictated by full theory of QCD.
- Strength of parton scatter given by coupling constant, momentum transfer
- Final state requires fragmentation.

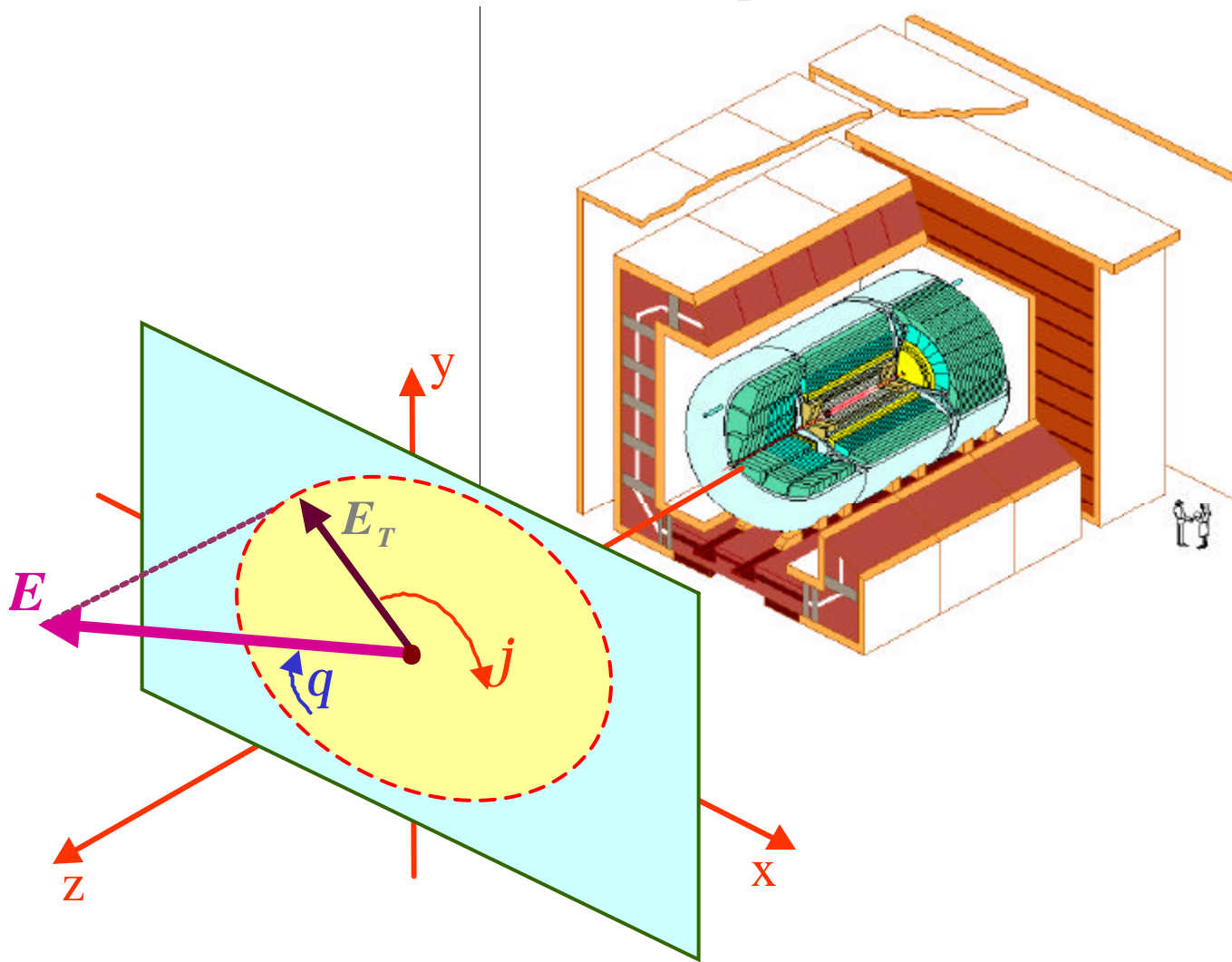
Hadron Collider Variables

- The incoming parton momenta x_1 and x_2 are unknown, and usually the beam particle remnants escape down the beam pipe
 - longitudinal motion of the centre of mass cannot be reconstructed



- Focus on transverse variables
 - Transverse Energy $E_T = E \sin q$ ($= p_T$ if mass = 0)
- And Longitudinally boost-invariant quantities
 - Pseudorapidity $h = -\log(\tan q/2)$ particle production typically scales per unit rapidity

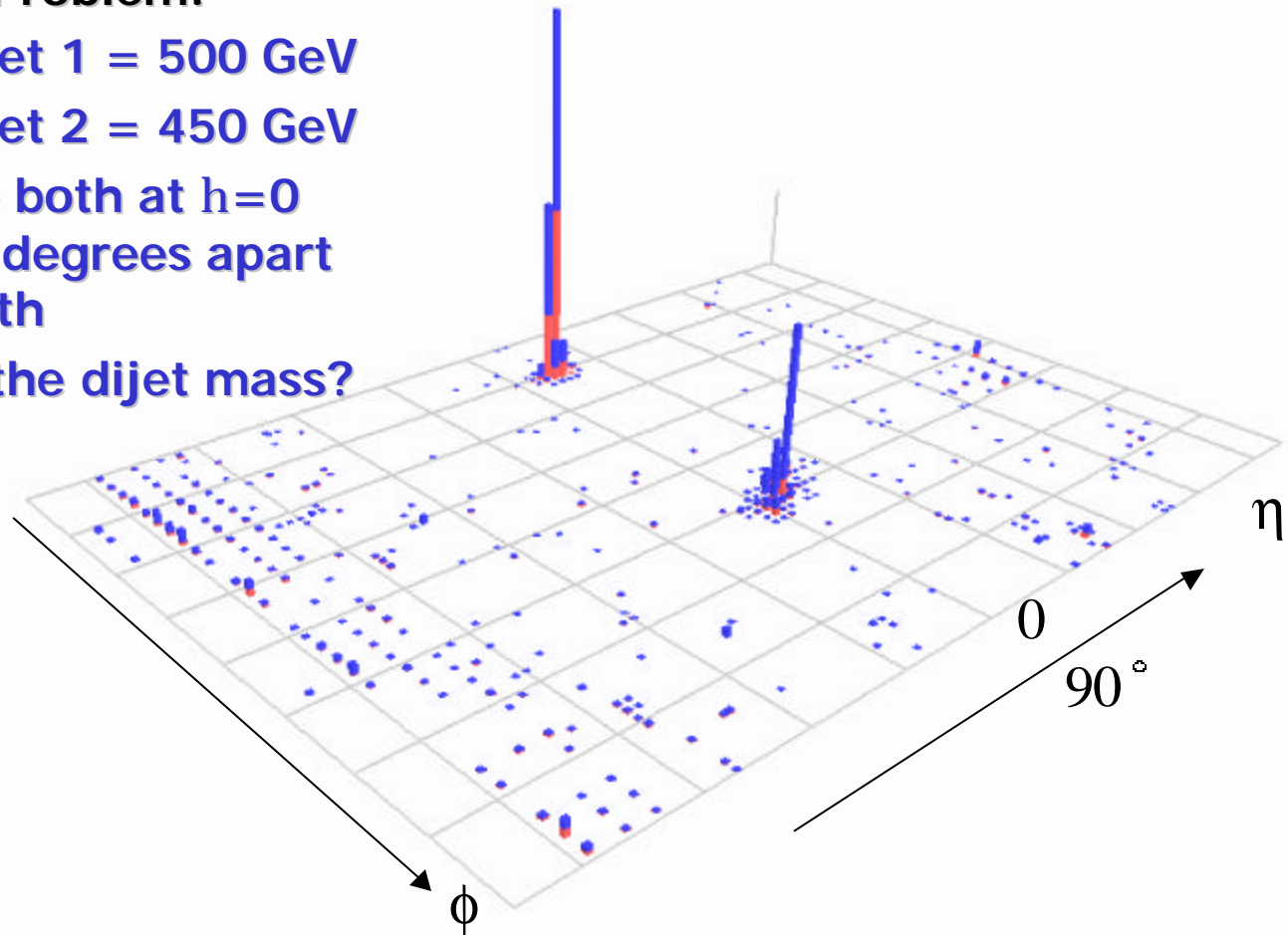
An Example:



A Real Event ! (see Slide 8)

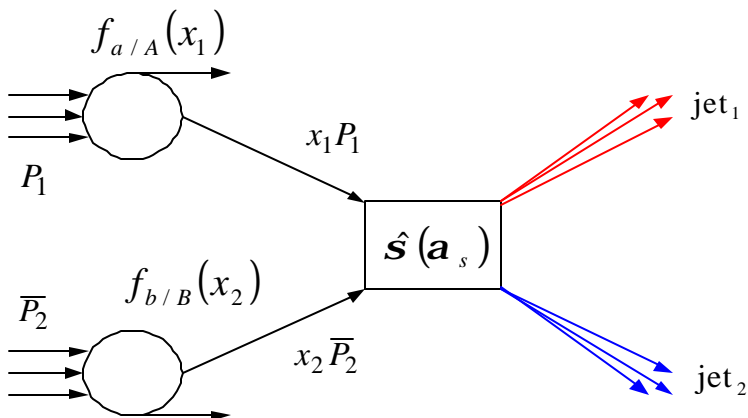
6. Homework Problem:

- Energy Jet 1 = 500 GeV
- Energy Jet 2 = 450 GeV
- They are both at $h=0$ and 180 degrees apart in azimuth
- What is the dijet mass?



Counting (q and g) Jets: Inclusive Jet Cross Section as a Test of the Standard Model (QCD)

Single Inclusive Jets: $pp \rightarrow jet + X$



$$s(p_1 \bar{p}_2 \rightarrow 2 \text{ jets}) = \int dx_1 dx_2 f_{a/A}(x_1) f_{b/B}(x_2) \hat{s}(ab \rightarrow cd)$$

$$\frac{1}{DE_T Dh} \frac{d^2 S}{dE_T dh} \stackrel{\sim}{=} \frac{N_{jet}}{DE_T Dhe} \stackrel{\sim}{=} \frac{1}{L dt} \text{ vs. } E_T$$

DE_T $\hat{=}$ E_T bin size

e $\hat{=}$ selection efficiency

Dh $\hat{=}$ h bin size

L $\hat{=}$ inst. Luminosity

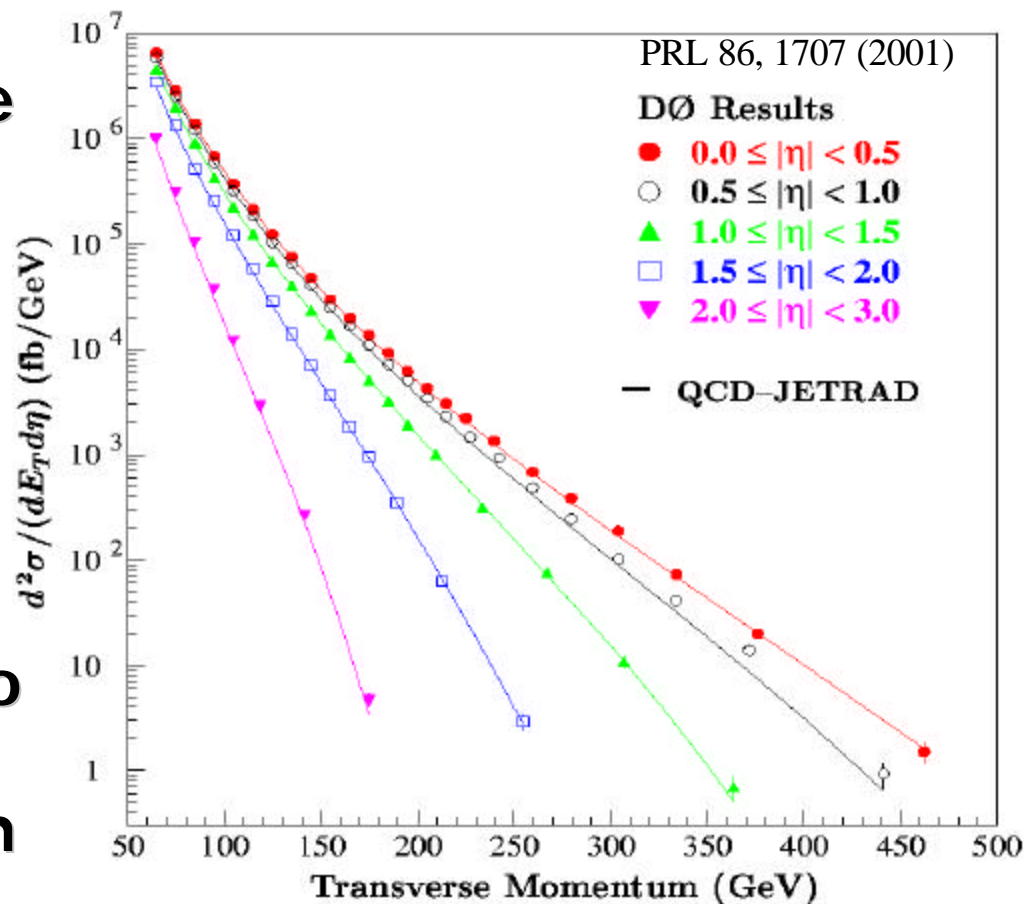
N_{jet} $\hat{=}$ # of jets in the bin

More Homework

7. Calculate the central jet cross section for 900 jets observed with 90% efficiency in the range $295 < ET < 305$ GeV and $-0.5 < h < 0.5$ and produced with a luminosity of 0.3 fb^{-1}
8. What would be the statistical error?

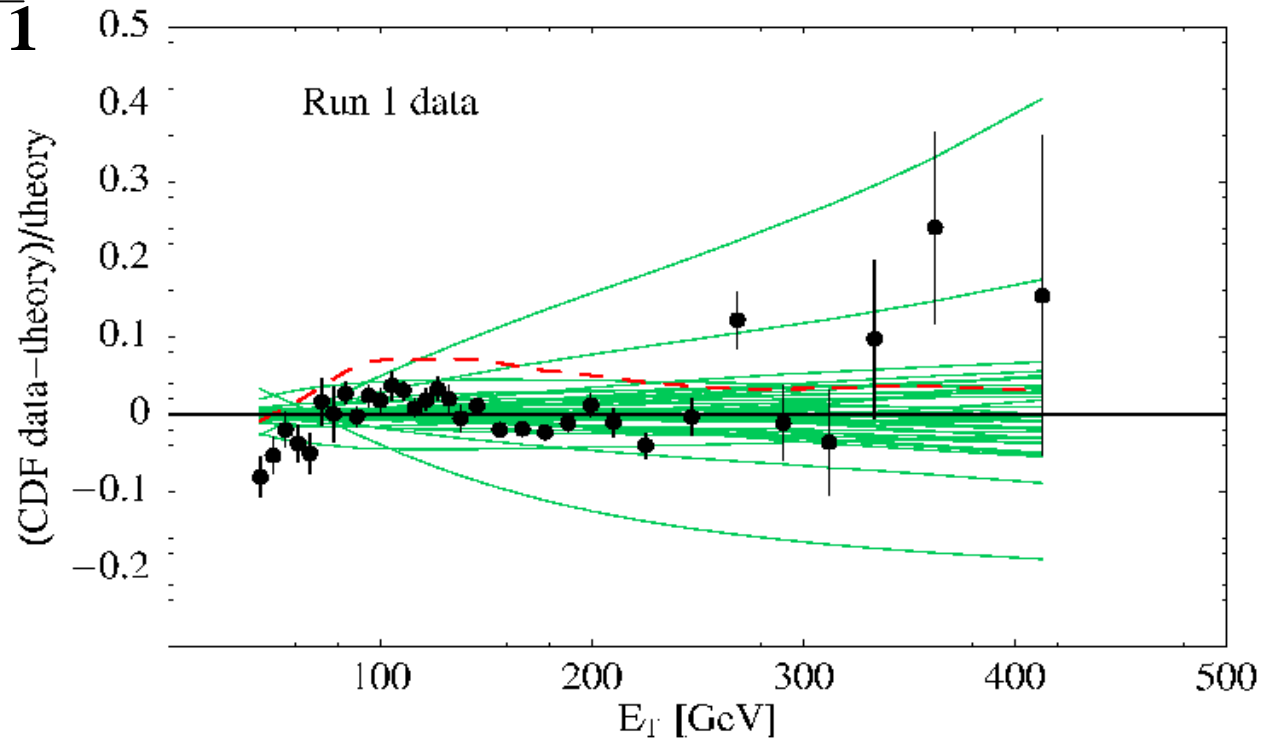
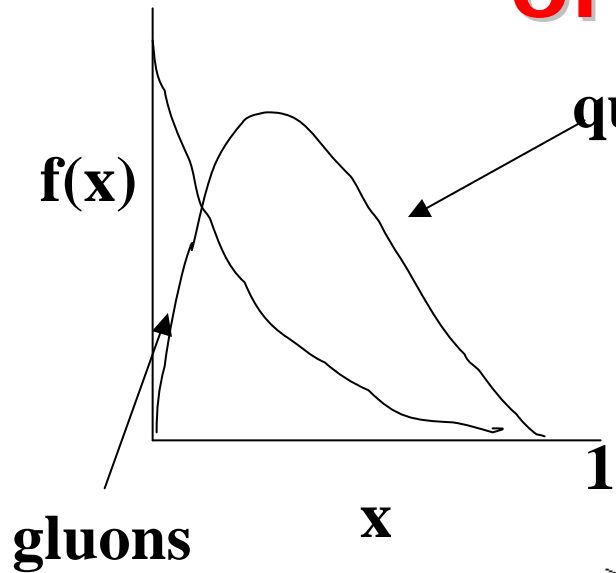
Inclusive Jet Cross Section

- ? Probes the hard interaction vertex over many decades in momentum exchange
- ? Probes for deviations from pQCD at small distance scales
- ? Sensitive to pdfs and running of a_s
- ? Cross section falls by seven orders of magnitude from 50 to 450 GeV
- ? Good agreement with QCD over the whole range



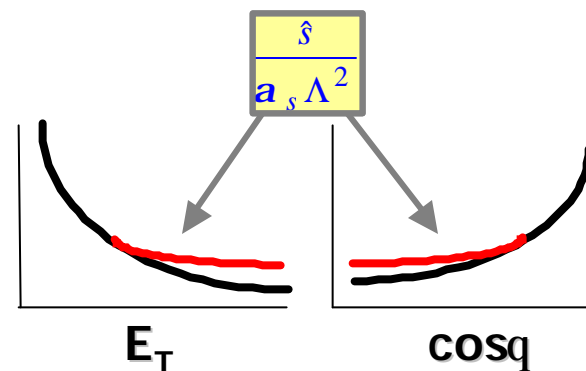
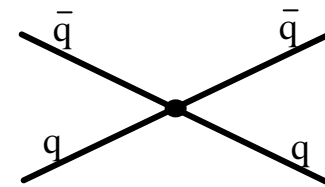
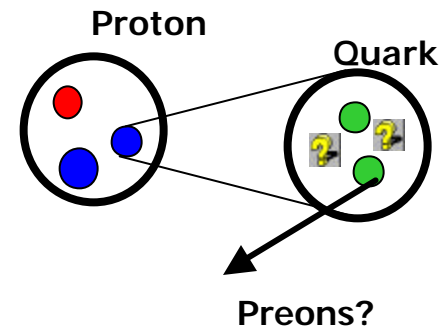
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Data Consistent with a Number of Possible PDFs

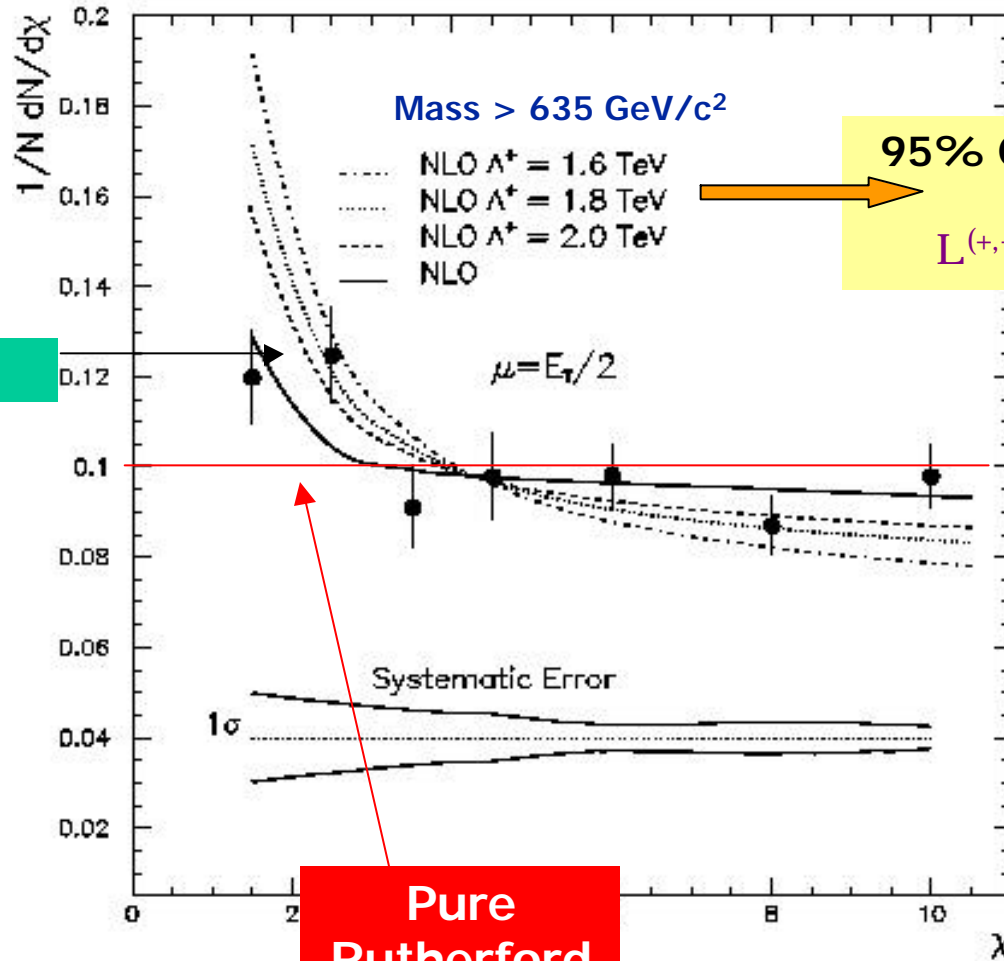


Jet Cross Sections Probe New Physics

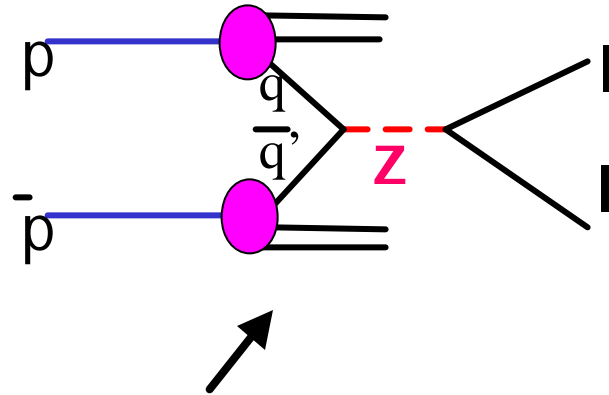
- If quarks contain smaller constituents
 - constituent interactions have a scale L
 - at momentum transfers $\ll L$, quarks appear pointlike and QCD is valid
 - as we approach scale L , interactions can be approximated by a four-fermion contact term:
 - at and above L , constituents interact directly



DØ Dijet Angular Distribution



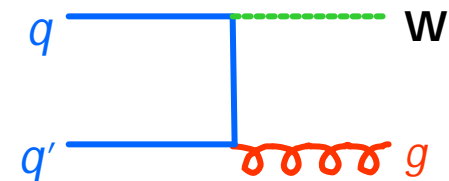
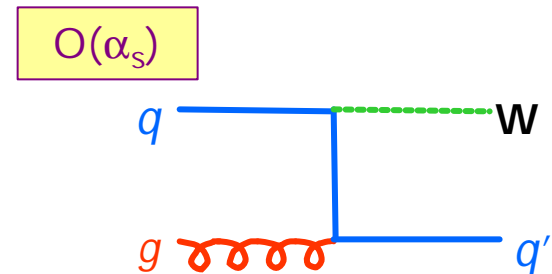
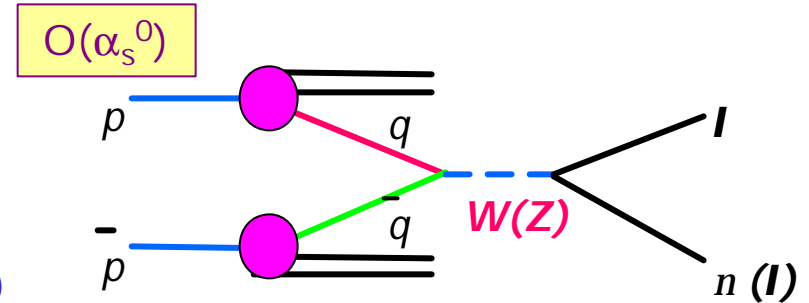
Electroweak Vector Boson Production



- Recall our introductory lecture?
- Especially interesting physics as the bosons couple to both quarks and leptons.
- W and Z boson masses test electroweak theory
- Lepton signals experimentally distinct.

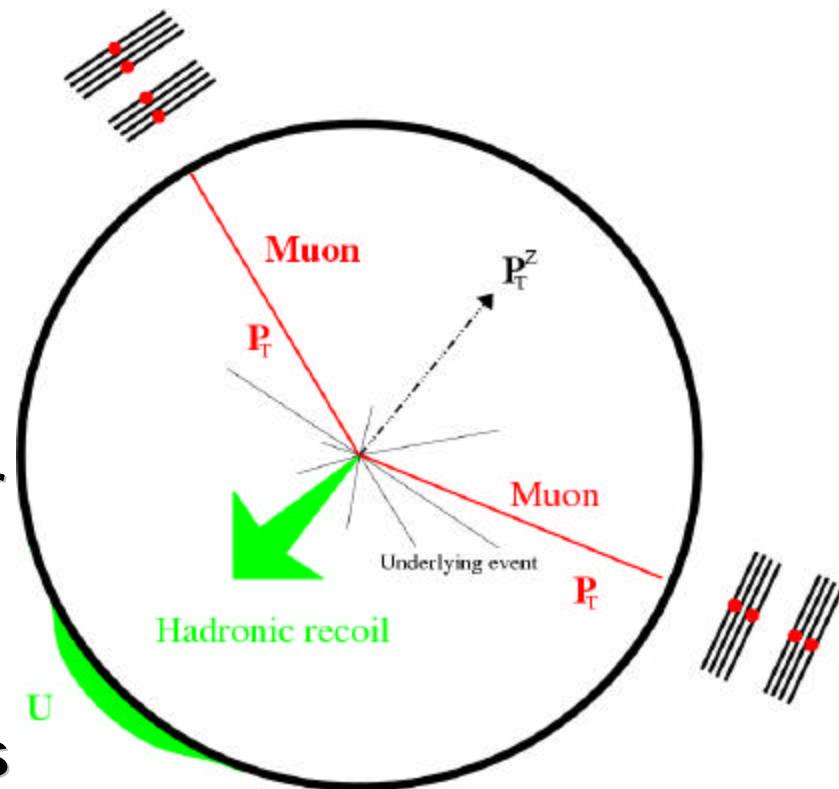
W/Z Production at Hadron Colliders

- **General Characteristics**
 - Production dominated by $q\bar{q}$ annihilation
(~60% valence-sea, ~20% sea-sea)
 - Due to very large $pp \rightarrow jj$ production, need to use leptonic decays to extract signal
- **Higher order QCD corrections:**
 - Boson produced with mean $p_T \sim 10$ GeV
 - Boson + jet events ($W + \text{jet} \sim 7\%$, $E_T^{\text{jet}} > 25$ GeV)



Experimental Signature: $Z \rightarrow m^+ m^-$

- A pair of charged leptons:
 - high p_T
 - isolated
 - opposite-charge
- Redundancy in trigger and offline selection
- Low backgrounds
- Control of systematics



Select Events, Calculate Dilepton Mass, Calculate Cross Sections

- Event selection:
 - Two central tracks:
 - μ -id with $p_T > 15$ GeV
 - opposite charge
 - $|\eta| < 1.8$
 - More than 1 isolated μ

- Calculate Cross Section:

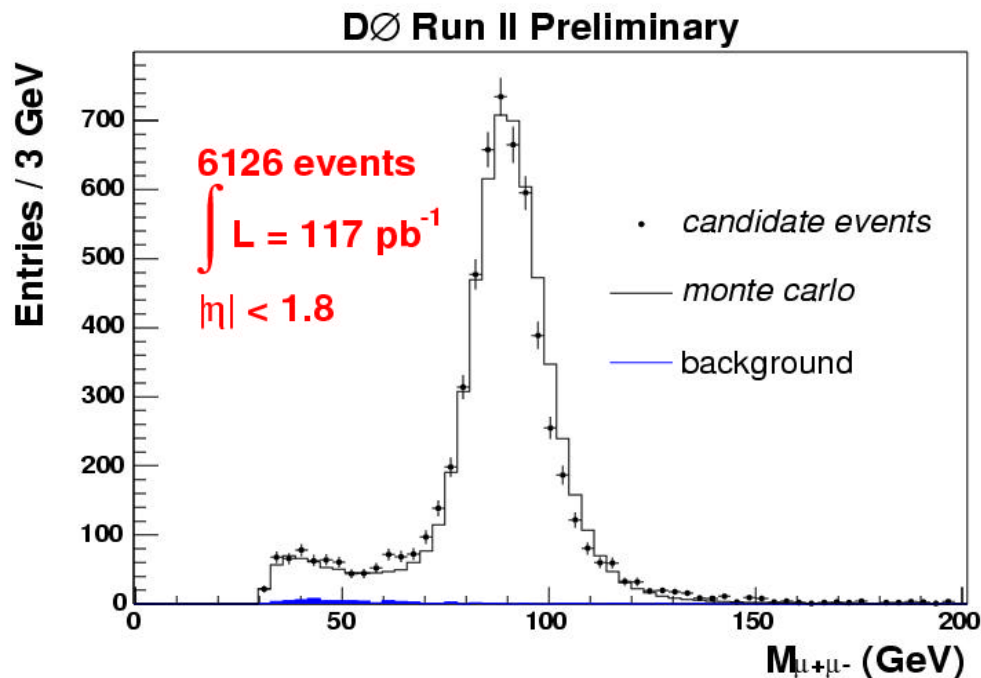
- $\epsilon_{\text{total}} = 19\%$

- For $M_{\mu\mu} > 30$ GeV
 $N_{\text{cand}} = 6126$

- $\mathcal{L} = 117 \text{ pb}^{-1}$

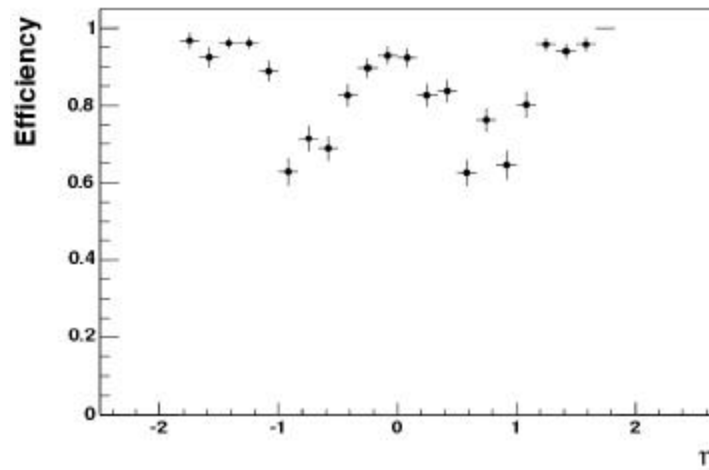
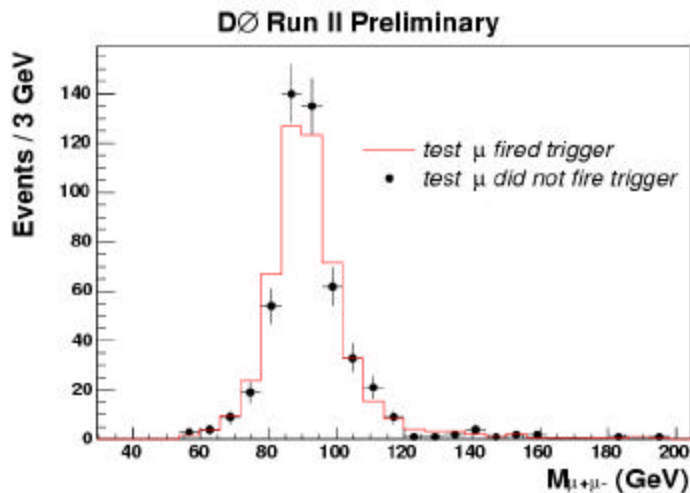
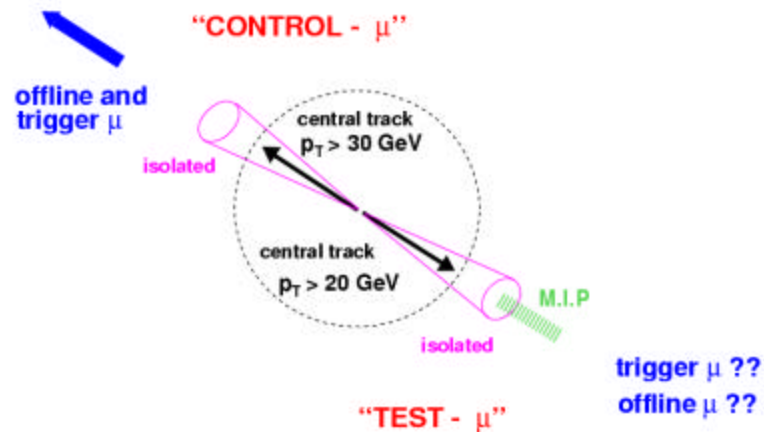
- Homework 9:

- Calculate the $Z \rightarrow \mu\mu$ cross section for yourself!
- From the figure can you tell why it's a few percent higher than the cross section on the next page?



Measuring efficiencies using the Z? $\mu^+\mu^-$ data

- There are two μ 's
- The backgrounds are low
- Can select pure Z sample with even looser cuts on one μ

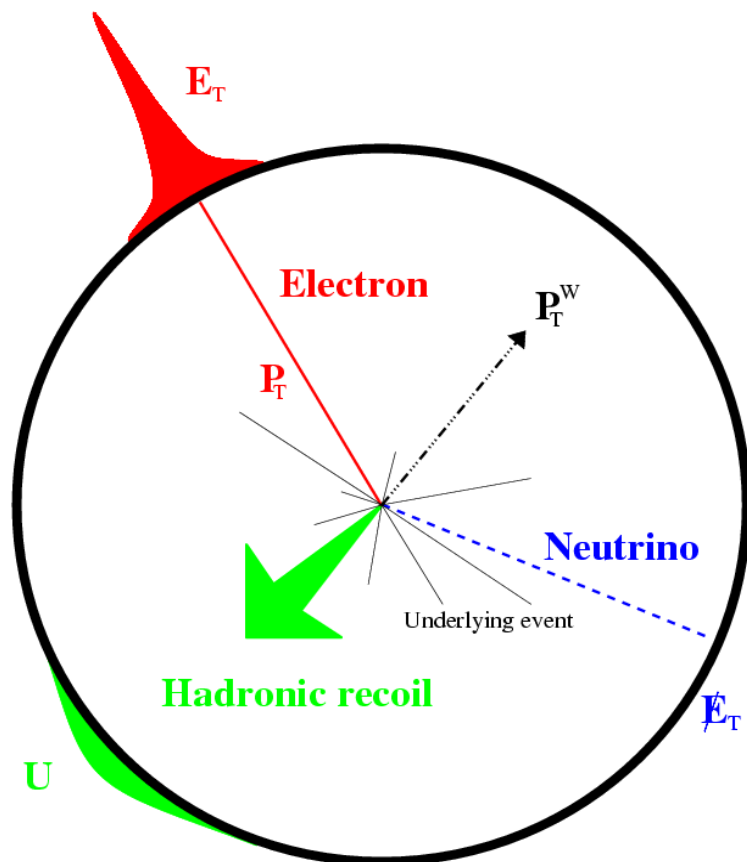


$$s_Z \cdot \text{Br}(Z? \mu^+\mu^-) = 261.8 \pm 5.0 \pm 8.9 \pm 26.2 \text{ pb}$$

stat. syst. lumi.

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Experimental Signature: $W^- \rightarrow e^- \bar{\nu}_e$

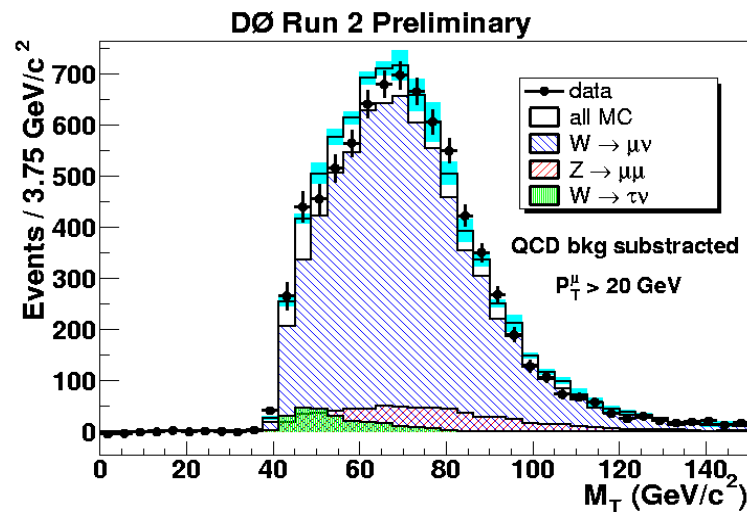
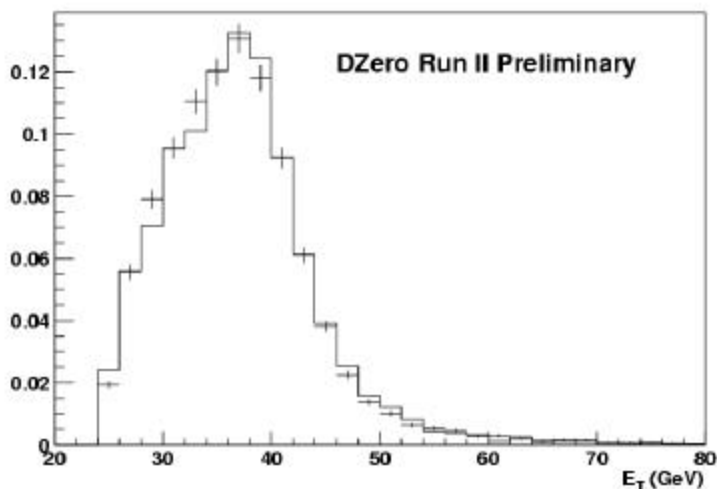


- **Single charged lepton:**
 - high p_T
 - isolated
- E_T^{miss} (from neutrino)
- Less redundancy in trigger and offline selection
- More difficult to control backgrounds and systematics
- need to understand hadronic recoil
- $\sigma \cdot \text{Br}$ 10 times larger than Z

W? en and W? $\mu\nu$

- $p_T(e) > 25 \text{ GeV}$
- $E_T^{\text{miss}} > 25 \text{ GeV}$
- $N_{\text{cand}} = 27370$
- $\mathcal{L} = 42 \text{ pb}^{-1}$

- $p_T(\mu) > 20 \text{ GeV}$
- $E_T^{\text{miss}} > 20 \text{ GeV}$
- $N_{\text{cand}} = 8302$
- $\mathcal{L} = 17 \text{ pb}^{-1}$



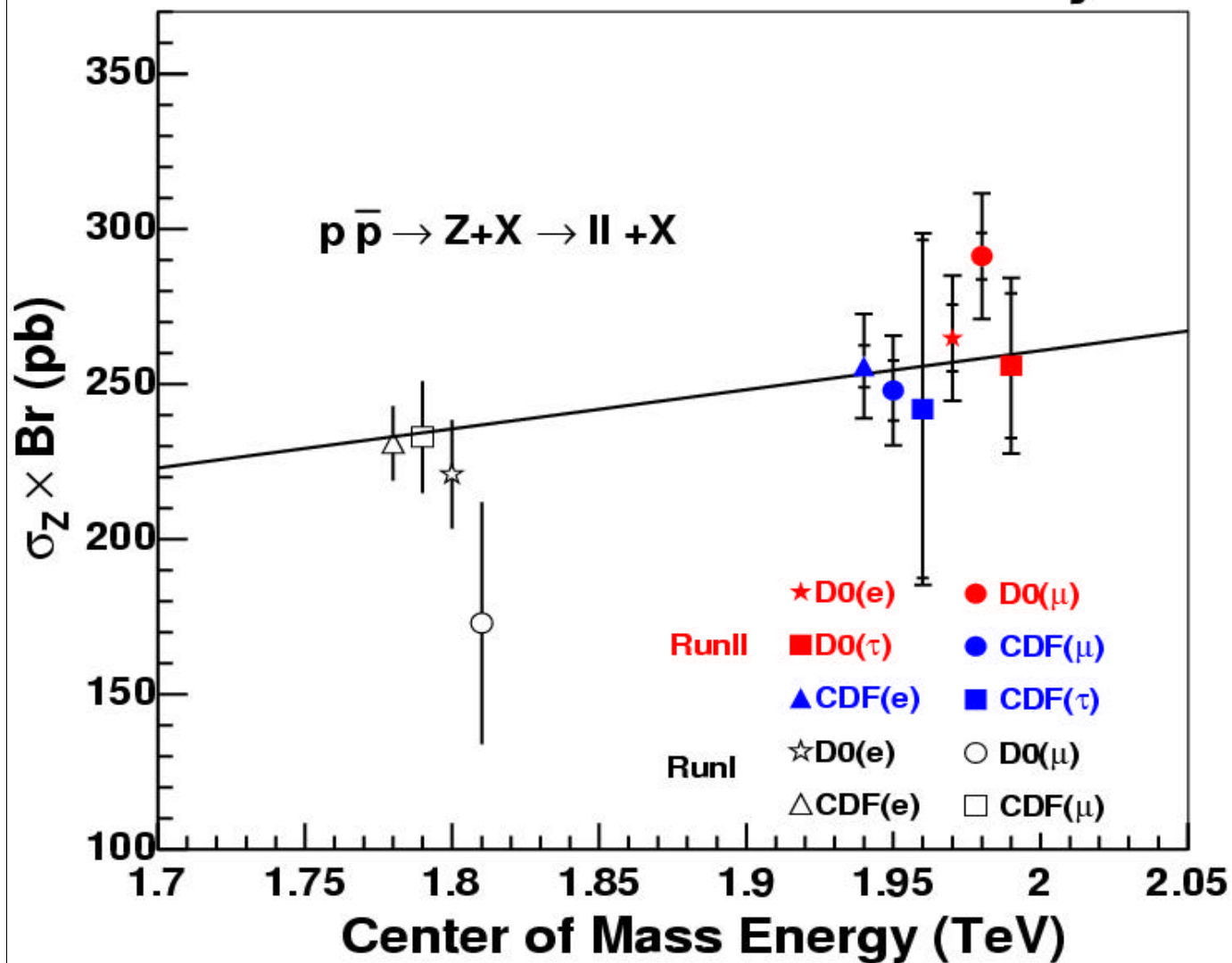
$$s_W \cdot \text{Br}(W? \text{ ev}) = 2.884 \pm 0.021 \pm 0.128 \pm 0.284 \text{ nb}$$

$$s_W \cdot \text{Br}(W? \mu\nu) = 3.226 \pm 0.128 \pm 0.100 \pm 0.322 \text{ nb}$$

stat. syst. lumi.

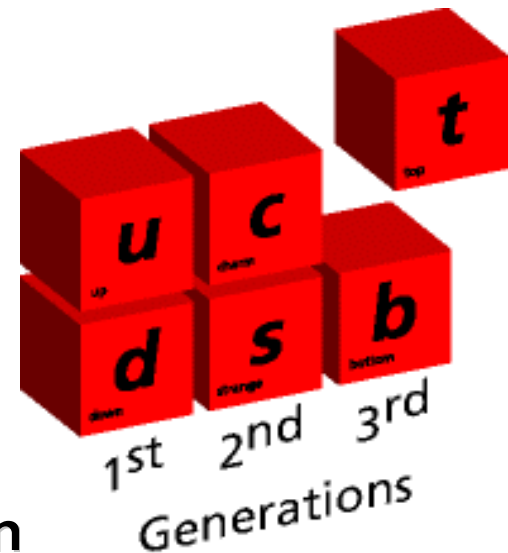
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CDF and D0 RunII Preliminary



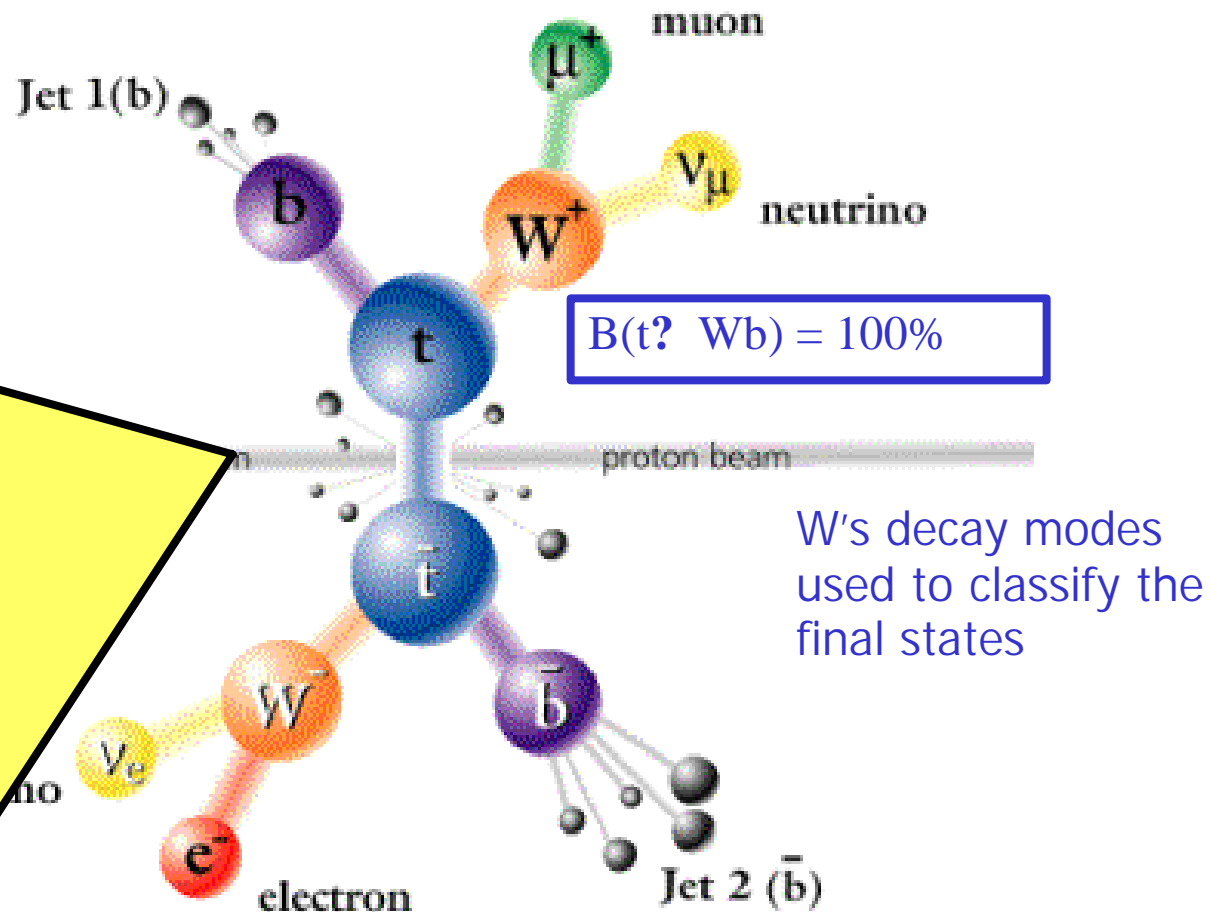
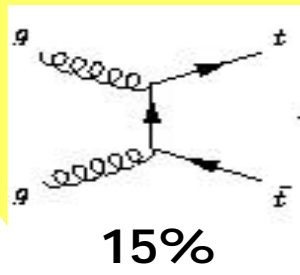
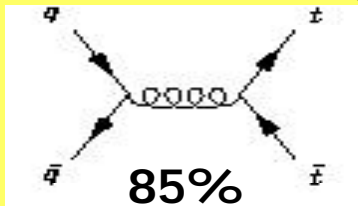
The Last Quark: an Introduction

- Top quark was expected in the Standard Model (SM) of electroweak interactions as a partner of b-quark for the third family of quarks
 - Evidence for top in 1994 (CDF)
 - Observation in 1995 (CDF&D0)
- In Run I statistical uncertainties dominated:
 - Overall consistency with the SM picture
 - but...still a few loose ends
- In anticipation of much increased statistics in Run II:
 - Rich physics menu
 - Increased luminosity → increased precision
 - Surprises?



Top Quarks at the Tevatron

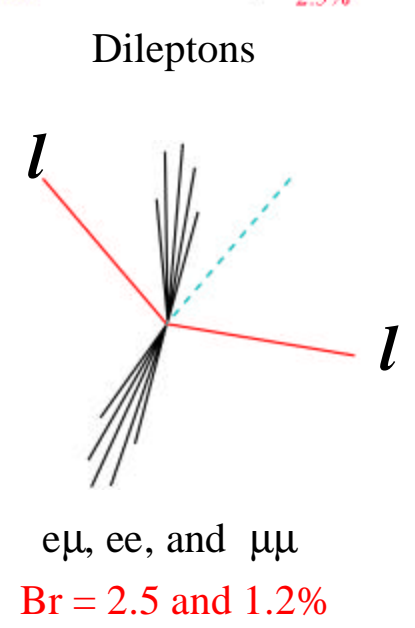
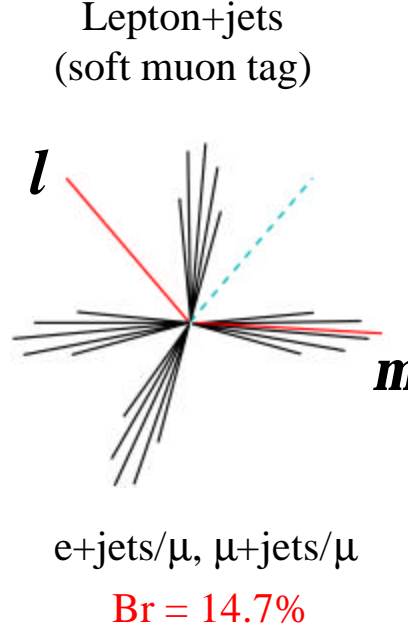
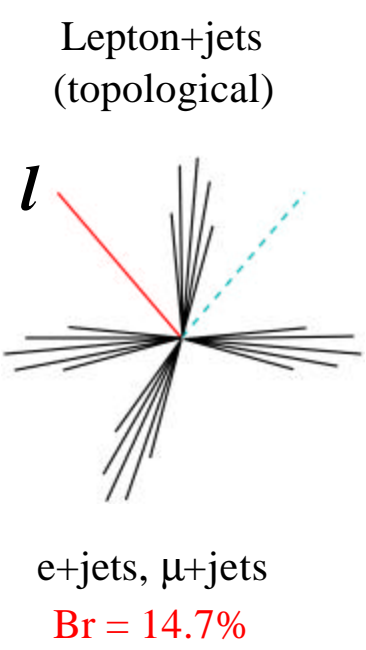
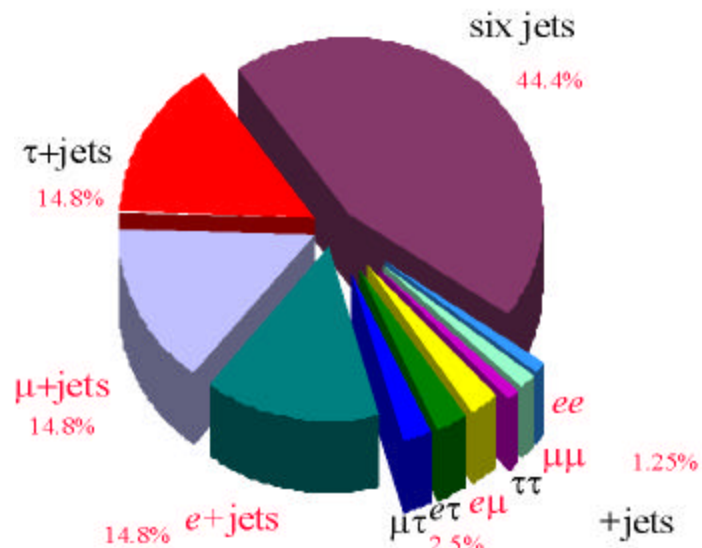
Pair production



Cross Section

- Why?
 - test of QCD predictions
 - Any discrepancy indicates possible new physics:
 - production via a high mass intermediate state
 - Non- Wb decay models
- How?
 - Measurement performed using various final states
 - Dilepton channels
 - ee, em, mm final states
 - Lepton + jets channels
 - e+jets, m+jets
 - topological analysis
 - b-tagging
 - All-jets channel
 - Use topological variables
 - exploit b-tagging using soft leptons
 - Combine them using a Neural network technique

- **W and b-quark decays specify final states**
 - Isolated high P_T leptons
 - Soft leptons in jets
 - detached vertices in jets



em + 2jet Candidate Event

Run 169920 Event 8545882 Wed Jan 22 16:09:11 2003

ET scale: 37 GeV

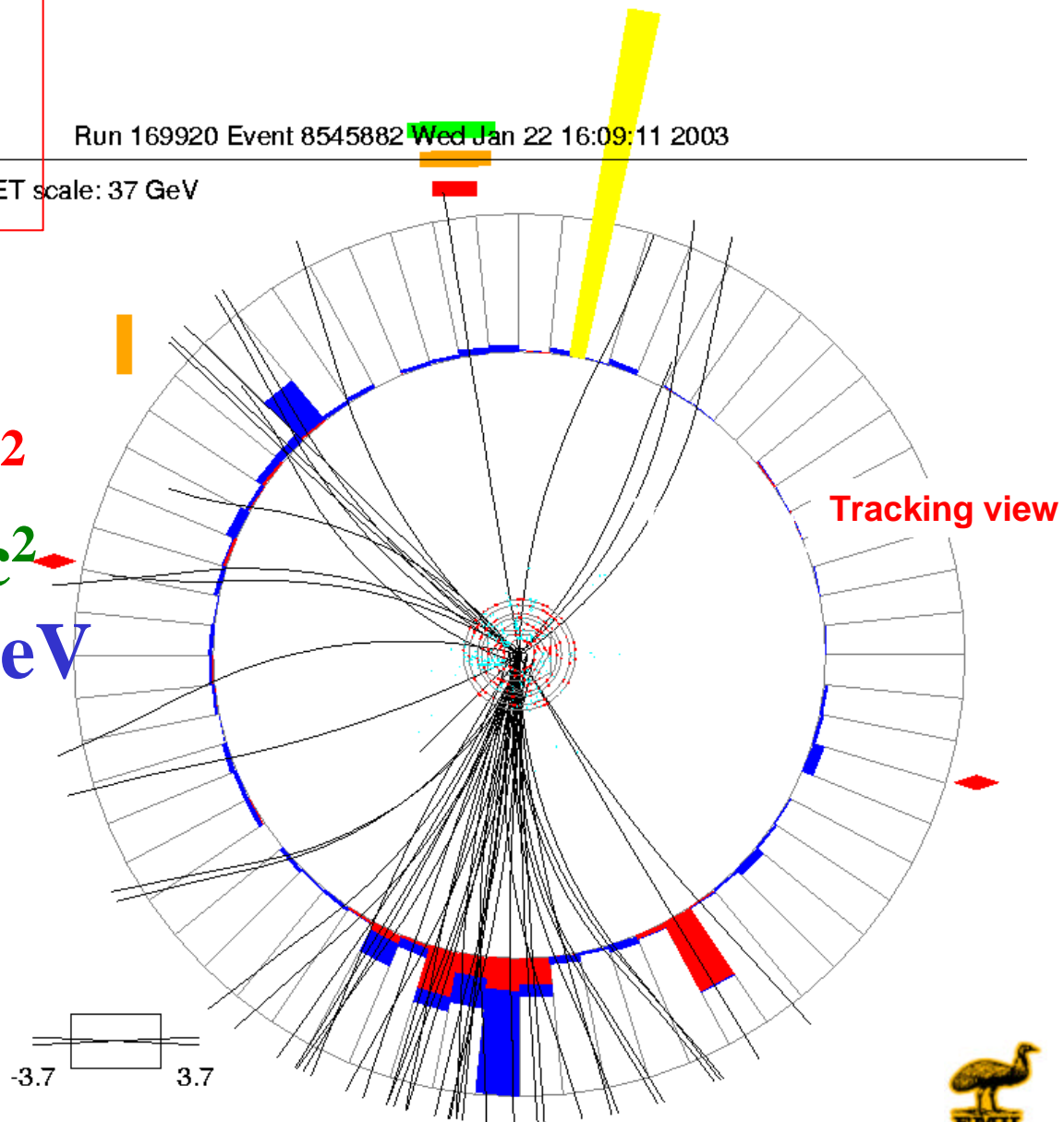
$$p_T(e) = 20.3 \text{ GeV}/c^2$$

$$p_T(m) = 58.1 \text{ GeV}/c^2$$

$$E_T^j = 141.0, 55.2 \text{ GeV}$$

$$\cancel{E}_T = 91 \text{ GeV}$$

$$H_T(e) = 216 \text{ GeV}$$



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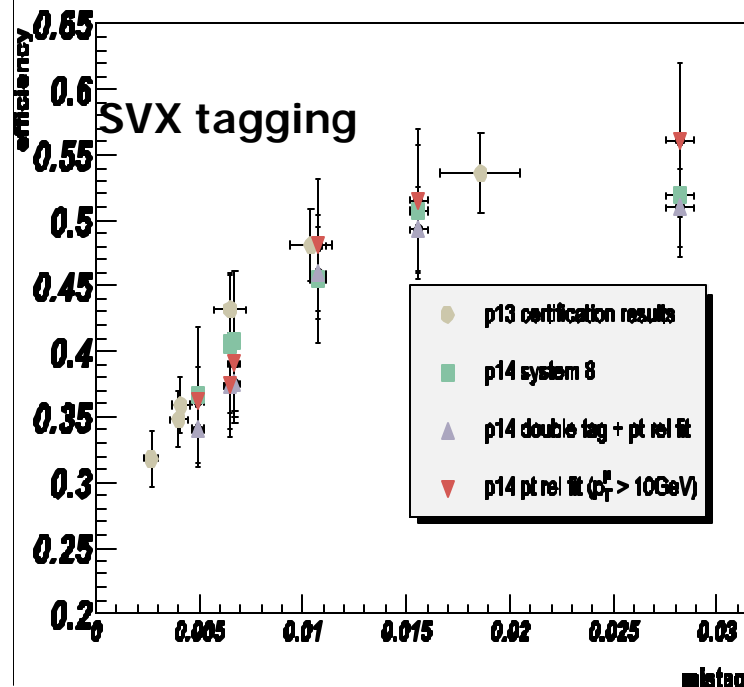
Lifetime b-tagging

Three different tagging algorithms

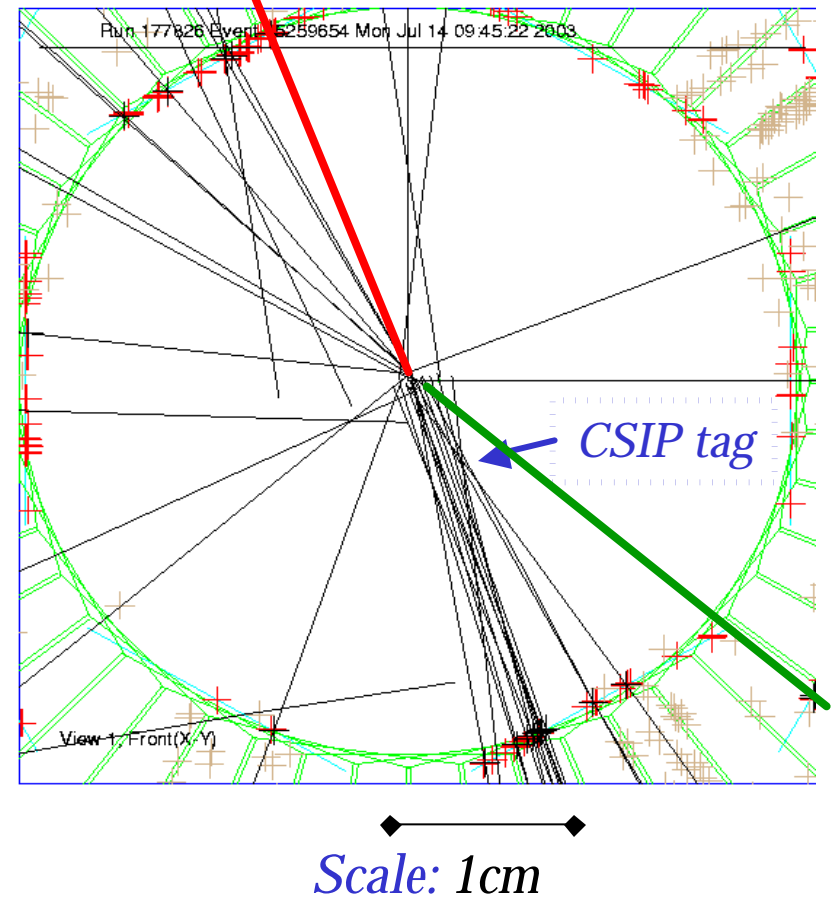
One based on secondary vertices

Two based on tracks with large IPs

Multiple methods to measure efficiencies and fake rates



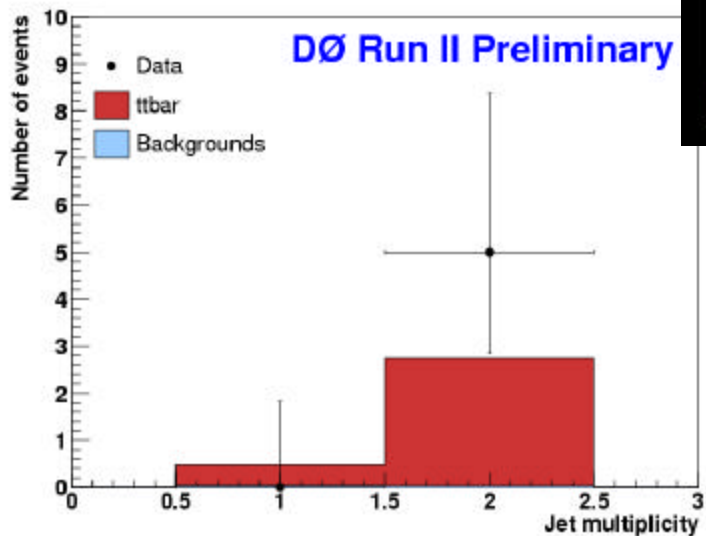
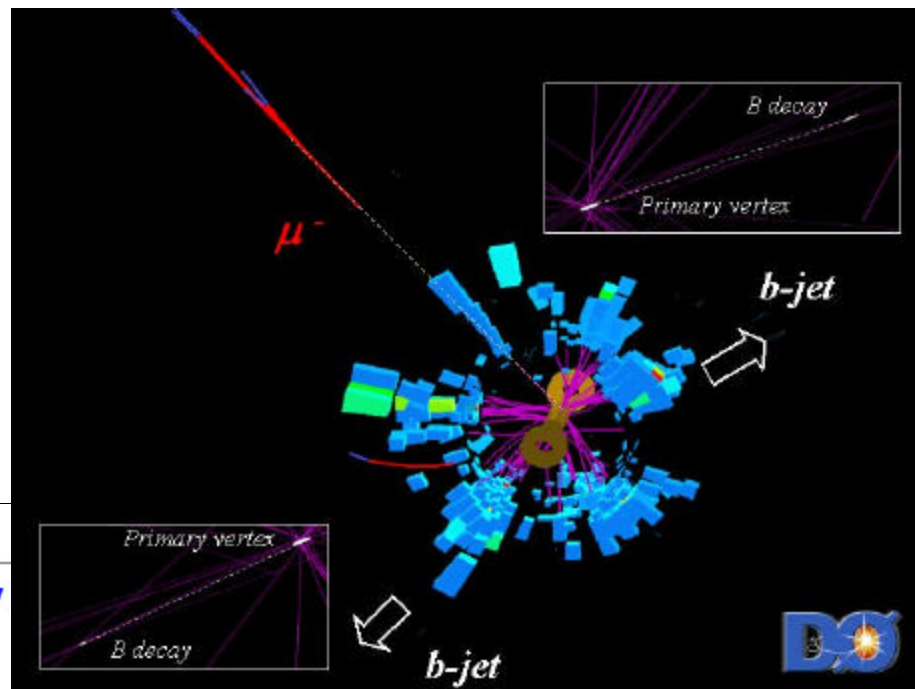
A tagged $e\mu$ top candidate



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The Golden Doubly Tagged $e\bar{e}$ Channel

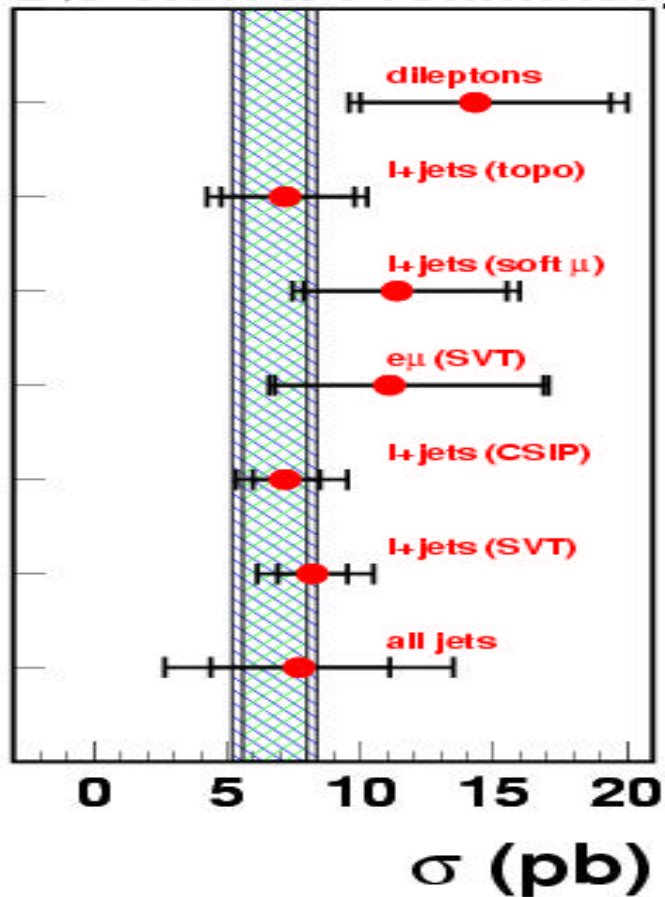
- Incredibly Clean
- Two leptons
- Two b-tags
- Signal/Noise > 70!



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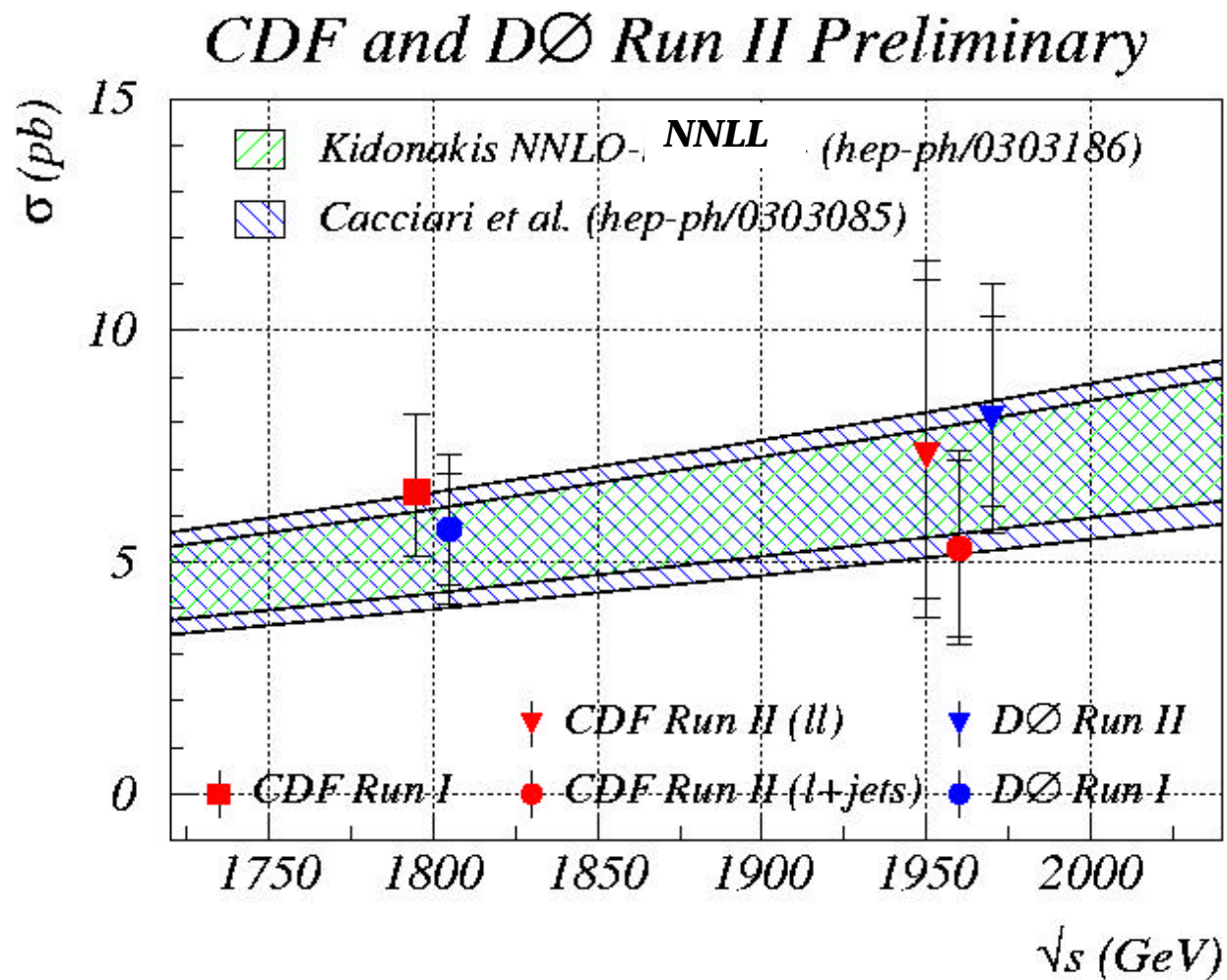
Top cross sections

$\text{D}\emptyset$ Run II Preliminary



146 pb ⁻¹	14.3 ^{+5.1+2.6} _{-4.3 -1.9}	pb
143 pb ⁻¹	7.2 ^{+2.6+1.6} _{-2.4 -1.7}	pb
93 pb ⁻¹	11.4 ^{+4.1+2.0} _{-3.5 -1.8}	pb
158 pb ⁻¹	11.1 ^{+5.8+1.4} _{-4.3 -1.4}	pb
164 pb ⁻¹	7.2 ^{+1.3+1.9} _{-1.2 -1.4}	pb
164 pb ⁻¹	8.2 ^{+1.3+1.9} _{-1.3 -1.6}	pb
162 pb ⁻¹	7.7 ^{+3.4+4.7} _{-3.3 -3.8}	pb

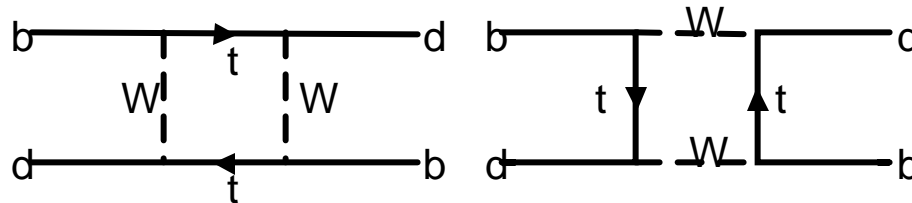
Cross section σ dependence



Equally Interesting: Top Quark Mass

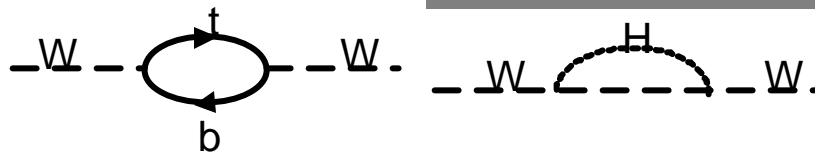
- Fundamental parameter of Standard Model (SM)
- Affects predictions of SM via radiative corrections:

- BB mixing



$$dM_W \propto m_t^2, (M_H)$$

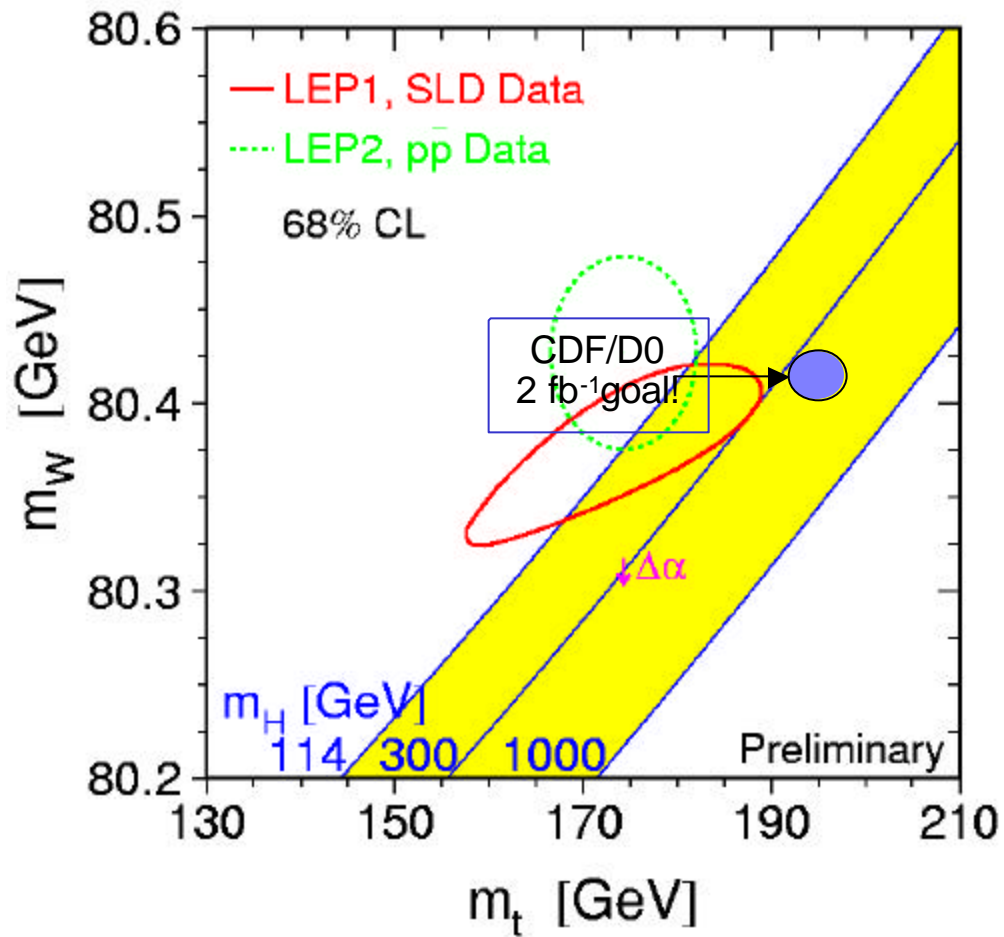
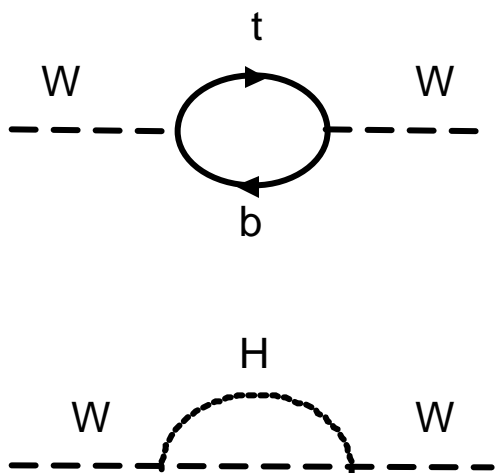
- W and Z mass



- measurements of M_W , m_t constrain M_H

- Large mass of top quark
 - Yukawa coupling $\gg 1$
 - may provide clues about electroweak symmetry breaking

Top Mass Constrains the Higgs

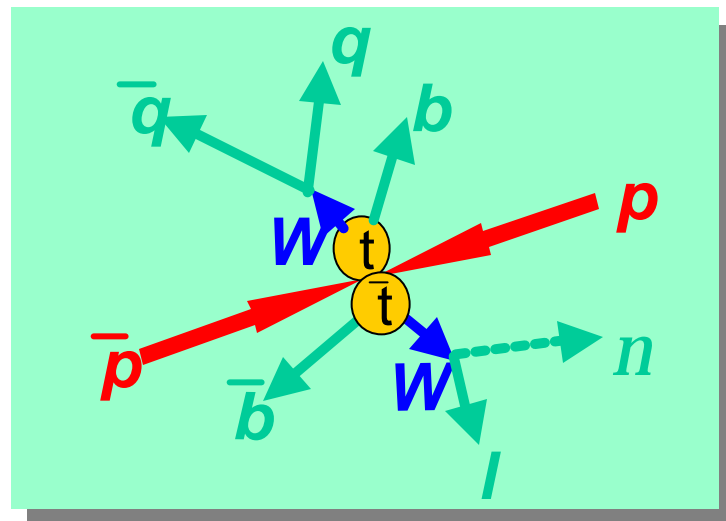


Mass: Lepton + Jets Channel

(there are other channels/methods)

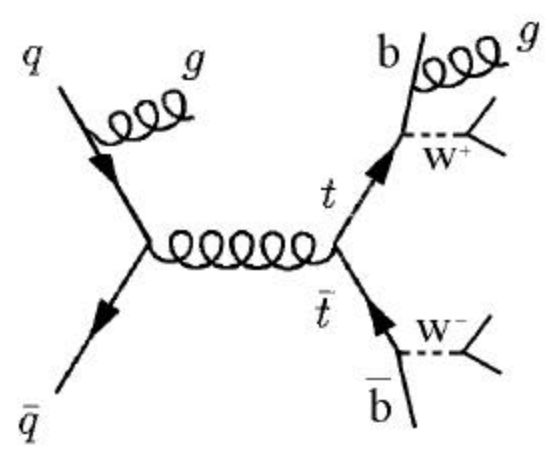
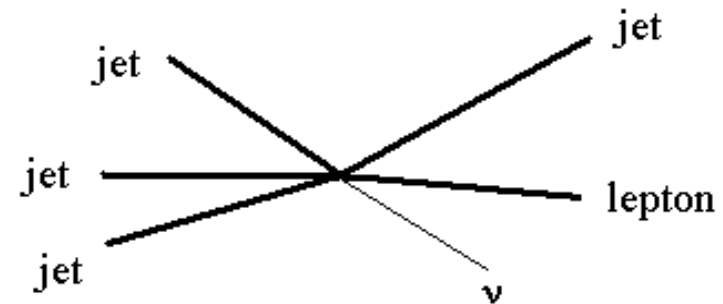
- 1 unknown (p_z^n)
- 3 constraints
 - $m(ln) = m(qq) = m_W$
 - $m(lnb) = m(qqb)$
- 2-constraint kinematic fit
- up to 24-fold combinatoric ambiguity
- compare to MC to measure m_t

$$t\bar{t} \text{ (R)} lnb\ qq\bar{q}\bar{b}$$



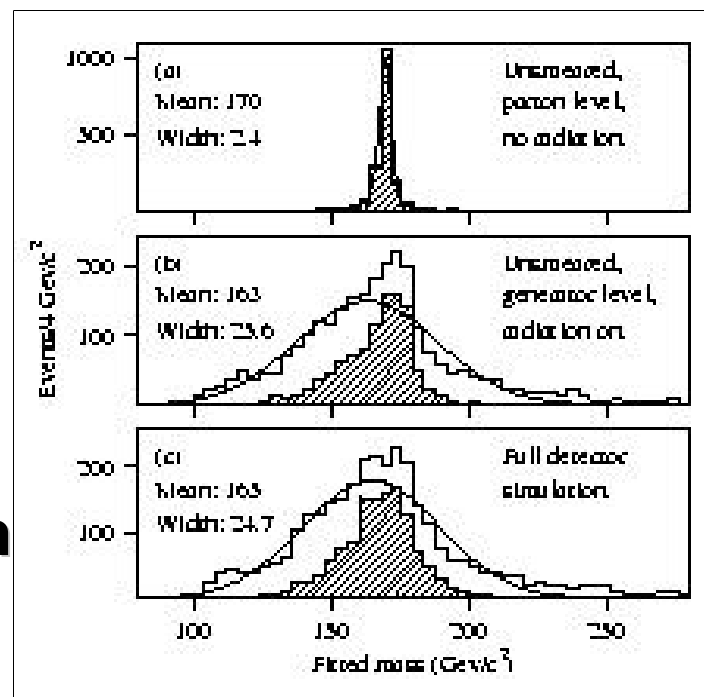
Mass: Complications

- **Combinatorics:**
 - 4 possible $j\bar{j}$ pairings
 - there are 12 possible assignments of the 4 jets to the 4 quarks (bbqq)
 - **only 6 if one of the jets is b-tagged**
 - **only 2 for events with double b-tagged jets**
 - **Homework 10:**
 - Work out the possibilities.
 - Although the em channel would have fewer combinatoric choices, why would it have less precision than lepton+jet channel?
- **Glueon radiation can add extra jets**



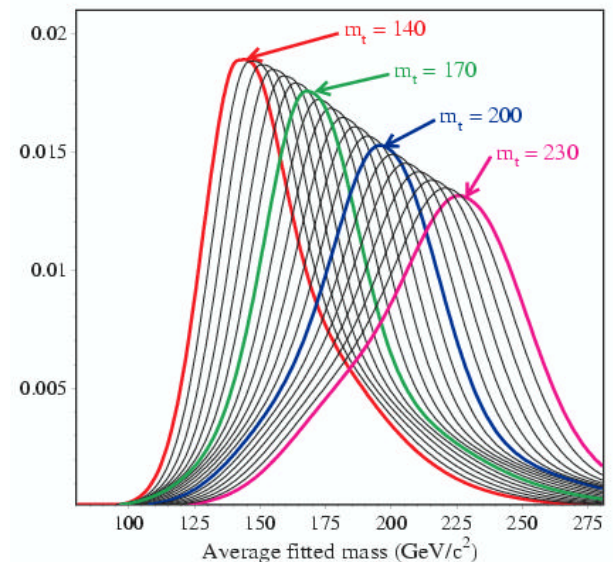
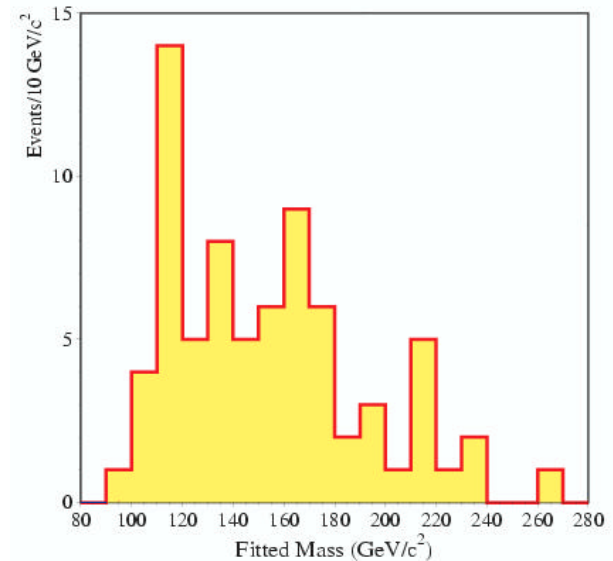
Mass: Combinatorics

- Monte Carlo tests:
 - shaded plots show correct combinations (Herwig MC, $m_t = 175$ GeV)
- The width and shape of the fitted mass distribution is due primarily to
 - jet combinatorics
 - QCD radiation



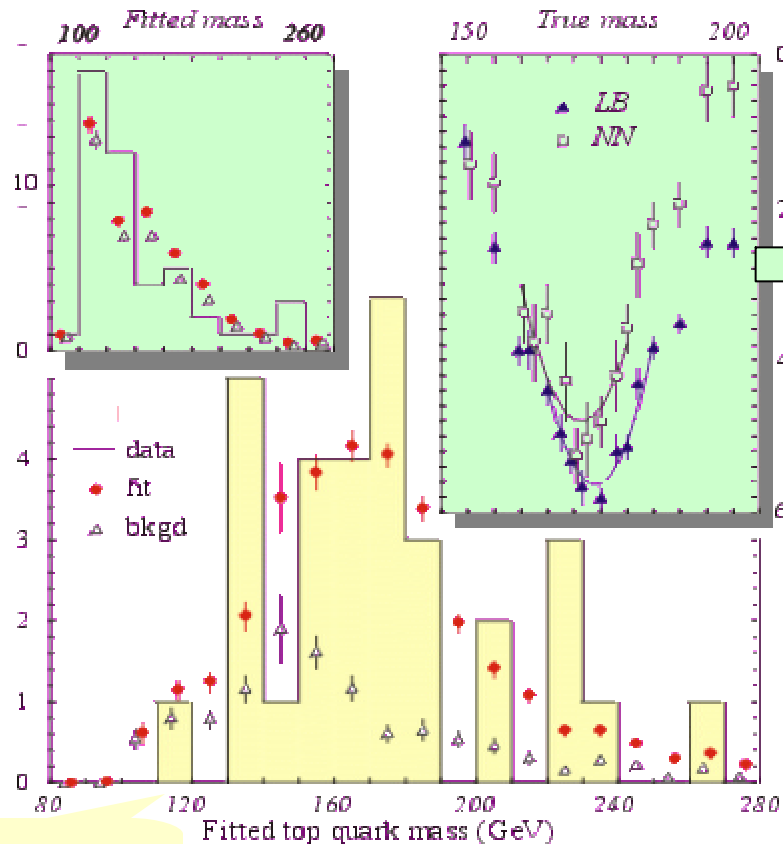
Mass: Basic Procedure

- In a sample of $t\bar{t}$ candidate events
 - For each candidate make a measurement of $X = f(m_t)$, where X is a suitable estimator for the top mass
 - e.g. result of the kinematic fit
 - This distribution contains signal and background.
- From MC determine shape of X as a function of m_t
 - Determine shape of X for background (MC & data).
 - Add these together and compare with data



Lepton + Jets Channel (DØ)

Background-rich sample



$$m_t = 173.3 \pm 5.6 \pm 5.5 \text{ GeV}$$

largest systematics

jet energy	4.0 GeV
MC generator	3.1 GeV
noise/pile-up	1.3 GeV

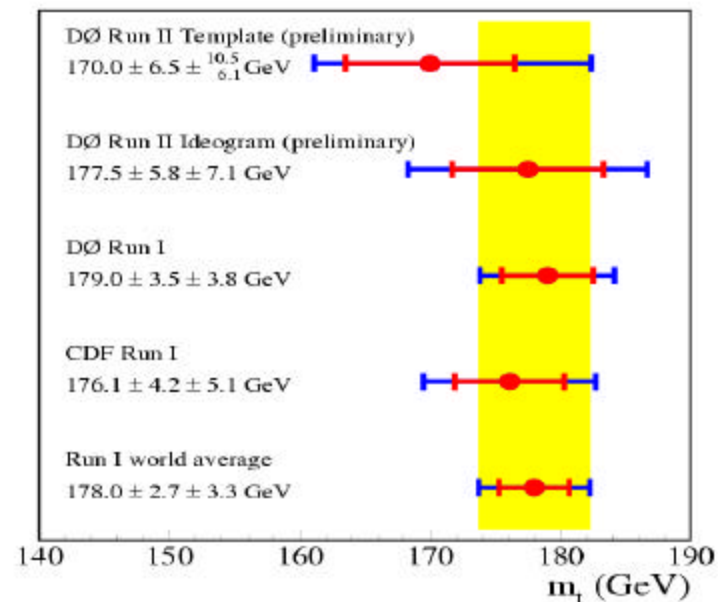
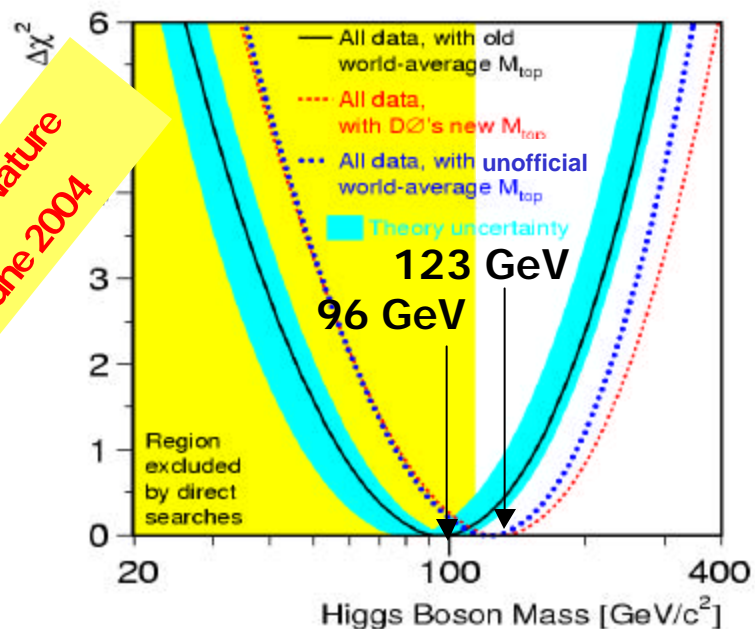
dominated by
jet energy scale and
gluon radiation

Signal-rich
sample

Run I and Run II Top mass

$m_{\text{top}} = 179.0 \pm 5.1 \text{ GeV}$ (DØ combined)

$m_{\text{top}} = 178.0 \pm 4.3 \text{ GeV}$ (official average)



First DZero Run II lepton + jets mass
160 pb⁻¹

Jerry Blazey
January 2005
Vietnam

A Disclaimer

- We've looked at QCD, Vector Boson, and Top Production.
- Just a sample of interesting Standard Model topics
- Many others:
 - Heavy Flavor
 - Di-boson Production
 - Diffractive Physics
 - And more....

Our Lecture Series

- ✓ **Overview**
- ✓ **History, Accelerators, Detectors, & Cross Sections**
- ✓ **QCD and Electroweak Physics**
- **Searches for New Physics**